# Engineering Functional Quantum Algorithms 

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#### Abstract

Suppose that a quantum circuit with $K$ elementary gates is known for a unitary matrix $U$, and assume that $U^{m}$ is a scalar matrix for some positive integer $m$. We show that a function of $U$ can be realized on a quantum computer with at most $O\left(m K+m^{2} \log m\right)$ elementary gates. The functions of $U$ are realized by a generic quantum circuit, which has a particularly simple structure. Among other results, we obtain efficient circuits for the fractional Fourier transform.


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Let $U$ be a unitary matrix, $U \in \mathcal{U}\left(2^{n}\right)$. Suppose that a fast quantum algorithm is known for $U$, which is given by a factorization of the form

$$
\begin{equation*}
U=U_{1} U_{2} \cdots U_{K} \tag{1}
\end{equation*}
$$

where the unitary matrices $U_{i}$ are realized by controllednot gates or by single qubit gates [1]. We are interested in the following question:

Are there efficient quantum algorithms for unitary matrices, which are functions of $U$ ?

The question is puzzling, because the knowledge of the factorization (II) of $U$ does not seem to be of much help in finding similar factorizations for, say, $V=U^{1 / 3}$. The purpose of this letter is to give an answer to the above question for a wide range of unitary matrices $U$.

Our solution to this problem is based on a generic circuit which implements arbitrary functions of $U$, assuming that $U^{m}$ is a scalar matrix for some positive integer $m$. If $m$ is small, then our method provides an efficient quantum circuit for $V$.

Notations. We denote by $\mathcal{U}(m)$ the group of unitary $m \times m$ matrices, by $\mathbf{1}$ the identity matrix, and by $\mathbf{C}$ the field of complex numbers.

## I. PRELIMINARIES

We recall some standard material on matrix functions, see [2] 3, 4] for more details. Let $U$ be a unitary matrix. The spectral theorem states that $U$ is unitarily equivalent to a diagonal matrix $D$, that is, $U=T D T^{\dagger}$ for some unitary matrix $T$. The elements $\lambda_{i}$ on the diagonal of $D=\operatorname{diag}\left(\lambda_{1}, \ldots, \lambda_{2^{n}}\right)$ are the eigenvalues of $U$.

[^0]Let $f$ be any function of complex scalars such that its domain contains the eigenvalues $\lambda_{i}, 1 \leq i \leq 2^{n}$. The matrix function $f(U)$ is then defined by

$$
f(U)=T \operatorname{diag}\left(f\left(\lambda_{1}\right), \ldots, f\left(\lambda_{2^{n}}\right)\right) T^{\dagger}
$$

where $T$ denotes the diagonalizing matrix of $U$, as above.
Notice that any two scalar functions $f$ and $g$, which take the same values on the spectrum of $U$, yield the same matrix value $f(U)=g(U)$. In particular, one can find an interpolation polynomial $g$, which takes the same values as $f$ on the eigenvalues $\lambda_{i}$. It is possible to assume that the degree of $g$ is smaller than the degree of the minimal polynomial of $U$. In other words, $V=f(U)$ can be expressed by a linear combination of integral powers of the matrix $U$,

$$
\begin{equation*}
V=f(U)=\sum_{i=0}^{m-1} \alpha_{i} U^{i} \tag{2}
\end{equation*}
$$

where $m$ is the degree of the minimal polynomial of the matrix $U$, and $\alpha_{i} \in \mathbf{C}$ for $i=0, \ldots, m-1$. In order for $V$ to be unitary, it is necessary and sufficient that the function $f$ maps the eigenvalues $\lambda_{i}$ of $U$ to elements on the unit circle.

Remark. There exist several different definitions for matrix functions. The relationshop between these definitions is discussed in detail in [5]. We have chosen the most general definition that allows to express the function values by polynomials.

## II. THE GENERIC CIRCUIT

Let $U$ be a unitary $2^{n} \times 2^{n}$ matrix with minimal polynomial of degree $m$. We assume that an efficient quantum circuit is known for $U$. How can we go about implementing the linear combination (2)? We will use an ancillary system of $\mu$ quantum bits, where $\mu$ is chosen such that $2^{\mu-1}<m \leq 2^{\mu}$ holds. This will allow us to create the linear combination by manipulating somewhat larger matrices, which on input $|0\rangle \otimes|\psi\rangle \in \mathbf{C}^{2^{\mu}} \otimes \mathbf{C}^{2^{n}}$ produce the state $|0\rangle \otimes V|\psi\rangle$.


FIG. 1: A quantum circuit realizing the block diagonal matrix $A=\operatorname{diag}\left(1, U, U^{2}, \ldots, U^{2^{\mu}-1}\right)$.

We first bring the ancillary system into a superposition of the first $m$ computational base states, such that an input state $|0\rangle \otimes|\psi\rangle \in \mathbf{C}^{2^{\mu}} \otimes \mathbf{C}^{2^{n}}$ is mapped to the state

$$
\begin{equation*}
\frac{1}{\sqrt{m}} \sum_{i=0}^{m-1}|i\rangle \otimes|\psi\rangle . \tag{3}
\end{equation*}
$$

This can be done by acting with a $2^{\mu} \times 2^{\mu}$ unitary matrix $B$ on the ancillary system, where the first column of $B$ is of the form $1 / \sqrt{m}(1, \ldots, 1,0, \ldots, 0)^{t}$. Efficient implementations of $B$ exist.

Notice that there exists an efficient implementation of the block diagonal matrix $A=\operatorname{diag}\left(1, U, U^{2}, \ldots, U^{2^{\mu}-1}\right)$. Indeed, $A$ can be composed of the matrices $U^{2^{\eta}}, 0 \leq$ $\eta<\mu$, conditioned on the $\mu$ ancillae bits. The resulting implementation is shown in Fig. 1. The state (3) is transformed by this circuit into the state

$$
\begin{equation*}
\frac{1}{\sqrt{m}} \sum_{i=0}^{m-1}|i\rangle \otimes U^{i}|\psi\rangle \tag{4}
\end{equation*}
$$

In the next step, we let a $2^{\mu} \times 2^{\mu}$ matrix $M$ act on the ancillae bits. We choose $M$ such that the state (4) is mapped to

$$
\begin{equation*}
\frac{1}{\sqrt{m}} \sum_{k=0}^{m-1}|k\rangle \otimes U^{k} V|\psi\rangle \tag{5}
\end{equation*}
$$

It turns out that $M$ can be realized by a unitary matrix, assuming that the minimal polynomial of $U$ is of the form $x^{m}-\tau, \tau \in \mathbf{C}$. This will be explained in some detail in the next section.

We apply the inverse $A^{\dagger}$ of the block diagonal matrix $A$. This transforms the state (5) to

$$
\begin{equation*}
\frac{1}{\sqrt{m}} \sum_{k=0}^{m-1}|k\rangle \otimes V|\psi\rangle \tag{6}
\end{equation*}
$$

We can clean up the ancillae bits by applying the $2^{\mu} \times 2^{\mu}$ matrix $B^{\dagger}$. This yields then the output state

$$
\begin{equation*}
|0\rangle \otimes V|\psi\rangle=|0\rangle \otimes f(U)|\psi\rangle \tag{7}
\end{equation*}
$$

The steps from the input state $|0\rangle \otimes|\psi\rangle$ to the final output state $|0\rangle \otimes V|\psi\rangle$ are illustrated in Fig. 2 for the case $\mu=2$.


FIG. 2: Generic circuit realizing a linear combination $V$. The case $\mu=2$ is shown.

The following theorem gives an upper bound on the complexity of the method. We use the number of elementary gates (that is, the number of single qubit gates and controlled-not gates) as a measure of complexity.

Theorem 1 Let $U$ be a $2^{n} \times 2^{n}$ unitary matrix with minimal polynomial $x^{m}-\tau, \tau \in \mathbf{C}$. Suppose that there exists a quantum algorithm for $U$ using $K$ elementary gates. Then a unitary matrix $V=f(U)$ can be realized with at most $O\left(m K+m^{2} \log m\right)$ elementary operations.

Proof. A matrix acting on $\mu \in O(\log m)$ qubits can be realized with at most $O\left(m^{2} \log m\right)$ elementary operations, cf. [1]. Therefore, the matrices $B, B^{\dagger}$, and $M$ can be realized with a total of at most $O\left(3 m^{2} \log m\right)$ operations.

If $K$ operations are needed to implement $U$, then at most $14 K$ operations are needed to implement $\Lambda_{1}(U)$, the operation $U$ controlled by a single qubit. The reason is that a doubly controlled NOT gate can be implemented with 14 elementary gates [6], and a controlled single qubit gate can be implemented with six or fewer elementary gates [1].

We observe that $2^{\mu}-1$ copies of $\Lambda_{1}(U)$ suffice to implement $A$. Indeed, we certainly can implement $\Lambda_{1}\left(U^{2^{k}}\right)$ by a sequence of $2^{k}$ circuits $\Lambda_{1}(U)$. This bold implementation yields the estimate for $A$. Typically, we will be able to find much more efficient implementations. Anyway, we can conclude that $A$ and $A^{\dagger}$ can both be implemented by at most $14\left(2^{\mu}-1\right) K \in O(14 m K)$ operations. Combining our counts yields the result.

## III. UNITARITY OF THE MATRIX $M$

It remains to show that the state (4) can be transformed into the state (5) by acting with a unitary matrix $M$ on the system of $\mu$ ancillae qubits. This is the crucial step in the previously described method.

Let $U$ be a unitary matrix with a minimal polynomial of degree $m$. A unitary matrix $V=f(U)$ can then be represented by a linear combination

$$
\begin{equation*}
V=\sum_{i=0}^{m-1} \alpha_{i} U^{i} \tag{8}
\end{equation*}
$$

We will motivate the construction of the matrix $M$ by examining in some detail the resulting linear combinations of the matrices $U^{k} V$. From (8), we obtain

$$
\begin{equation*}
U^{k} V=\sum_{i=0}^{m-1} \alpha_{i} U^{i+k} \tag{9}
\end{equation*}
$$

Suppose that the minimal polynomial of $U$ is of the form $m(x)=x^{m}-g(x)$, with $g(x)=\sum_{i=0}^{m-1} g_{i} x^{i}$. The right hand side of (9) can be reduced to a polynomial in $U$ of degree less than $m$ using the relation $U^{m}=g(U)$ :

$$
U^{k} V=\sum_{i=0}^{m-1} \beta_{k i} U^{i}
$$

The coefficients $\beta_{k i}$ are explicitly given by

$$
\left(\beta_{k 0}, \beta_{k 1}, \ldots, \beta_{k(m-1)}\right)=\left(\alpha_{0}, \alpha_{1}, \ldots, \alpha_{m-1}\right) P^{k}
$$

where $P$ denotes the companion matrix of $m(x)$, that is,

$$
P=\left(\begin{array}{ccccc}
0 & 1 & 0 & \cdots & 0 \\
0 & 0 & 1 & \cdots & 0 \\
\vdots & \vdots & \vdots & \ddots & \vdots \\
0 & 0 & 0 & \cdots & 1 \\
g_{0} & g_{1} & g_{2} & \cdots & g_{m-1}
\end{array}\right)
$$

The $2^{\mu} \times 2^{\mu}$ matrix $M$ is defined by

$$
M=\left(\begin{array}{cc}
C & 0 \\
0 & \mathbf{1}
\end{array}\right)
$$

where $C=\left(\beta_{k i}\right)_{k, i=0, \ldots, m-1}$, and $\mathbf{1}$ is a $\left(2^{\mu}-m\right) \times\left(2^{\mu}-m\right)$ identity matrix. Under the assumptions of Theorem 1, it turns out that the matrix $M$ is unitary. Before proving this claim, let us formally check that the matrix $M$ transforms the state (4) into the state (5). If we apply the matrix $M$ to the ancillary system, then we obtain from (4) the state

$$
\begin{aligned}
\frac{1}{\sqrt{m}} \sum_{i=0}^{m-1} M|i\rangle & \otimes U^{i}|\psi\rangle=\frac{1}{\sqrt{m}} \sum_{k, i=0}^{m-1} \beta_{k i}|k\rangle \otimes U^{i}|\psi\rangle \\
& =\frac{1}{\sqrt{m}} \sum_{k=0}^{m-1}|k\rangle \otimes \sum_{i=0}^{m-1} \beta_{k i} U^{i}|\psi\rangle \\
& =\frac{1}{\sqrt{m}} \sum_{k=0}^{m-1}|k\rangle \otimes U^{k} V|\psi\rangle
\end{aligned}
$$

which coincides with (5), as claimed.

Lemma 2 Let $U$ be a unitary matrix with minimal polynomial $m(x)=x^{m}-\tau$. Let $V$ be a matrix satisfying (因). If $V$ is unitary, then $M$ is unitary.

Proof. It suffices to show that the matrix $C$ is unitary. Notice that the assumption on the minimal polynomial $m(x)$ implies that $C$ is of the form

$$
C=\left(\begin{array}{ccccc}
\alpha_{0} & \alpha_{1} & \cdots & \alpha_{m-2} & \alpha_{m-1} \\
\tau \alpha_{m-1} & \alpha_{0} & \cdots & \alpha_{m-3} & \alpha_{m-2} \\
\ddots & \ddots & & \ddots & \ddots \\
\tau \alpha_{1} & \tau \alpha_{2} & \cdots & \tau \alpha_{m-1} & \alpha_{0}
\end{array}\right)
$$

that is, $C$ is obtained from a circulant matrix by multiplying every entry below the diagonal by $\tau$. In other words, we have

$$
C=\left([\tau]_{i>j} \alpha_{j-i \bmod m}\right)_{i, j=0, \ldots, m-1}
$$

where $[\tau]_{i>j}=\tau$ if $i>j$, and $[\tau]_{i>j}=1$ otherwise.
Note that the inner product of row $a$ with row $b$ of matrix $C$ is the same as the inner product of row $a+1$ with row $b+1$. Thus, to prove the unitarity of $C$, it suffices to show that

$$
\begin{equation*}
\delta_{a, 0} \stackrel{!}{=}\langle\text { row a|row } 0\rangle=\sum_{j=0}^{a-1} \bar{\tau} \overline{\alpha_{j-a}} \alpha_{j}+\sum_{j=a}^{m-1} \overline{\alpha_{j-a}} \alpha_{j} \tag{10}
\end{equation*}
$$

holds, where $\delta_{a, 0}$ denotes the Kronecker delta and the indices of $\alpha$ are understood modulo $m$.

Consider the equation

$$
\begin{equation*}
\mathbf{1}=V^{\dagger} V=\left(\sum_{i=0}^{m-1} \overline{\alpha_{i}} U^{-i}\right)\left(\sum_{i=0}^{m-1} \alpha_{i} U^{i}\right) \tag{11}
\end{equation*}
$$

The right hand side can be simplified to a polynomial in $U$ of degree less than $m$ using the identity $\bar{\tau} U^{m}=\mathbf{1}$. The coefficient of $U^{a}$ in (11) is exactly the right hand side of equation (10). Since the minimal polynomial of $U$ is of degree $m$, it follows that the matrices $U^{0}, U^{1}, \ldots, U^{m-1}$ are linearly independent. Thus, comparing coefficients on both sides of equation (11) shows (10). Hence the rows of $C$ are pairwise orthogonal and of unit norm.

A Simple Example. Let $F_{n}$ be the discrete Fourier transform matrix

$$
F_{n}=2^{-n / 2}\left(\exp \left(-2 \pi i k \ell / 2^{n}\right)\right)_{k, \ell=0, \ldots, 2^{n}-1}
$$

with $i^{2}=-1$. Recall that the Cooley-Tukey decomposition yields a fast quantum algorithm, which implements $F_{n}$ with $O\left(n^{2}\right)$ elementary operations. The minimal polynomial of $F_{n}$ is $x^{4}-1$ if $n \geq 3$. Thus, any unitary matrix $V$, which is a function of $F_{n}$, can be realized with $O\left(n^{2}\right)$ operations.

For instance, if $n \geq 3$, then the fractional power $F_{n}^{x}$, $x \in \mathbf{R}$, can be expressed by

$$
F_{n}^{x}=\alpha_{0}(x) I+\alpha_{1}(x) F_{n}+\alpha_{2}(x) F_{n}^{2}+\alpha_{3}(x) F_{n}^{3}
$$

where the coefficients $\alpha_{i}(x)$ are given by (cf. 7):

$$
\begin{array}{ll}
\alpha_{0}(x)=\frac{1}{2}\left(1+e^{i x}\right) \cos x, & \alpha_{1}(x)=\frac{1}{2}\left(1-i e^{i x}\right) \sin x \\
\alpha_{2}(x)=\frac{1}{2}\left(-1+e^{i x}\right) \cos x, & \alpha_{3}(x)=\frac{1}{2}\left(-1-i e^{i x}\right) \sin x
\end{array}
$$

In this case, $F_{n}^{x}$ is realized by the circuit in Fig. 2 with $U=F_{n}$ and $M=\left(\alpha_{j-i}(x)\right)_{i, j=0, \ldots, 3}$. The circuit can be implemented with $O\left(n^{2}\right)$ operations.

## IV. LIMITATIONS

The previous sections showed that a unitary matrix $f(U)$ can be realized by a linear combination of the powers $U^{i}$, $0 \leq i<m$, if the minimal polynomial $m(x)$ of $U$ is of the form $x^{m}-\tau, \tau \in \mathbf{C}$. One might wonder whether the restriction to minimal polynomials of this form is really necessary. The next lemma explains why we had this limitation:

Lemma 3 Let $U$ be a unitary matrix with minimal polynomial $m(x)=x^{m}-g(x), \operatorname{deg} g(x)<m$. If $g(x)$ is not a constant, then the matrix $M$ is in general not unitary.

Proof. Suppose that $g(x)=\sum_{i=0}^{m-1} g_{i} x^{i}$. We may choose for instance $V=U^{m}=g(U)$. Then the norm of first row in $M$ is greater than 1 . Indeed, we can calculate this norm to be $\left|g_{0}\right|^{2}+\left|g_{1}\right|^{2}+\cdots+\left|g_{m-1}\right|^{2}$. However, $\left|g_{0}\right|^{2}=1$, because $g_{0}$ is a product of eigenvalues of $U$. By assumption, there is another nonzero coefficient $g_{i}$, which proves the result.

## V. EXTENSIONS

We describe in this section one possibility to extend our approach to a larger class of unitary matrices $U$. We assumed so far that $f(U)$ is realized by a linear combination (2) of linearly independent matrices $U^{i}$. The exponents were restricted to the range $0 \leq i<m$, where $m$ is degree of the minimal polynomial of $U$. We can circumvent the problem indicated in the previous section by allowing $m$ to be larger than the degree of the minimal polynomial.

Theorem 4 Let $U \in \mathcal{U}\left(2^{n}\right)$ be a unitary matrix such that $U^{m}$ is a scalar matrix for some positive integer $m$. Suppose that there exists a quantum circuit which implements $U$ with $K$ elementary gates. Then a unitary matrix $V=f(U)$ can be realized with $O\left(m K+m^{2} \log m\right)$ elementary operations.

Proof. By assumption, $U^{m}=\tau \mathbf{1}$ for some $\tau \in \mathbf{C}$. This means that the minimal polynomial $m(x)$ of $U$ divides the polynomial $x^{m}-\tau$, that is, $x^{m}-\tau=m(x) m_{2}(x)$ for some $m_{2}(x) \in \mathbf{C}[x]$.

We may assume without loss of generality that the function $f$ is defined at all roots of $x^{m}-\tau$. Indeed, we can replace $f$ by an interpolation polynomial $g$ satisfying $f(U)=g(U)$ if this is necessary.

Choose any unitary matrix $A \in U\left(2^{n}\right)$ with minimal polynomial $m_{2}(x)$. The minimal polynomial of the block diagonal matrix $U_{A}=\operatorname{diag}(U, A)$ is $x^{m}-\tau$, the least common multiple of the polynomials $m(x)$ and $m_{2}(x)$. Express $f\left(U_{A}\right)$ by powers of the block diagonal matrix $U_{A}$ :

$$
\begin{equation*}
f\left(U_{A}\right)=\operatorname{diag}(f(U), f(A))=\sum_{i=0}^{m-1} \alpha_{i} \operatorname{diag}\left(U^{i}, A^{i}\right) \tag{12}
\end{equation*}
$$

The approach detailed in Section III yields a unitary matrix $M$ to realize this linear combination. On the other hand, we obtain from (12) the relation

$$
f(U)=\sum_{i=0}^{m-1} \alpha_{i} U^{i}
$$

by ignoring the auxiliary matrices $A^{i}, 0 \leq i<m$. It is clear that a circuit of the type shown in Fig. 2 with $\mu$ chosen such that $2^{\mu-1}<m \leq 2^{\mu}$ implements this linear combination of the matrices $\bar{U}^{i}, 0 \leq i<m$, provided we use the matrix $M$ constructed above.

## VI. CONCLUSIONS

Few methods are currently known that facilitate the engineering of quantum algorithms. Linear algebra allowed us to derive efficient quantum circuits for $f(U)$, given an efficient quantum circuit for $U$, as long as $U^{m}$ is a scalar matrix for some small integer $m$. This method can be used in conjuction with the Fourier sampling techniques by Shor $\|8\|$, the eigenvalue estimation technique by Kitaev [9], and the probability amplitude amplification method by Grover [10], to design more elaborate quantum algorithms.
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