

# Energy as a detector of nonlocality of many-body spin systems

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Seattle, Washington



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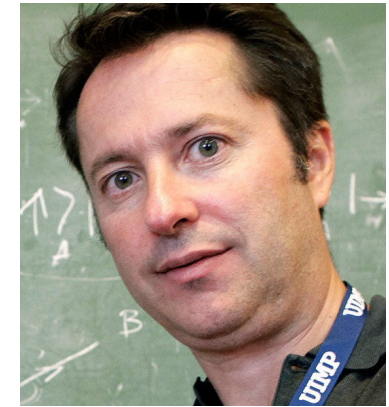
joint work with



Gemma de las Cuevas



Maciej Lewenstein



J. Ignacio Cirac



Remigiusz Augusiak



Antonio Acín

The paper is  
available on  
[arXiv:1607.06090]  
(with referees in  
Phys. Rev. X)



# Outline



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- Motivation





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- Motivation
- The idea, the setting



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- Motivation
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- Quantum optimization



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- Assigning a Bell inequality to a Hamiltonian



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- Conclusions and outlook





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- **Bell correlations are stronger than entanglement**



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[Tura et al, Science **344** 1256 (2014), Schmied et al, Science **352** 441(2016)]



# Why Bell correlations in the many-body regime?

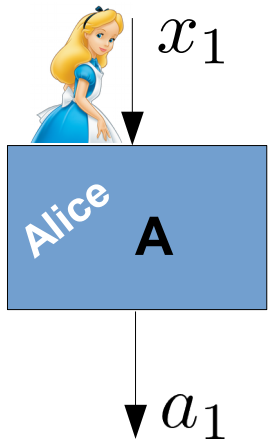
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- This talk: spin systems in one spatial dimension



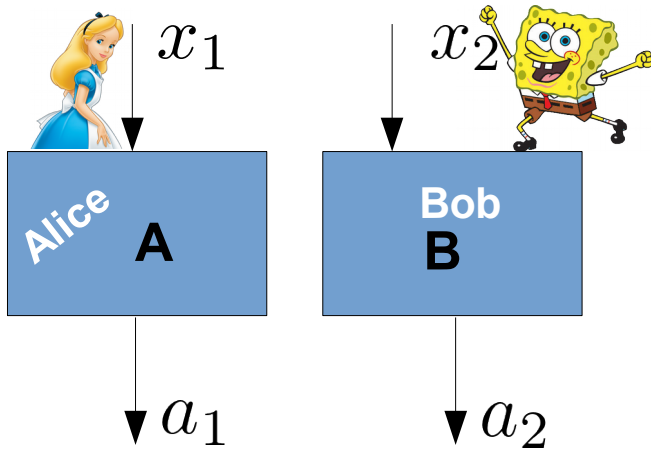
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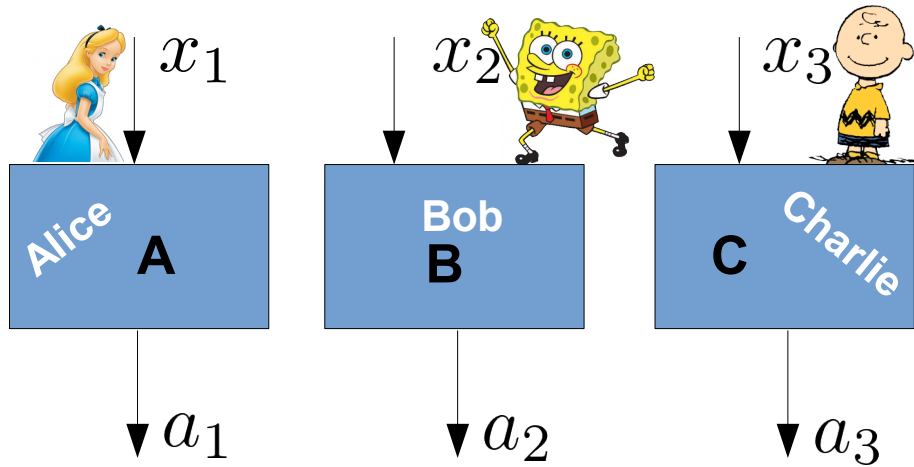


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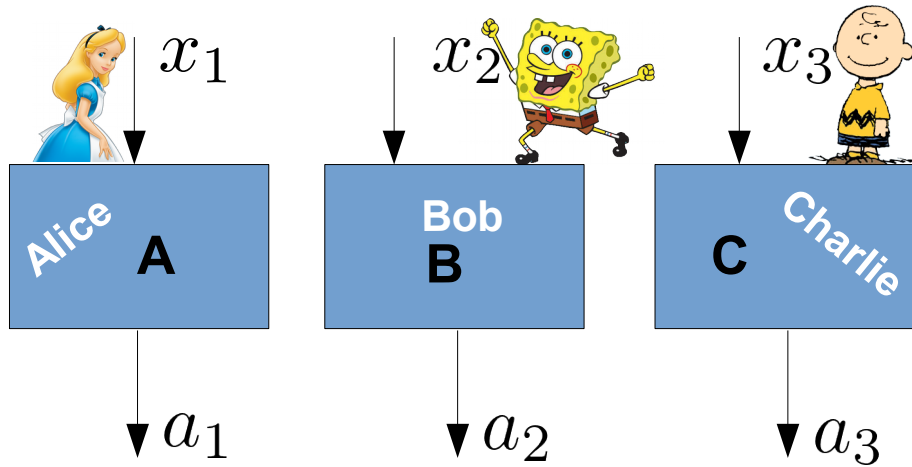




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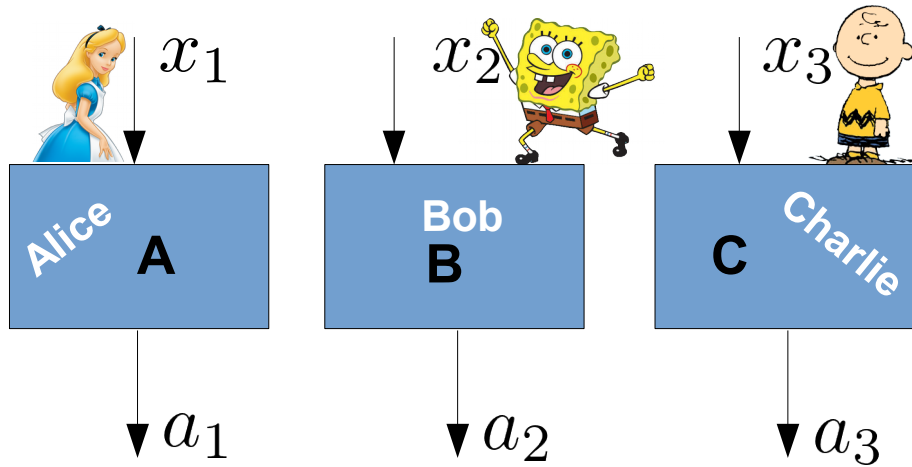


# *A crash course on nonlocality*



$$P(a_1 \dots a_n | x_1 \dots x_n)$$

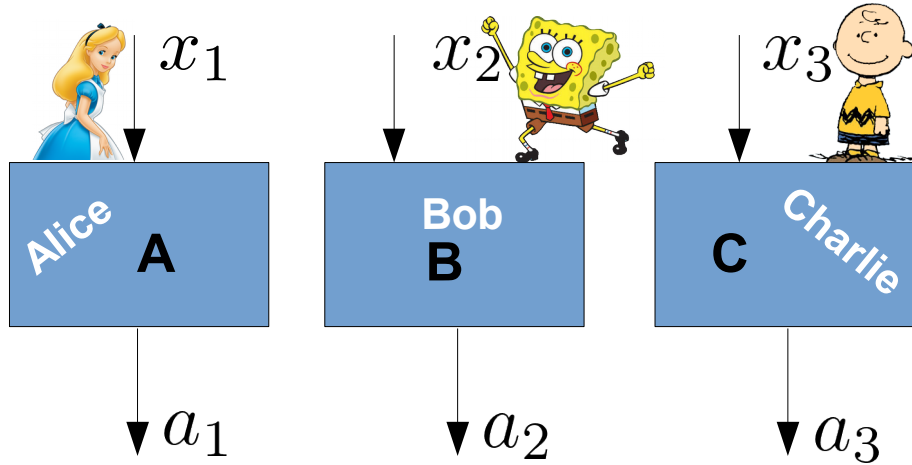
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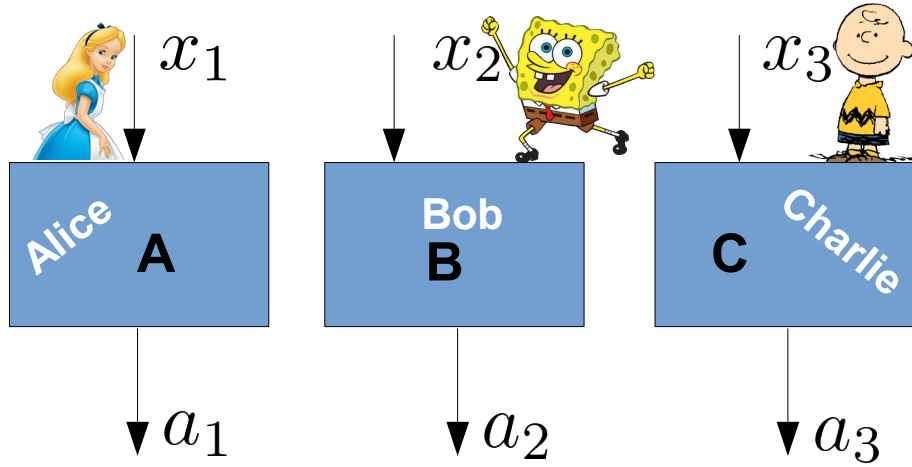
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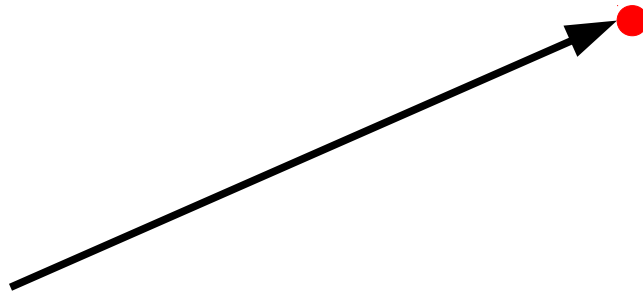
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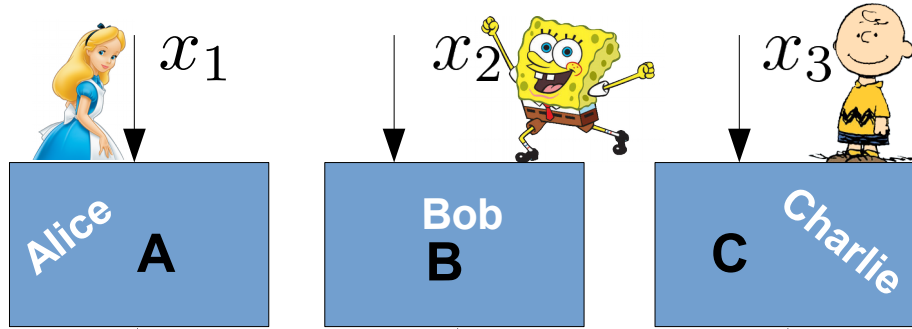
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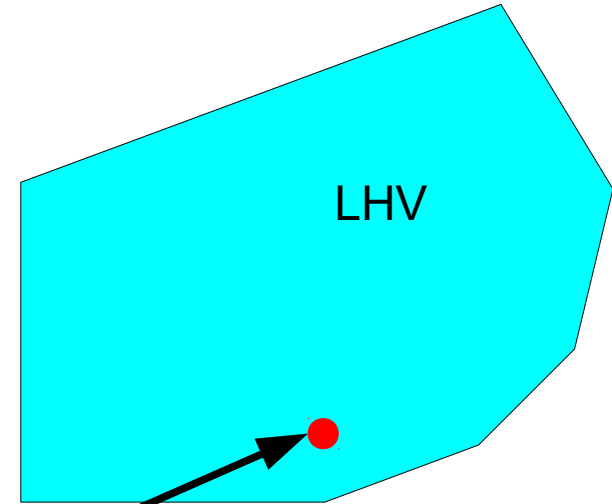
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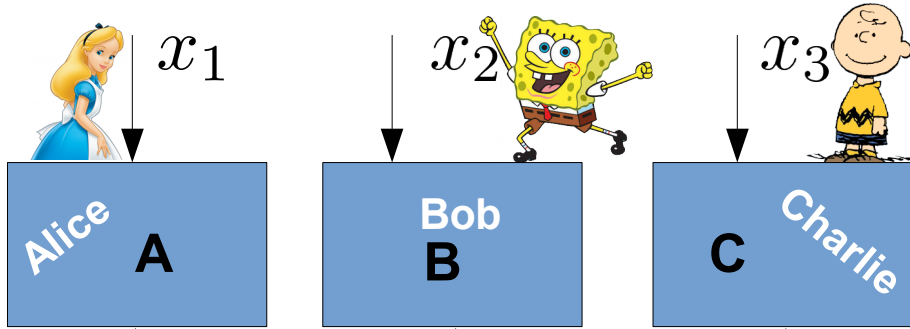


Local Polytope

$\mathbb{P}_L$



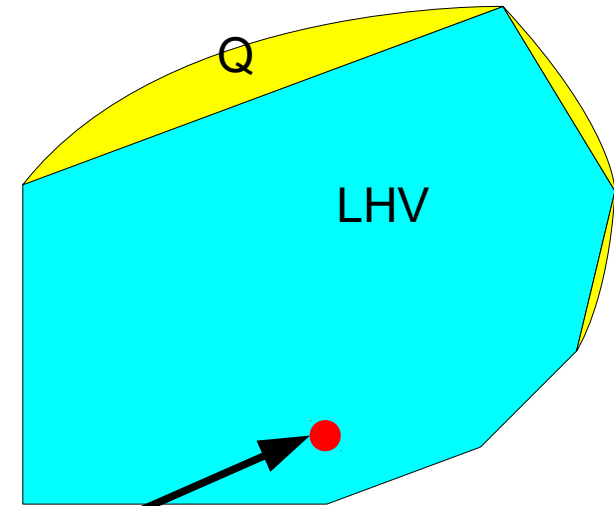
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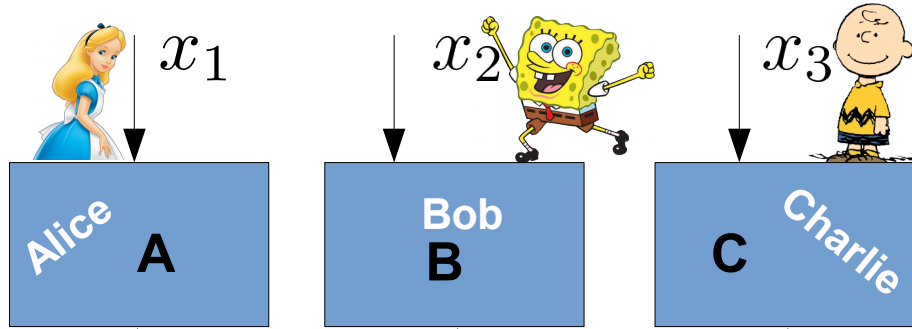


Local Polytope  $\subset$  Quantum Set

$$\mathbb{P}_L \subset Q$$



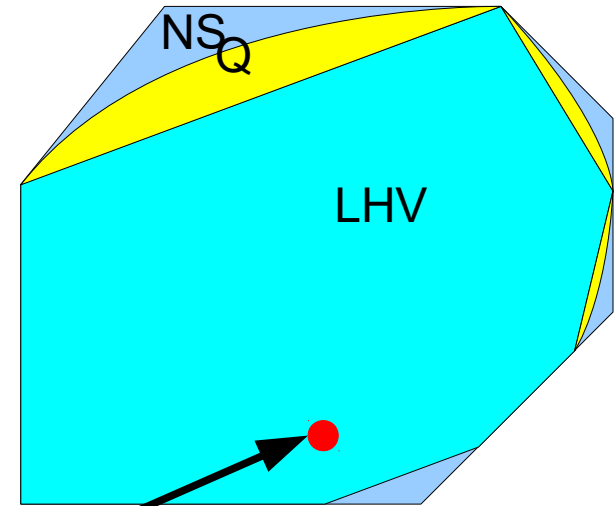
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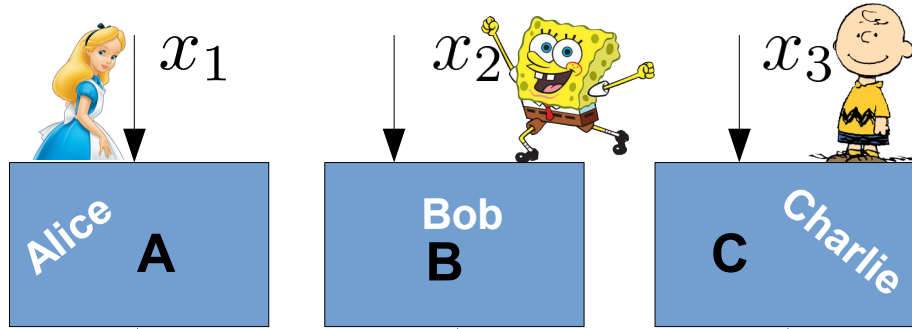


Local Polytope  $\subset$  Quantum Set  $\subset$  NS Polytope  
 $\mathbb{P}_L$                        $\mathcal{Q}$                        $\mathbb{P}_{NS}$





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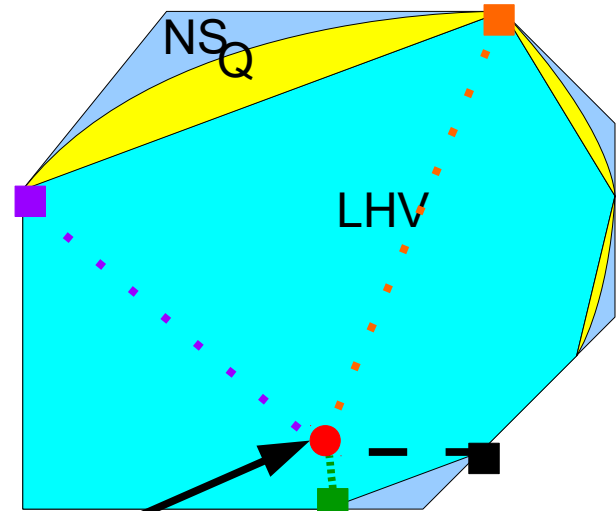


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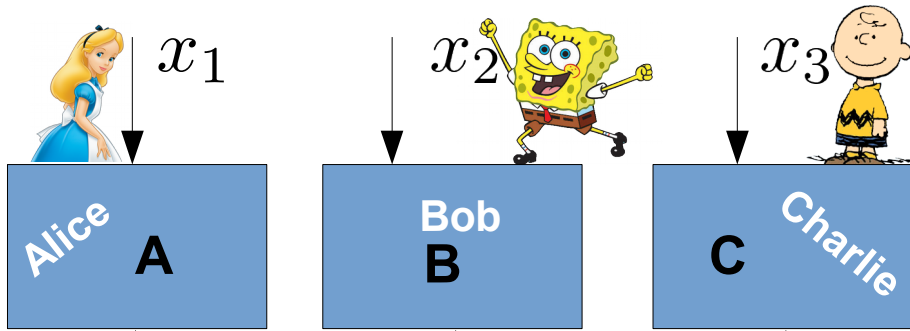


Example:

**Charlie's Instructions**  
 $\lambda = \{1, 3, 1, 2, 4, 3, 1, 1, \dots\}$   
**Output**  
 $0, x_3, 0, 1, \overline{x_3}, x_3, 0, 0, \dots$



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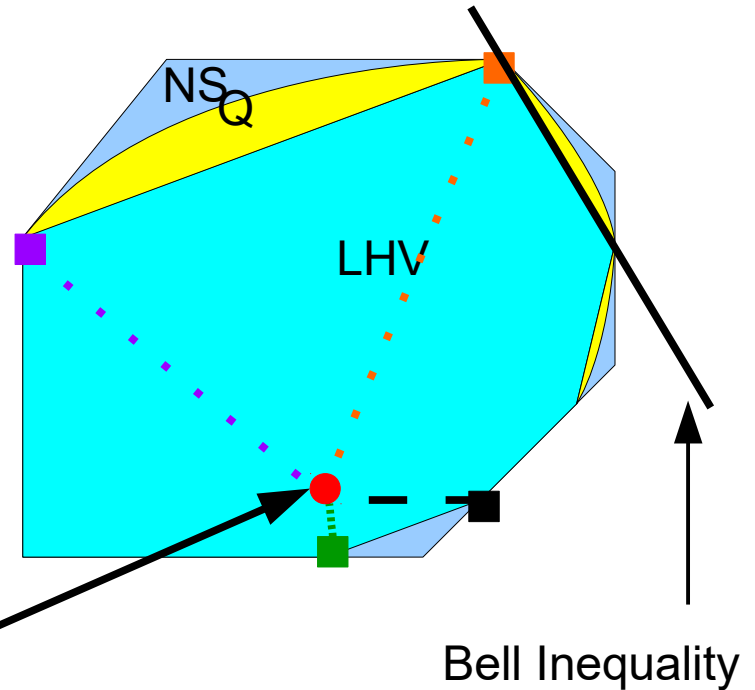


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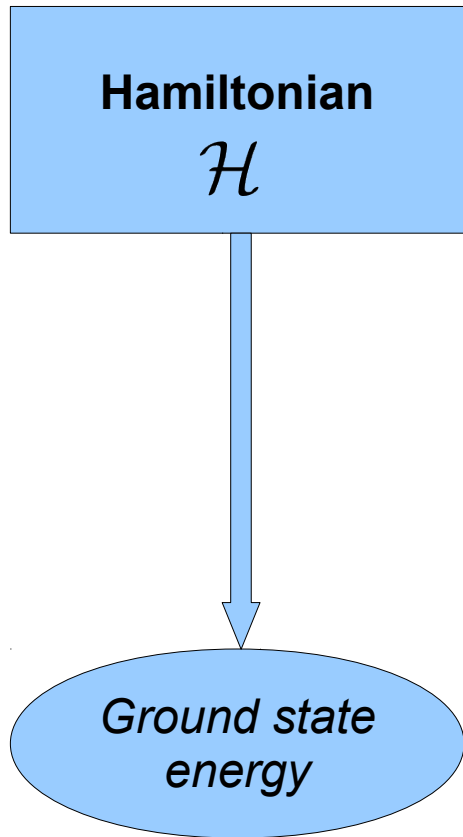
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**Hamiltonian**

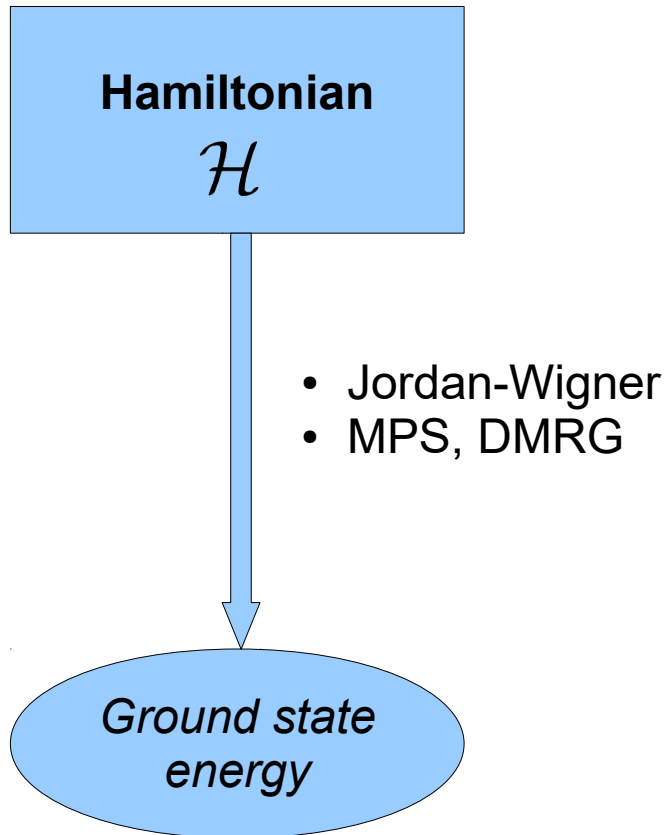
$\mathcal{H}$



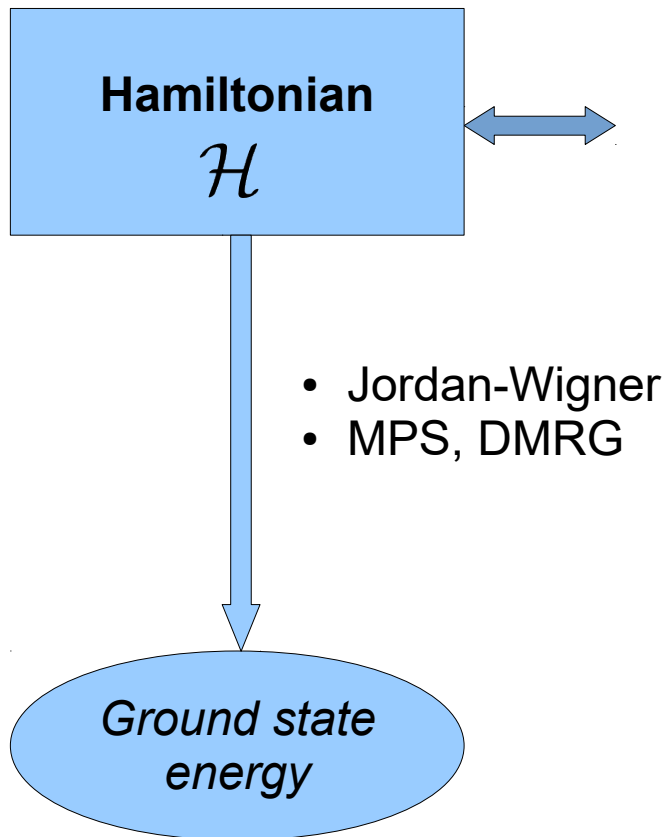
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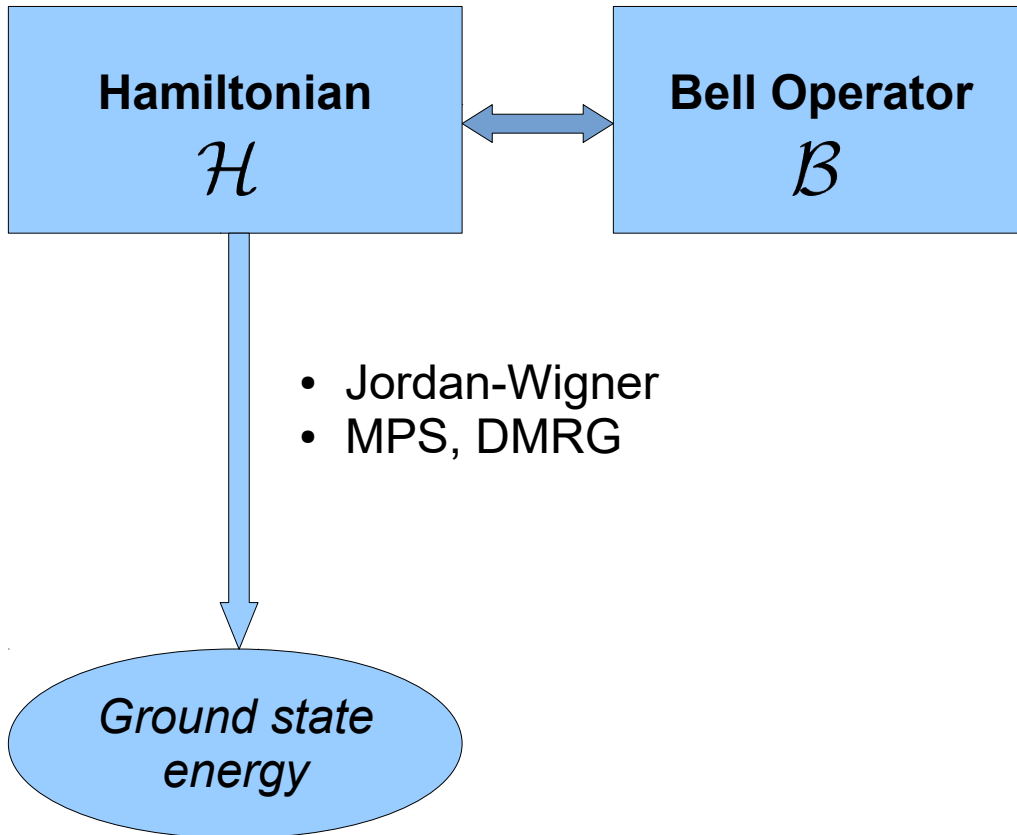
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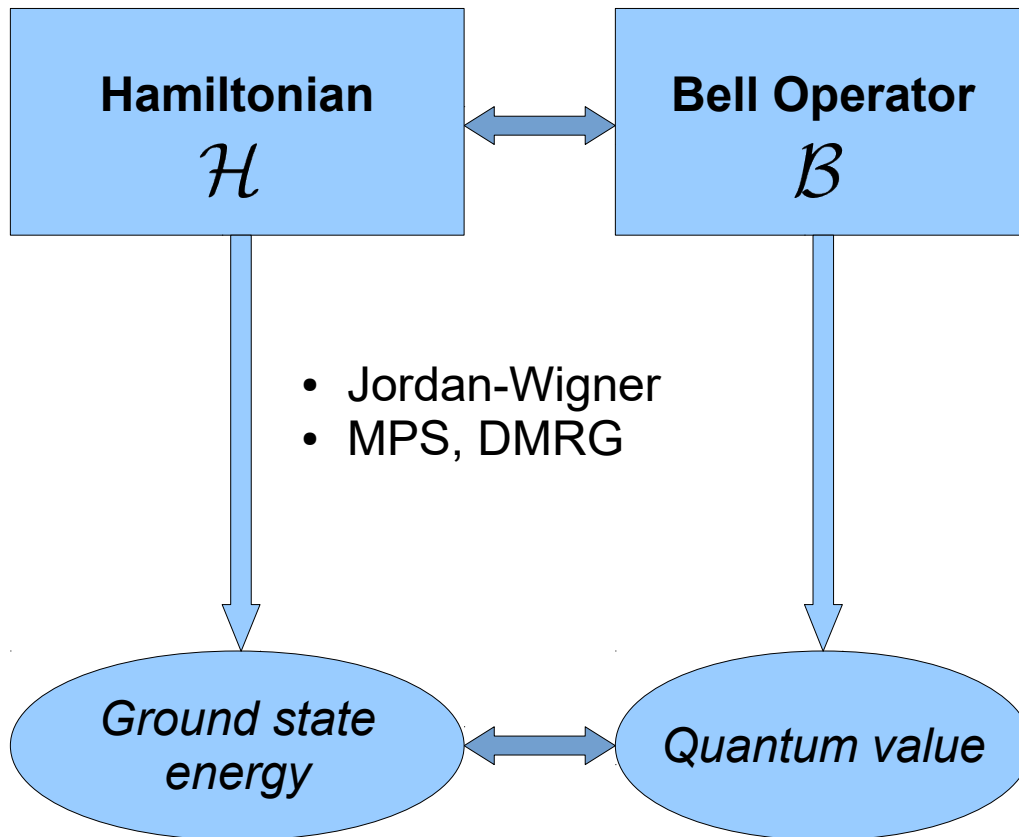


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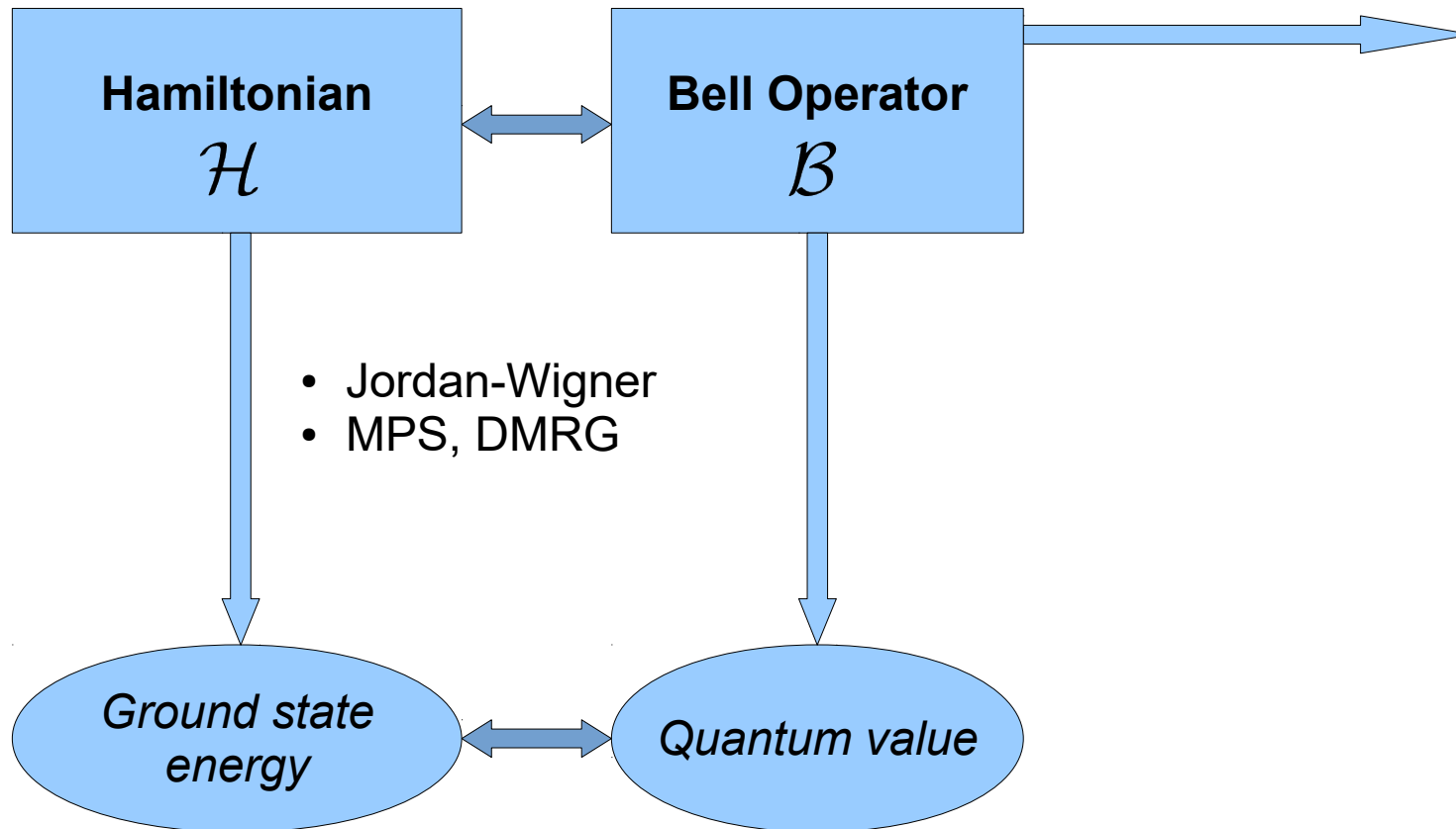




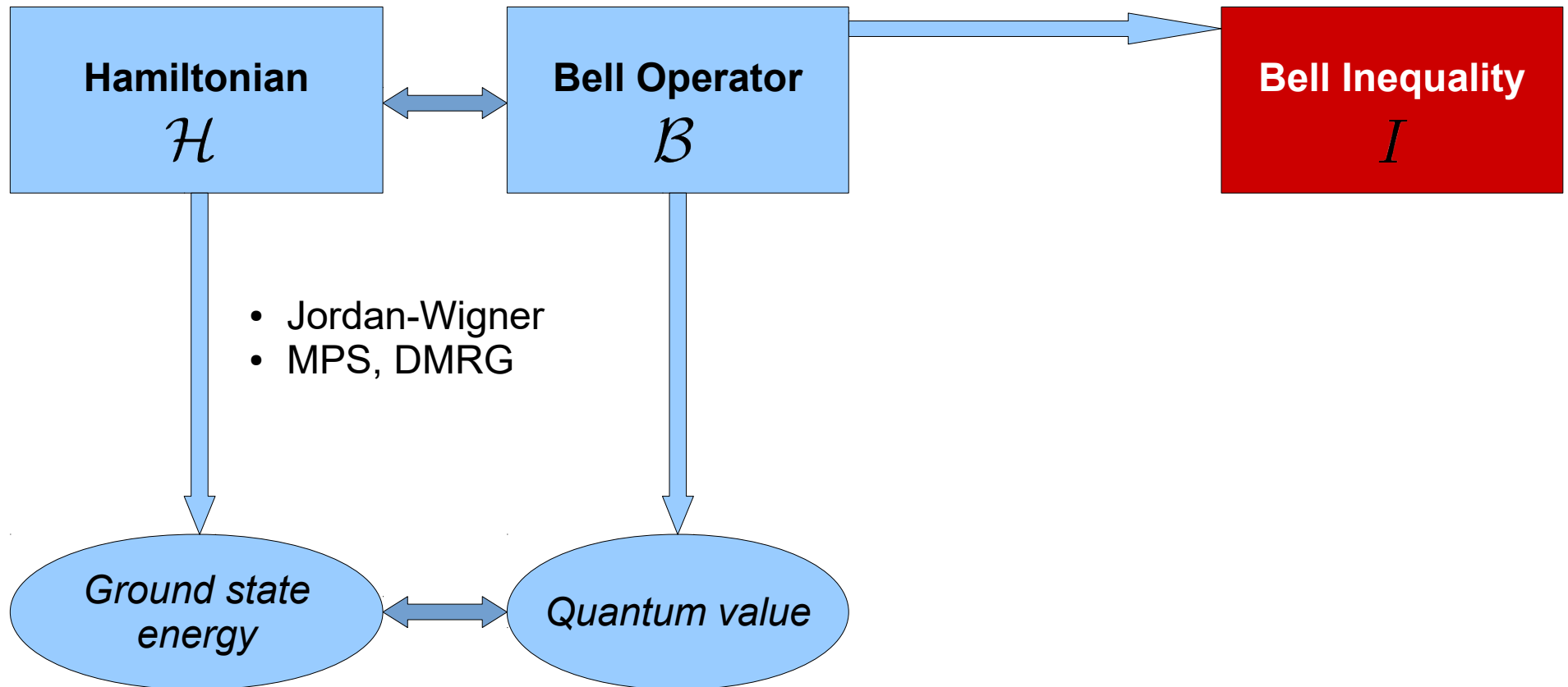
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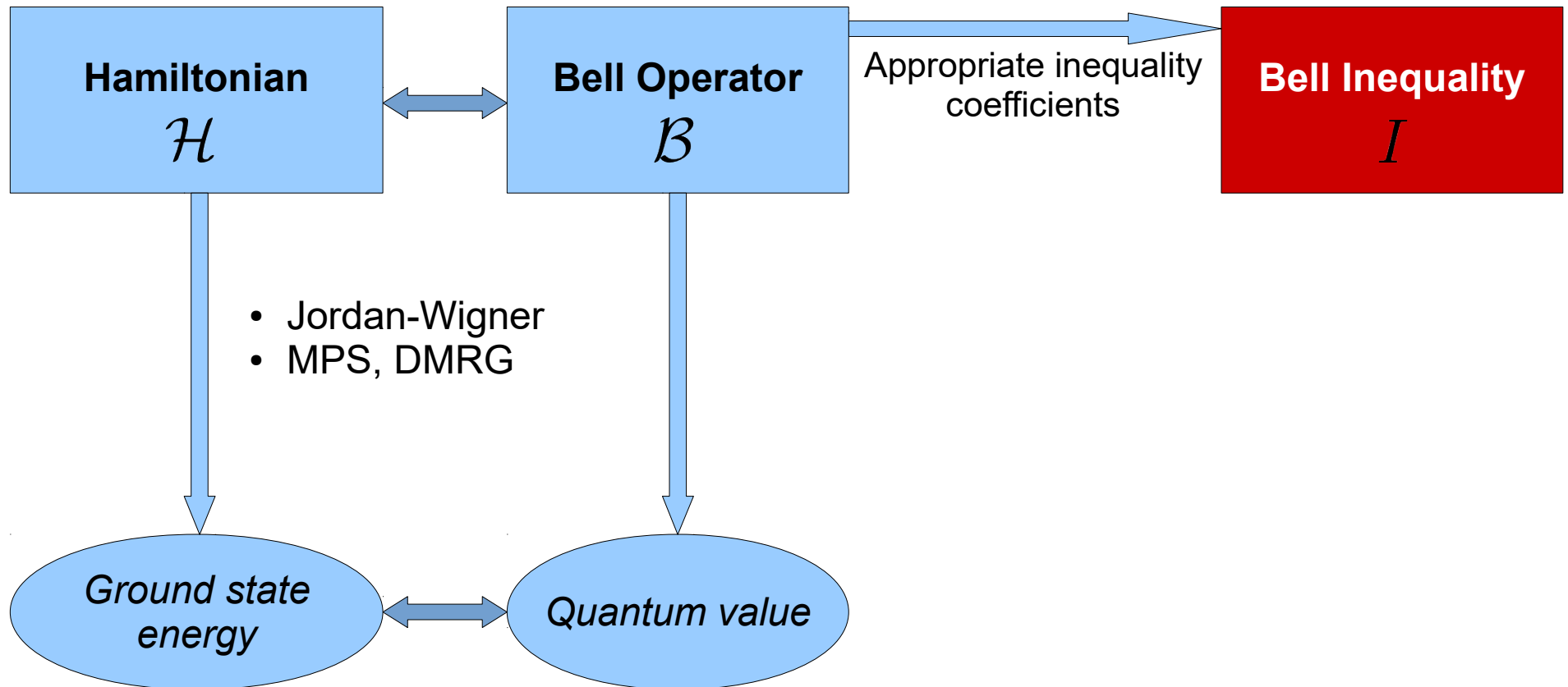
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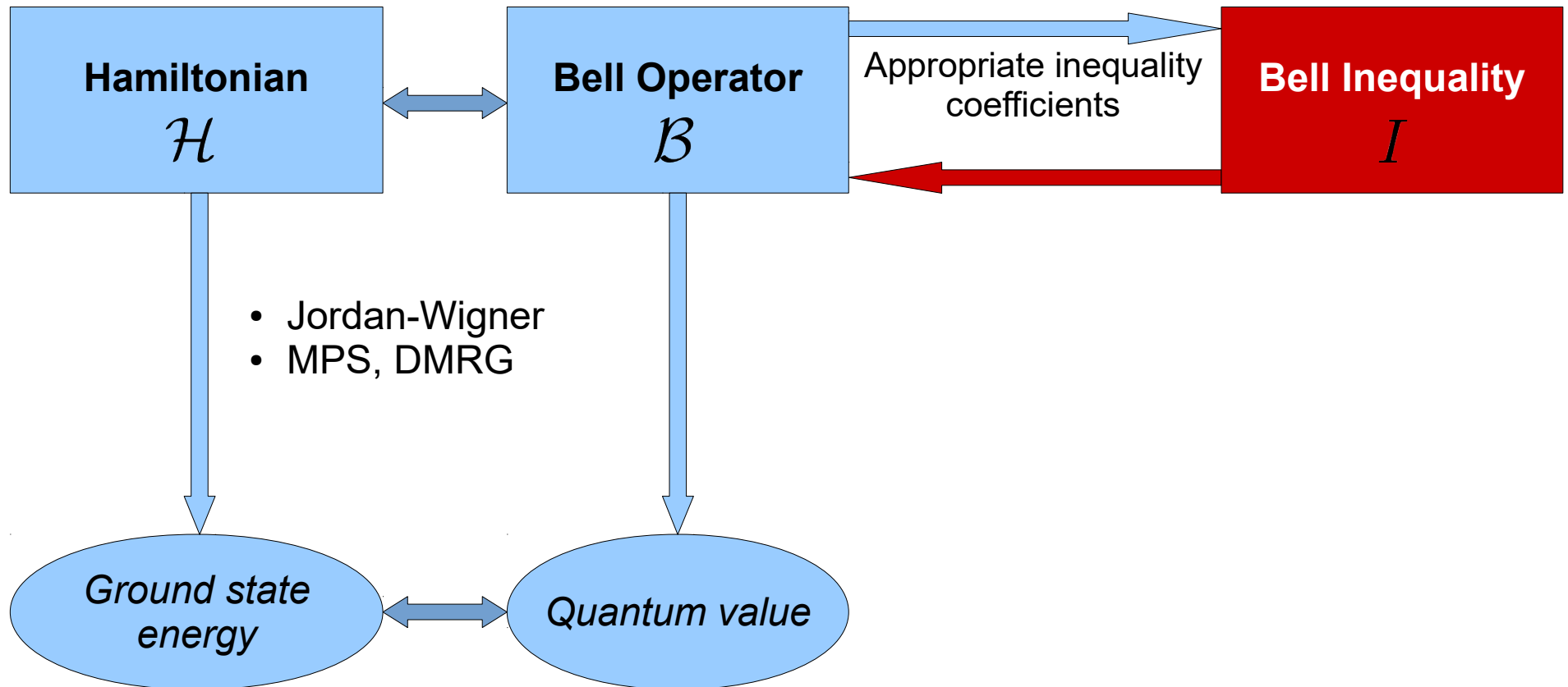
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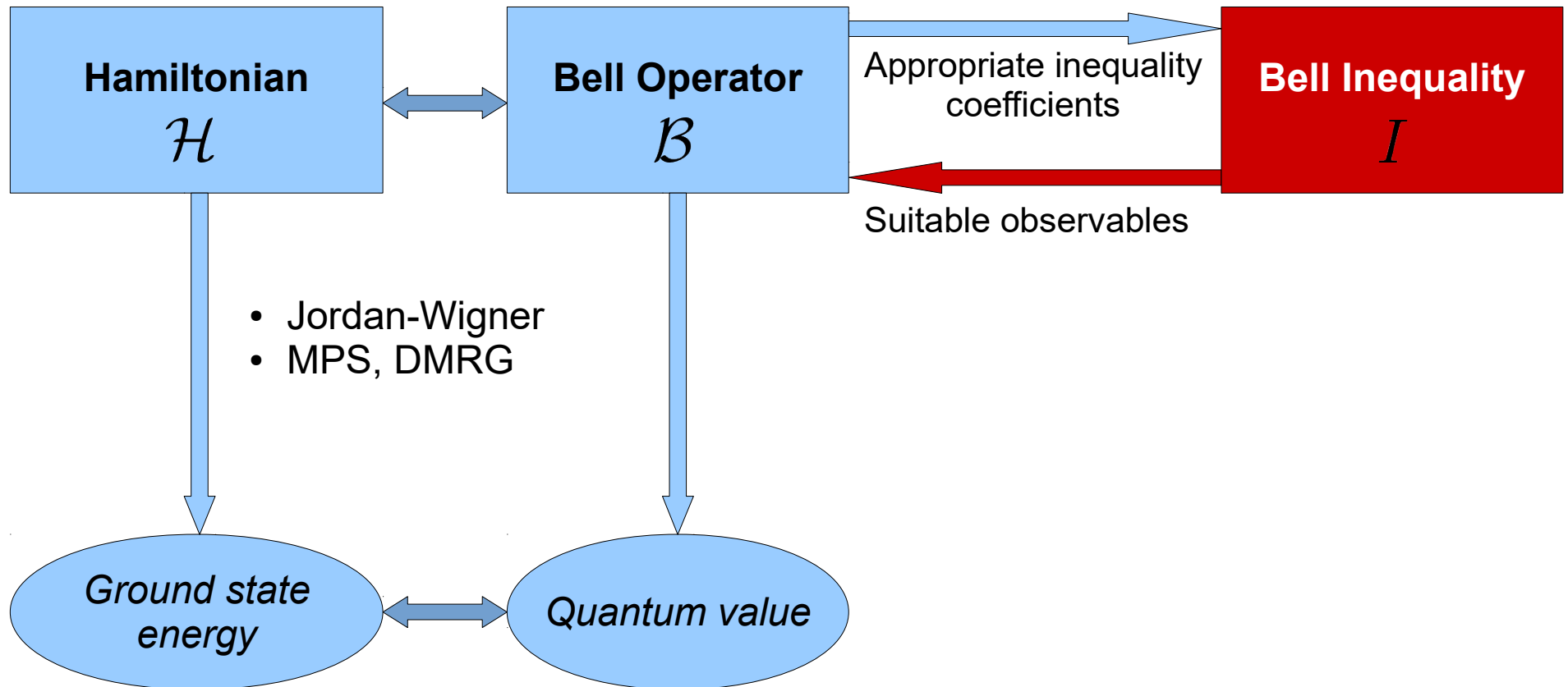
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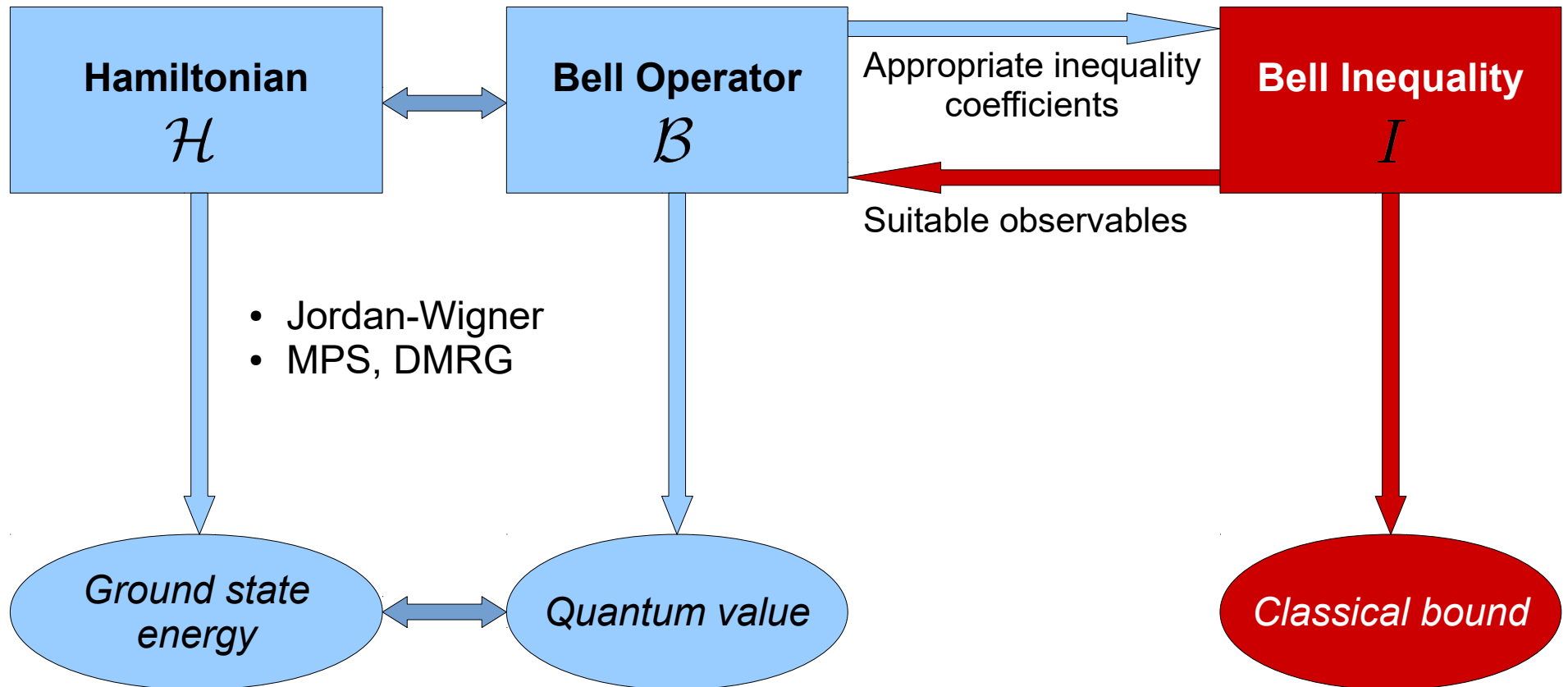
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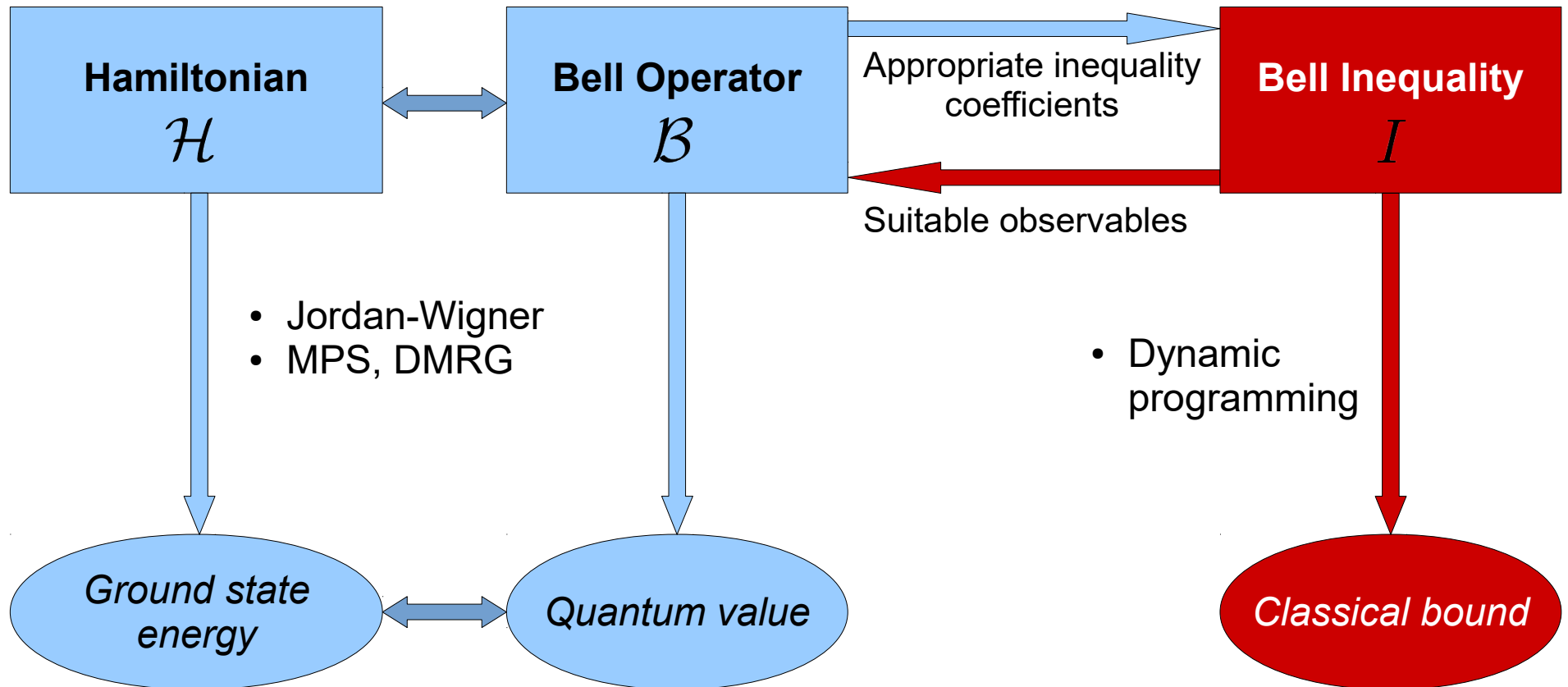
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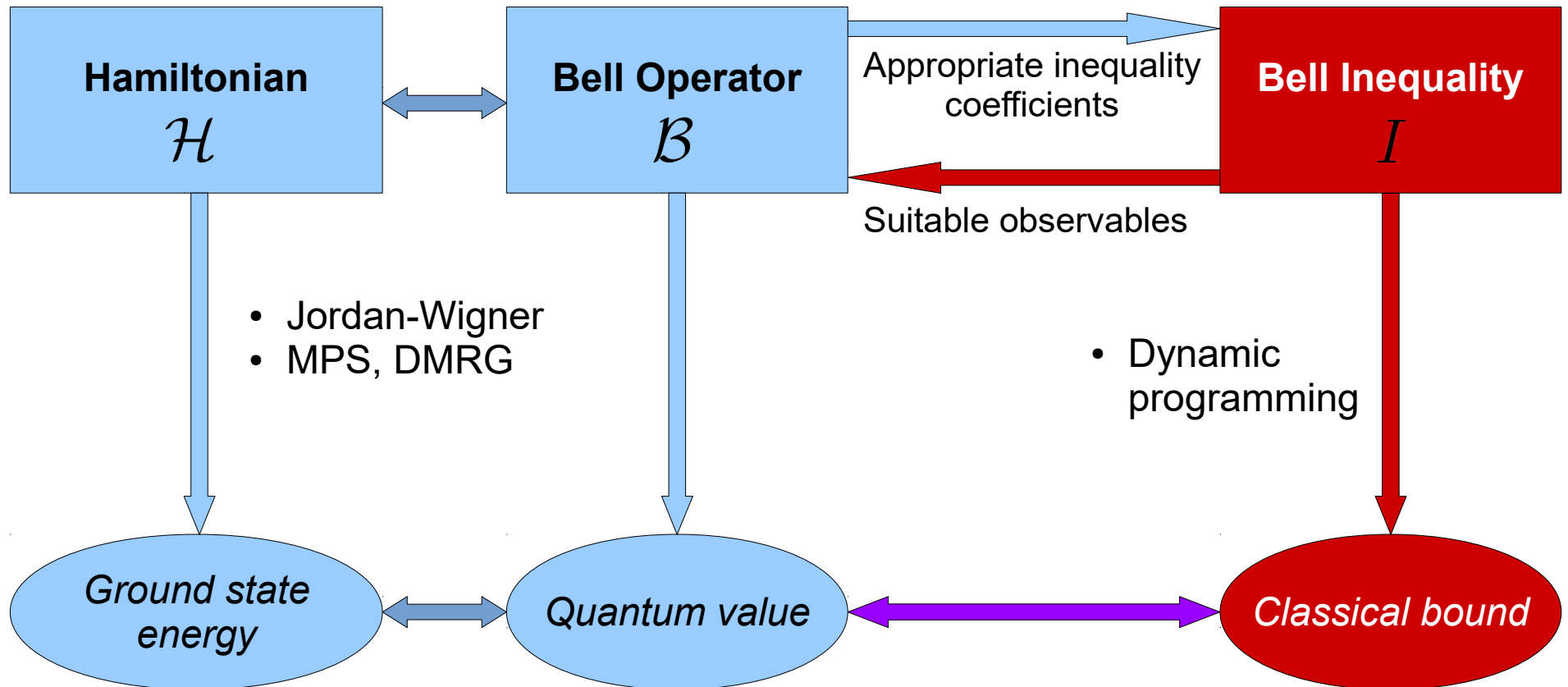


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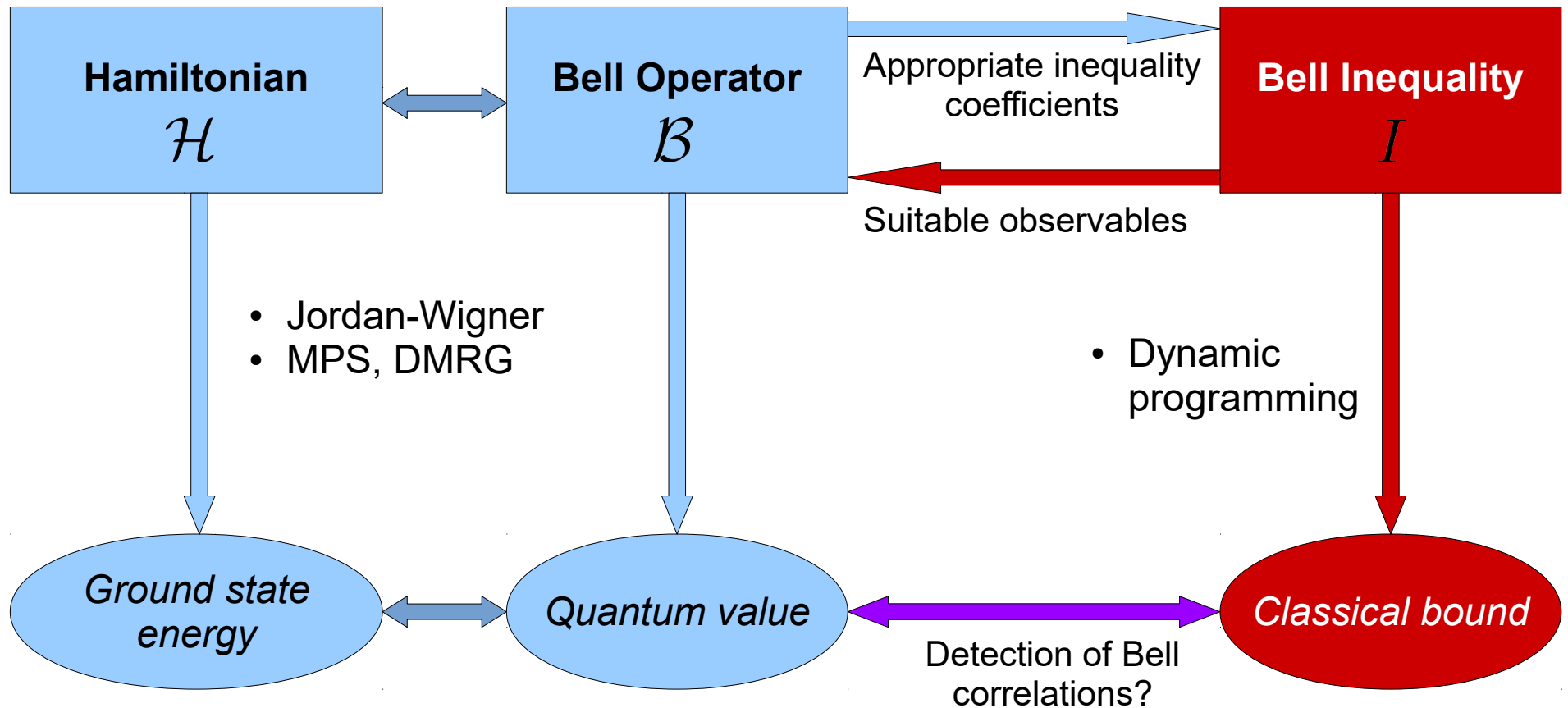




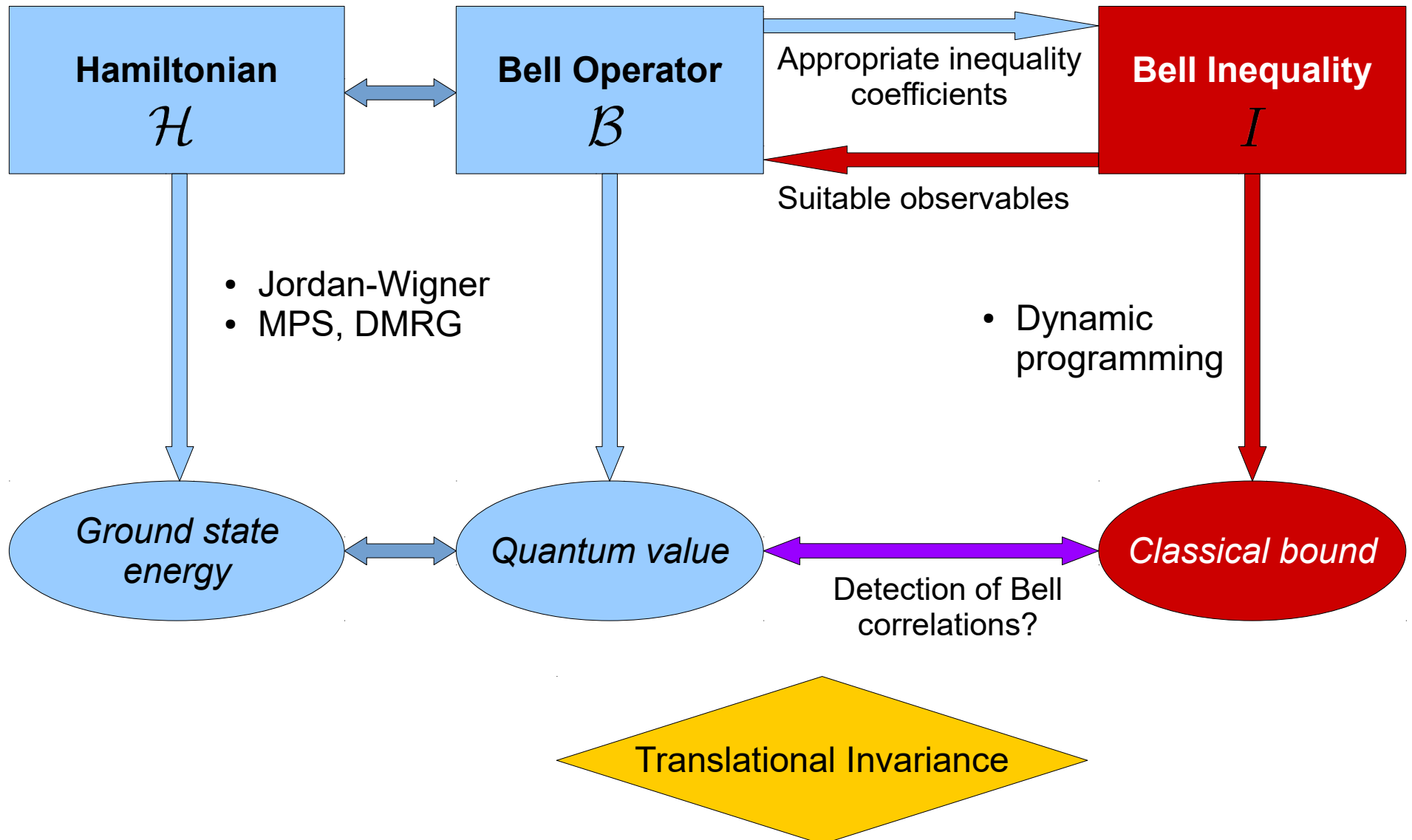
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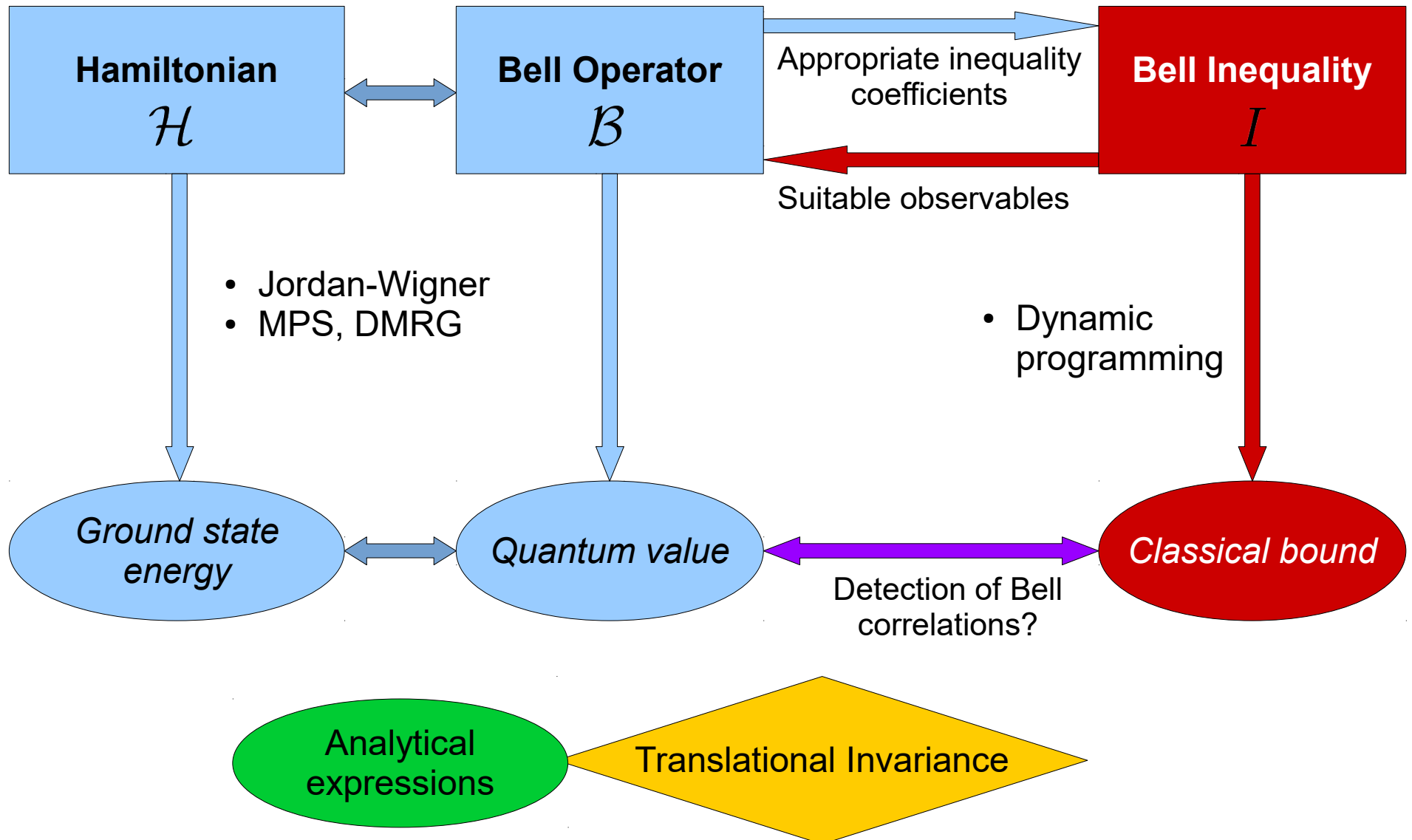
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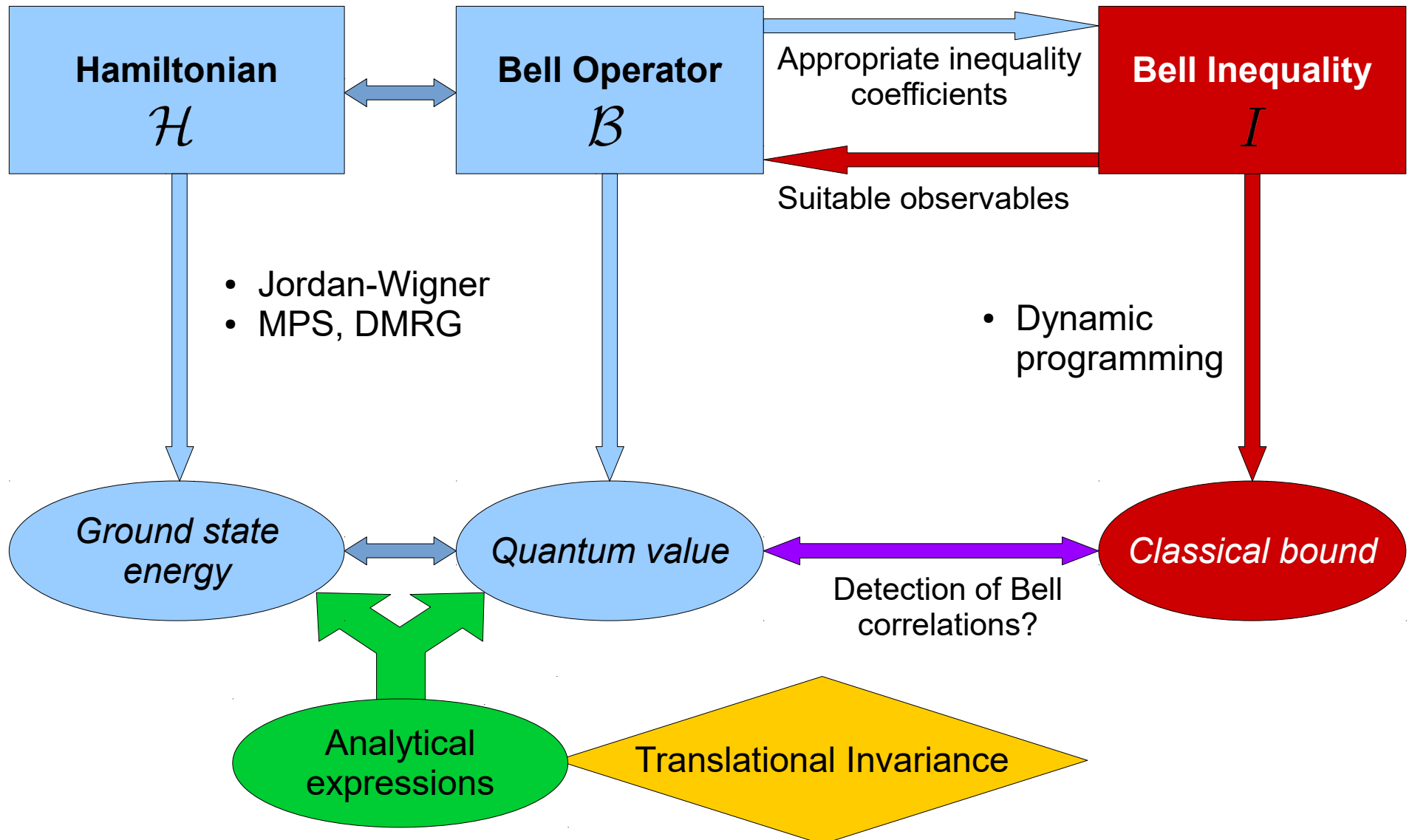
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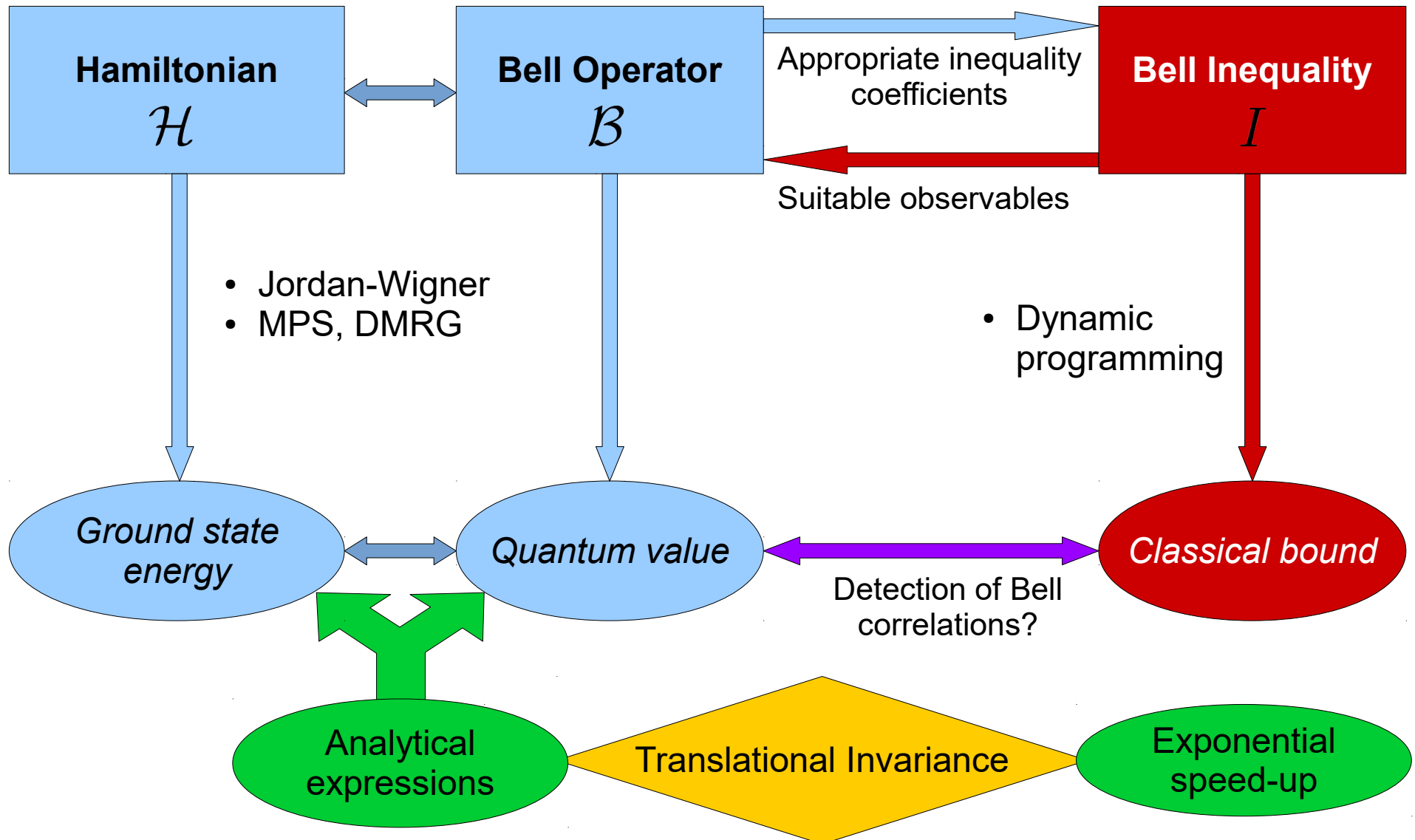
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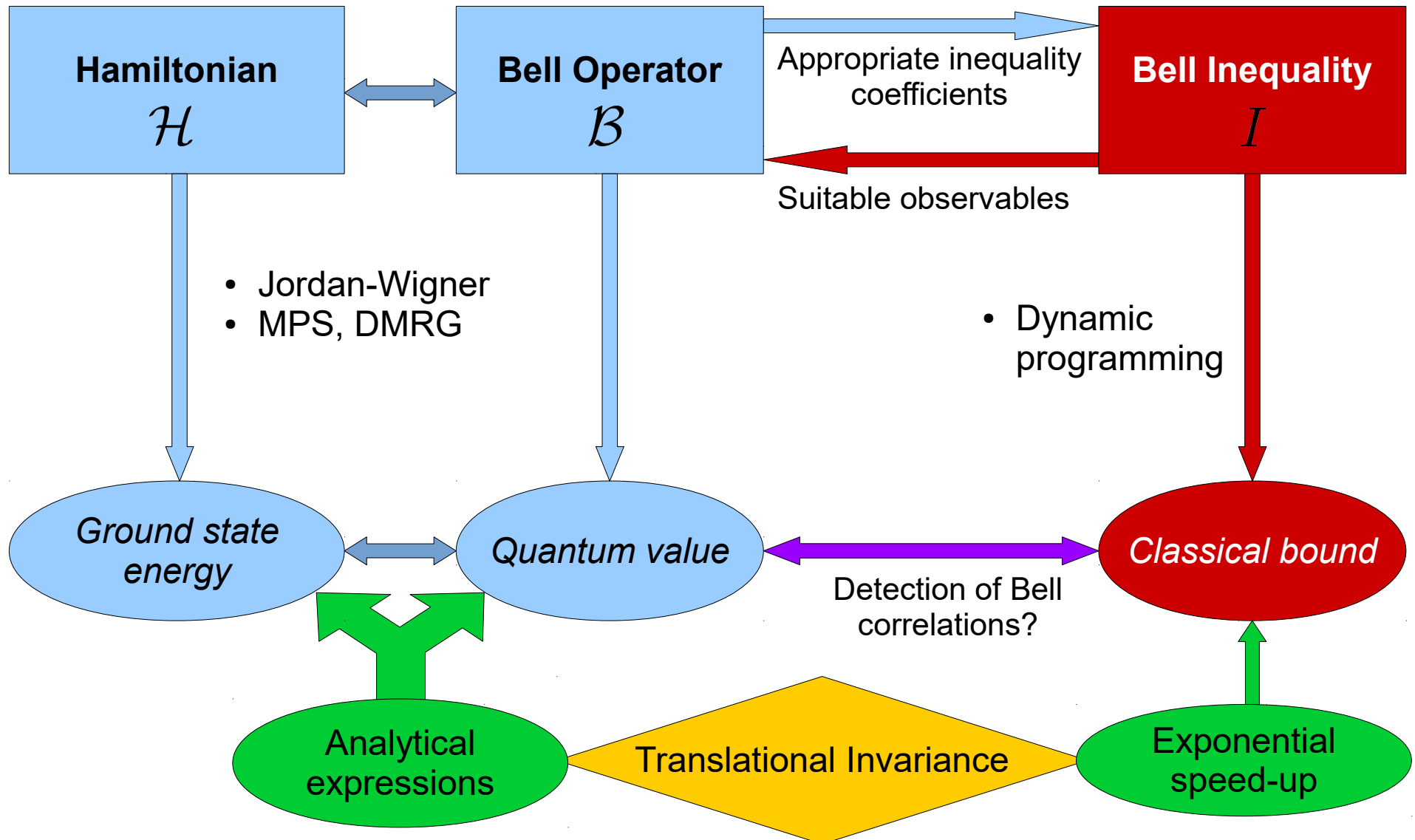
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# The setting





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- Spin – 1/2 Hamiltonians



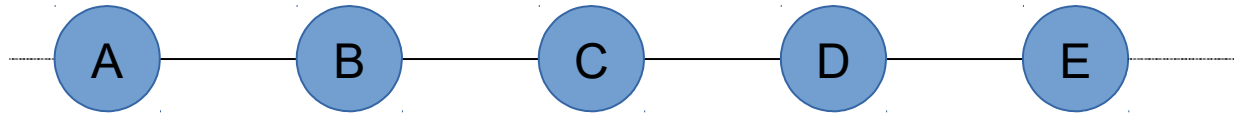
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- Spin – 1/2 Hamiltonians
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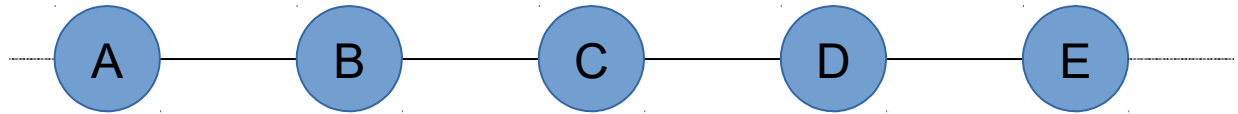
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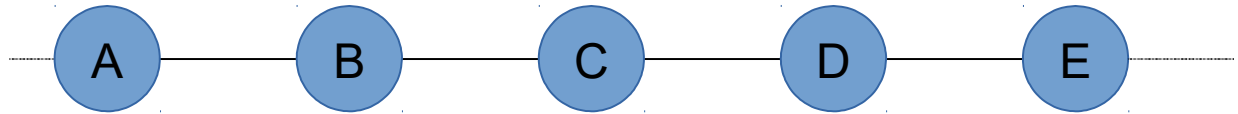
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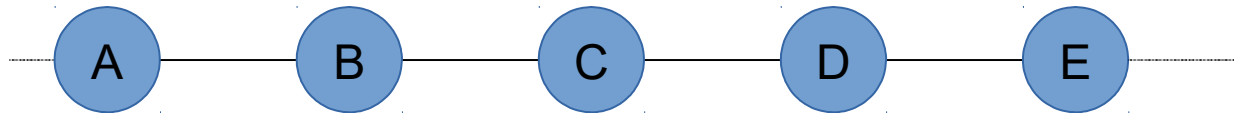
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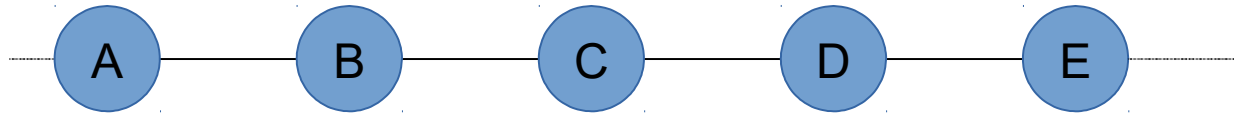
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$$\mathcal{H} = \sum_{i=0}^{n-1} \left( t^{(i)} \sigma_z^{(i)} + \sum_{r=1}^R \sum_{\alpha, \beta \in \{x, y\}} t_{\alpha, \beta}^{(i, r)} \text{Str}_{\alpha, \beta}^{(i, r)} \right)$$

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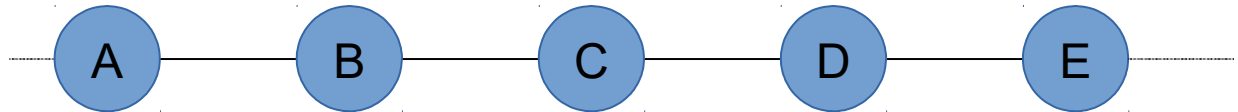


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- Every Hamiltonian of this form

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# Finding the ground state energy (I)

- Exact diagonalization

- Jordan – Wigner transformation: Spins to fermions

$$\hat{c}_{i,0} \leftrightarrow \prod_{j=0}^{i-1} \sigma_z^{(j)} \sigma_x^{(i)}, \quad \hat{c}_{i,1} \leftrightarrow - \prod_{j=0}^{i-1} \sigma_z^{(j)} \sigma_y^{(i)}$$

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becomes quadratic

$$\hat{\mathcal{H}} = \frac{\hat{\mathbb{1}}}{2} \sum_{i,j=0}^{n-1} \sum_{\alpha,\beta=0}^1 H_{i,\alpha;j,\beta} \hat{c}_{i,\alpha} \hat{c}_{j,\beta}$$



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Real, antisymmetric

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# Finding the ground state energy (II)





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- Williamson eigendecomposition



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$$H = OJO^T \quad O \in \mathcal{O}(2n) \quad J = \bigoplus_{k=0}^{n-1} \begin{pmatrix} 0 & \varepsilon_k \\ -\varepsilon_k & 0 \end{pmatrix}$$



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Mutually commuting



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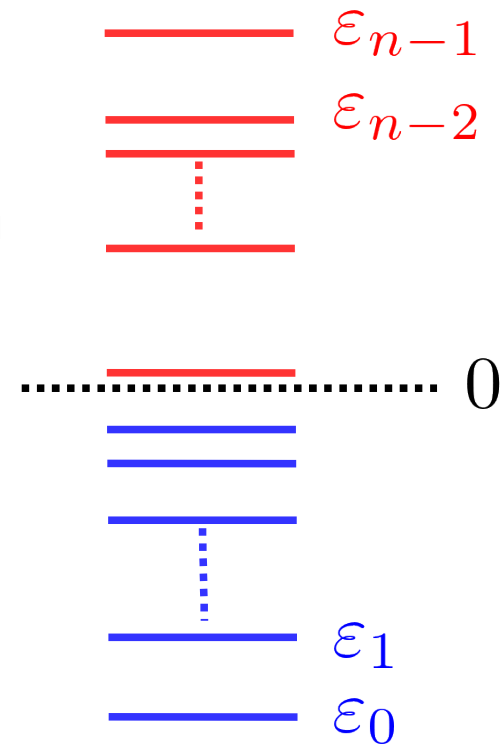
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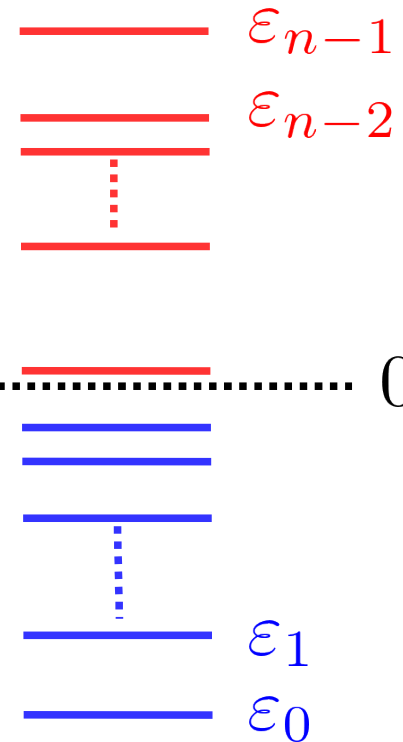
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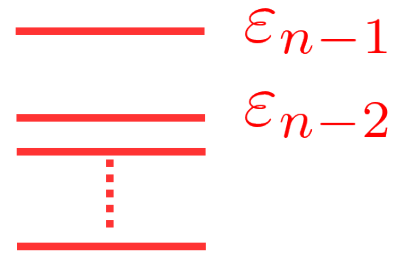
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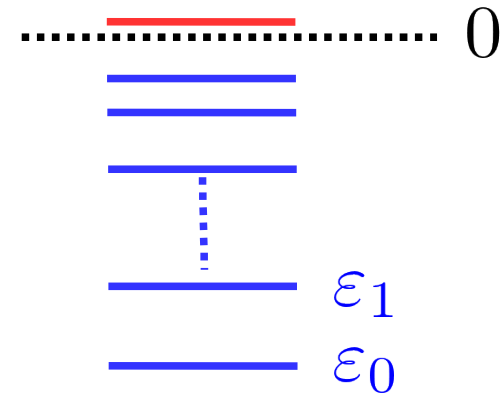
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Ground state energy

$$\beta_Q = \sum_{k=0}^{n-1} s_k \varepsilon_k$$





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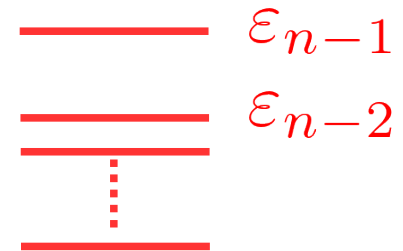
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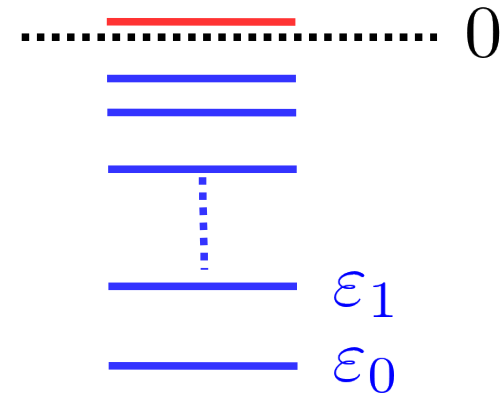


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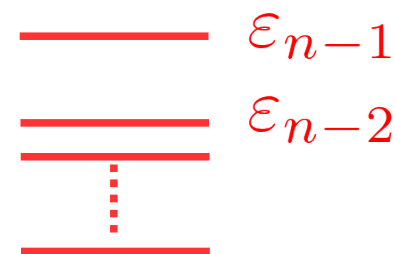
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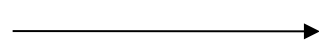
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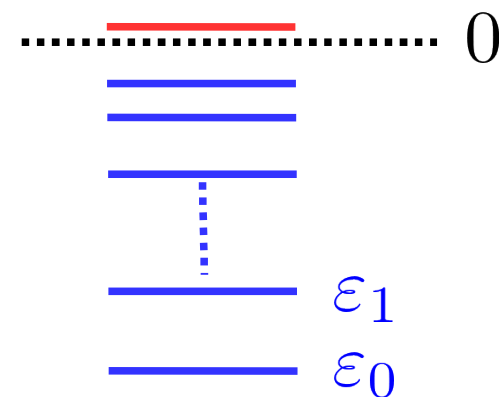
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The parity imposes a superselection rule



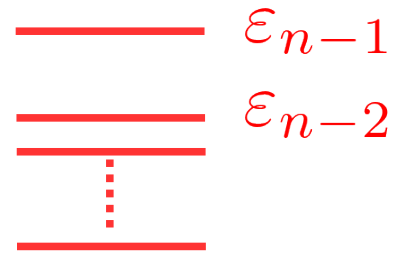
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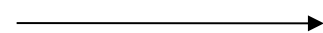
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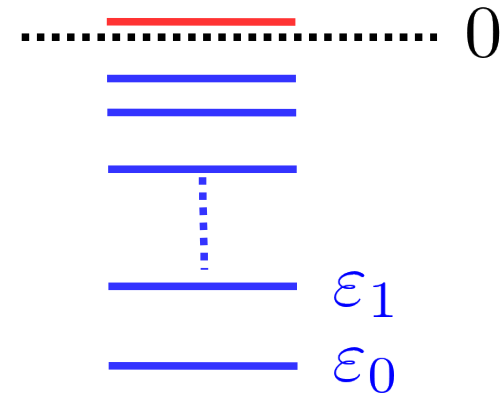
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$$p = (\det O) \prod_{k=0}^{n-1} s_k$$



# Assigning a Bell inequality



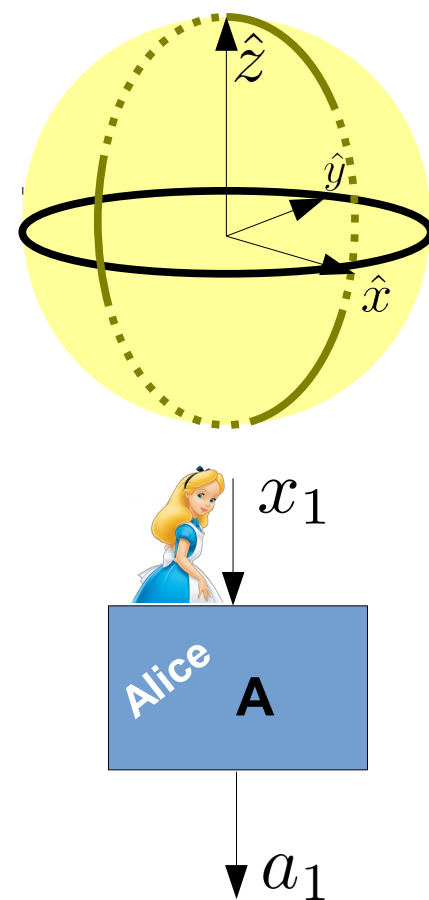
# Assigning a Bell inequality

- We want a Bell operator of the form  $\mathcal{B} = \beta_C \mathbb{1} + \mathcal{H}$



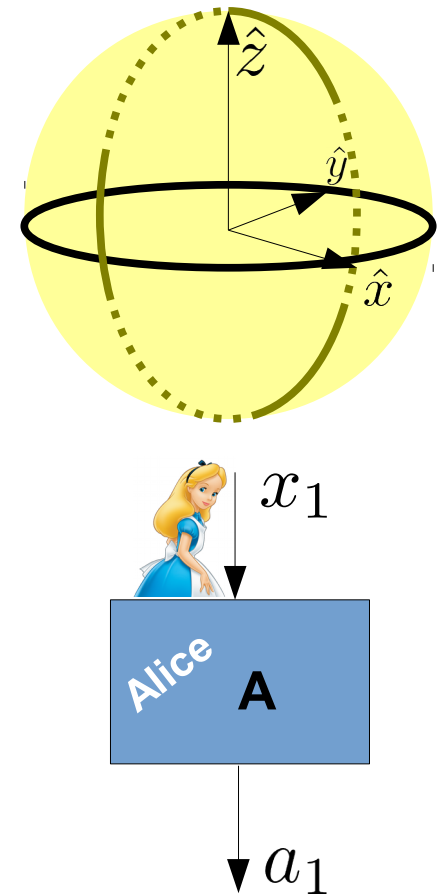
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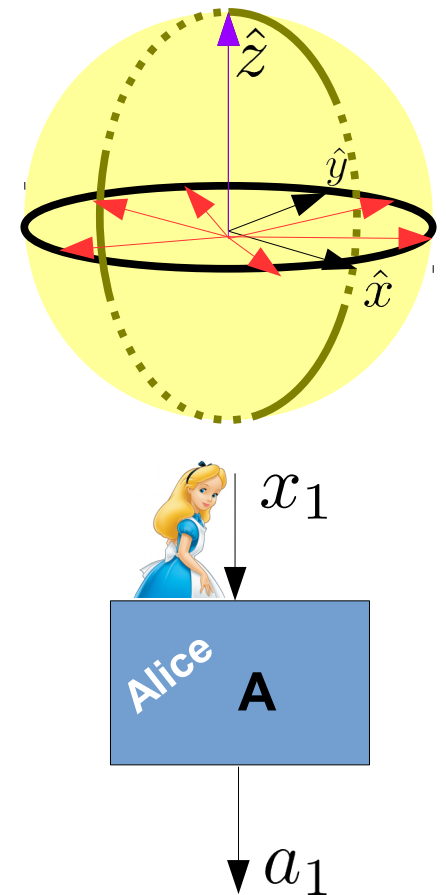
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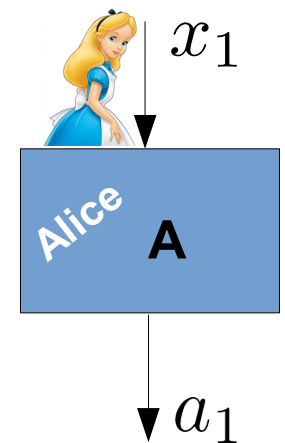
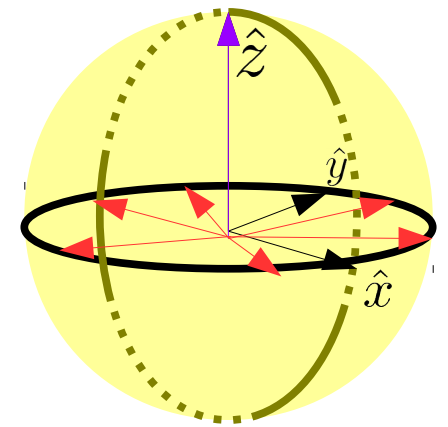




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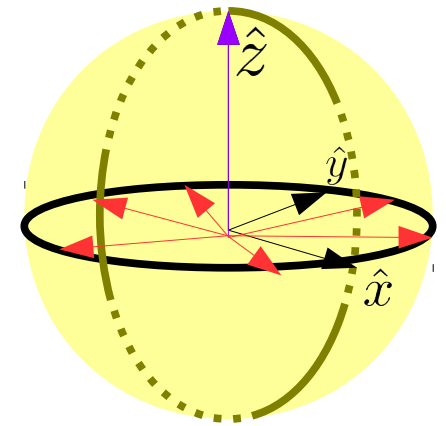
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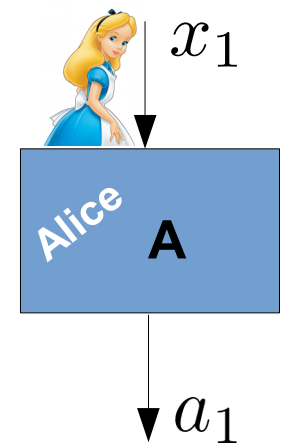
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$$I = \sum_{i=0}^{n-1} \left( \gamma^{(i)} M_m^{(i, 0)} + \sum_{r=1}^R \sum_{k, l=0}^{m-1} M_{(k, m, \dots, m, l)}^{(i, r)} \right)$$



# Finding the classical bound



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- We want a Bell inequality of the form  $I + \beta_C \geq 0$



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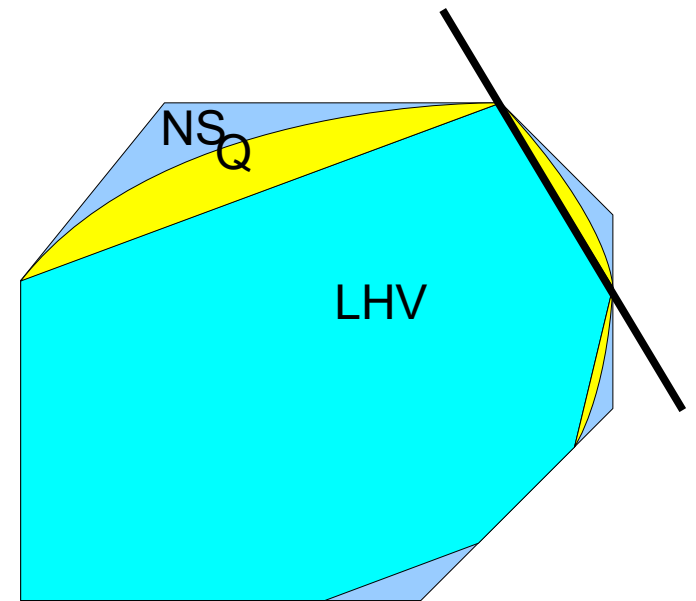
$$\beta_C = - \min_{\text{LHVM}} I$$



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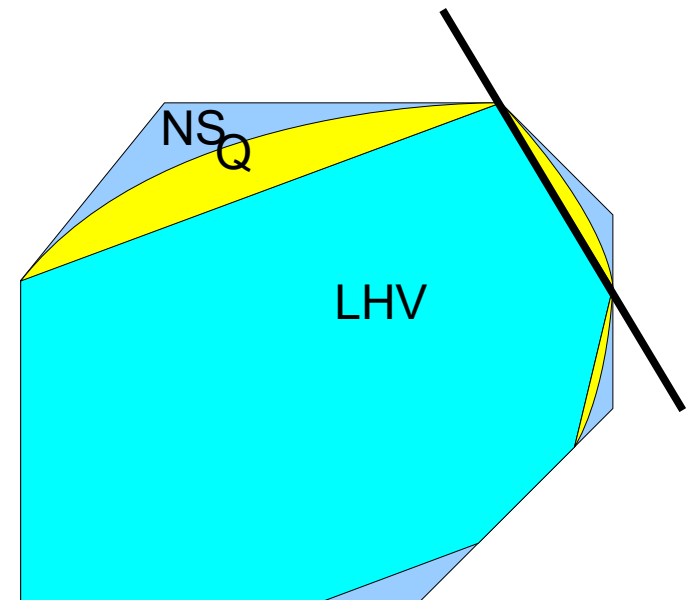
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- Fine's Theorem:

[A. Fine, Phys. Rev. Lett. **48**, 291 (1982)]



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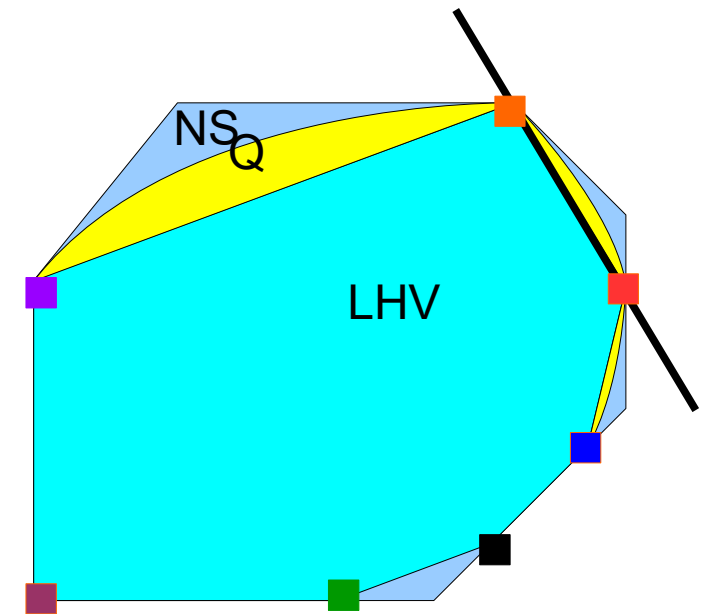
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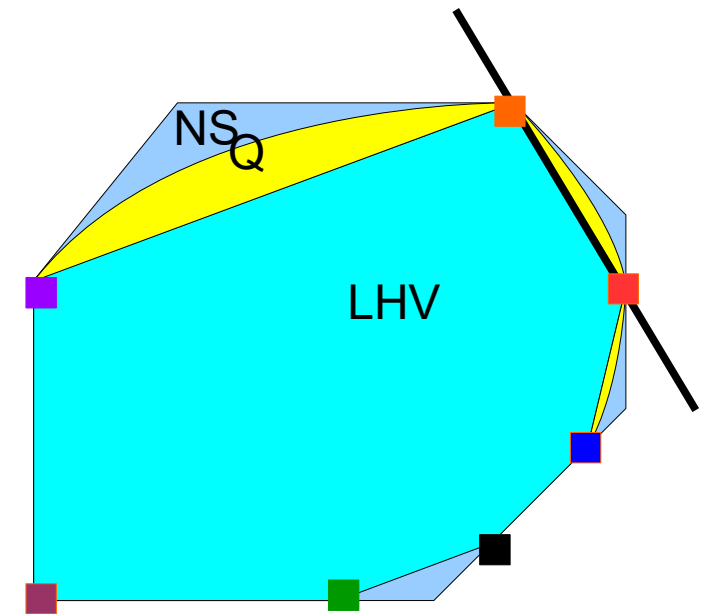
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$$M_0^{(i)}$$

$$M_1^{(i)}$$

$$M_2^{(i)}$$



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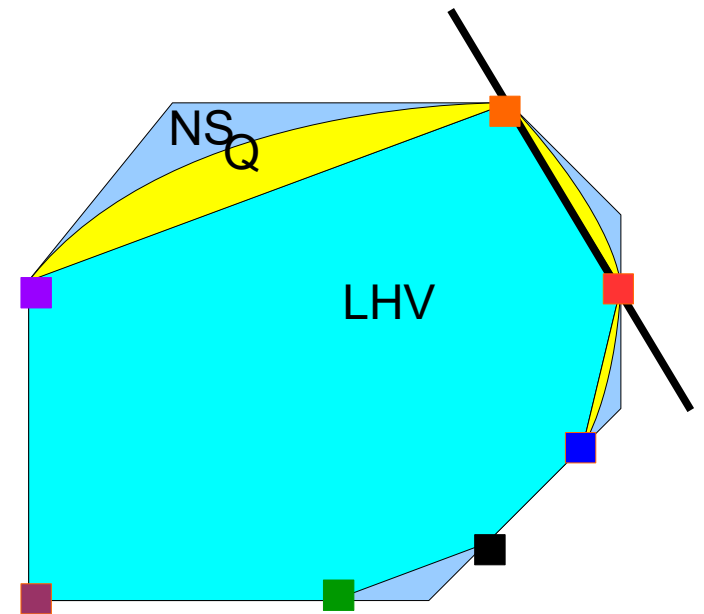
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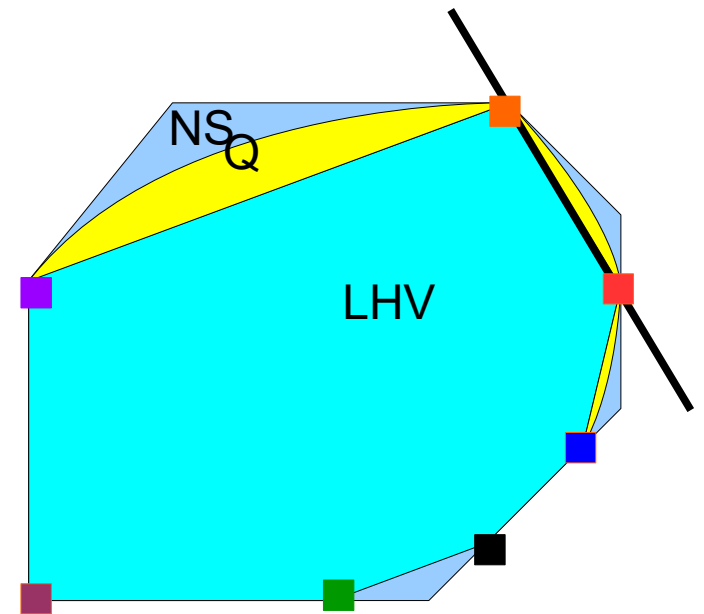
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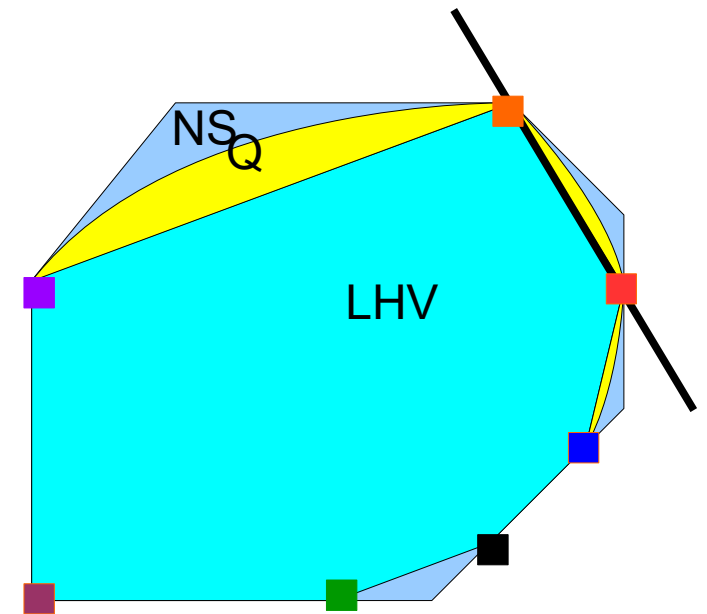
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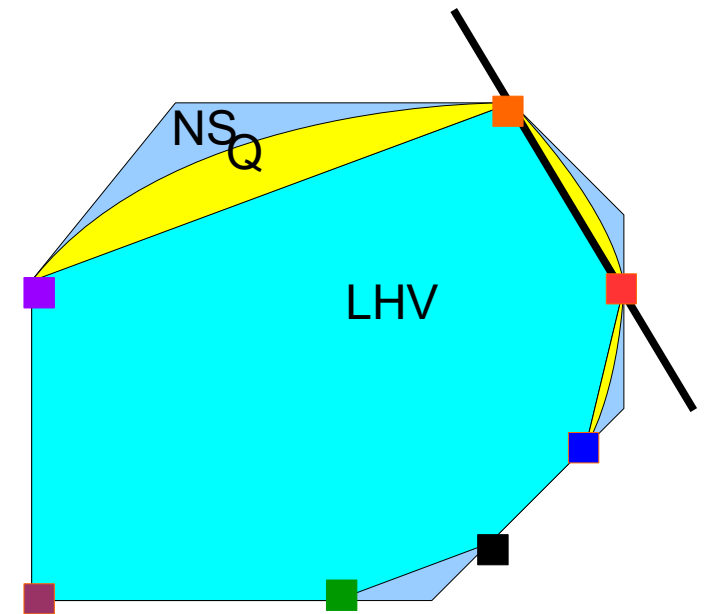
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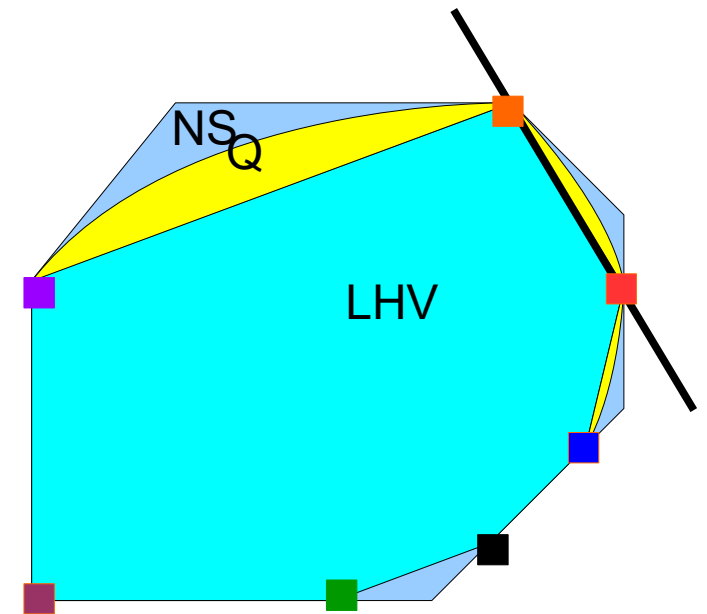
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$M_0^{(i)}$	●	●	●	●	...
$M_1^{(i)}$	●	●	●	●	...
$M_2^{(i)}$	●	●	●	●	...

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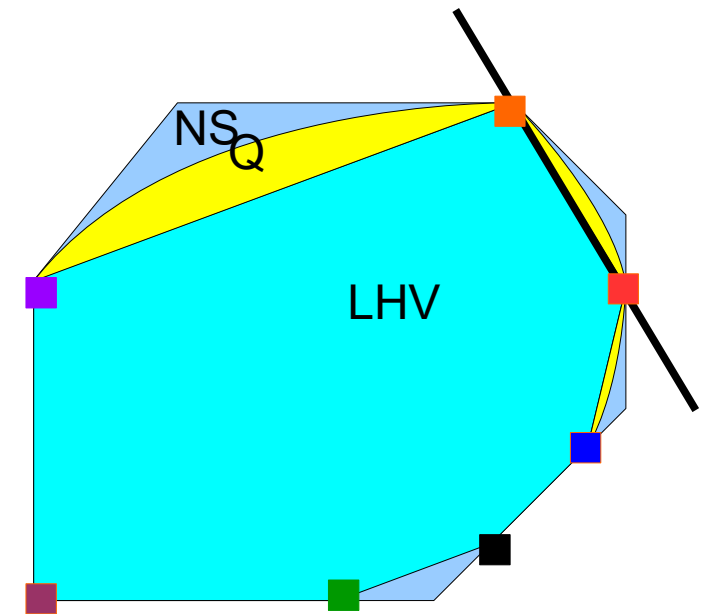
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Problem

					...
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$M_1^{(i)}$	●	●	●	●	...
$M_2^{(i)}$	●	●	●	●	...



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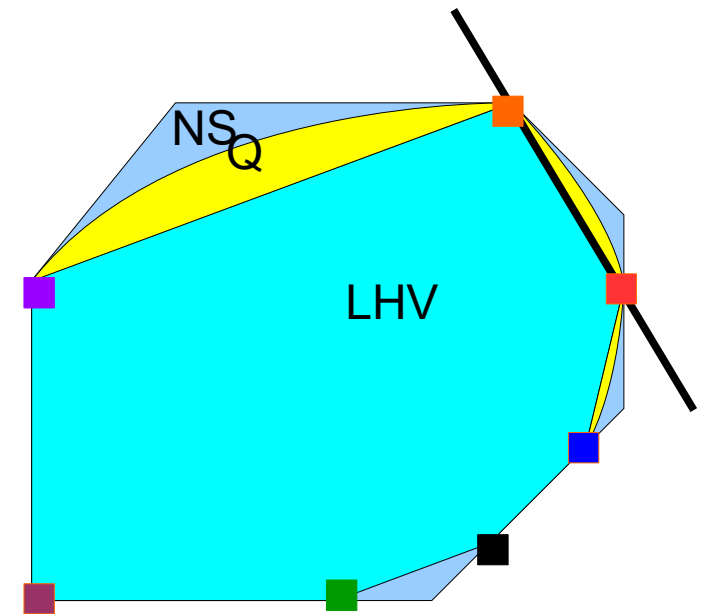
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Problem  
 $2^{mn}$  vertices

					...
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$M_1^{(i)}$	●	●	●	●	...
$M_2^{(i)}$	●	●	●	●	...



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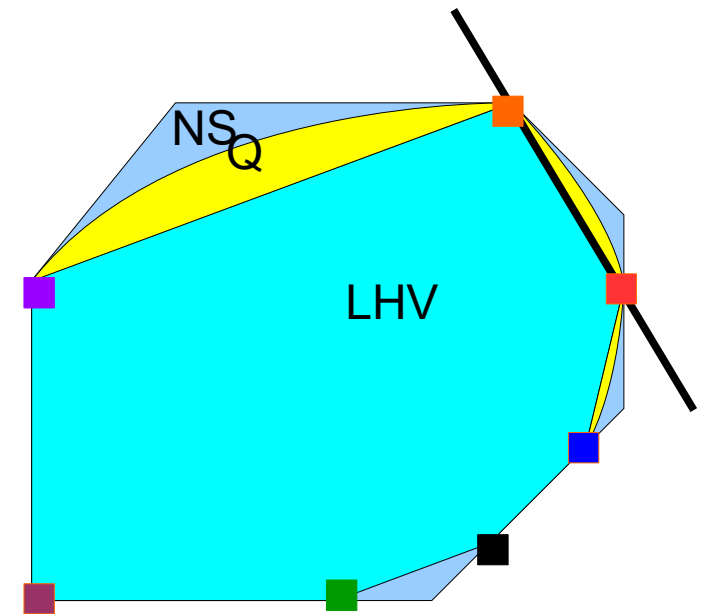
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Problem

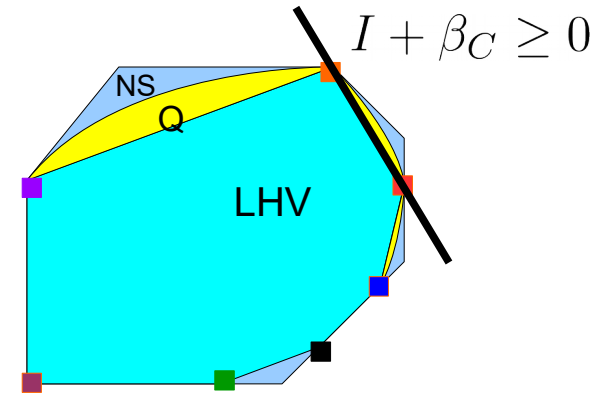
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For our Bell inequalities

Dynamic programming

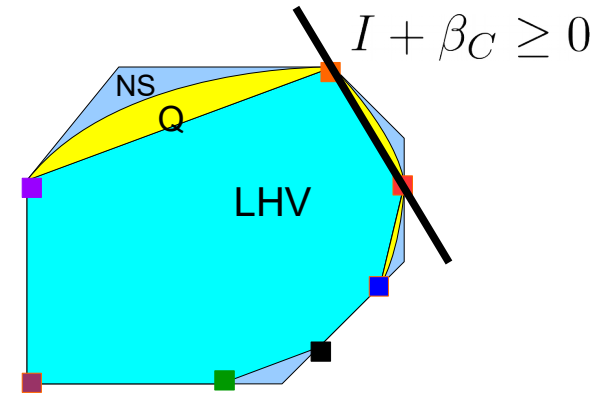


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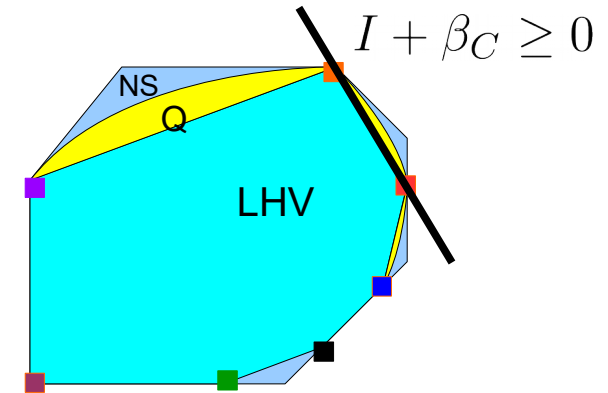
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- Optimization over all LHV models



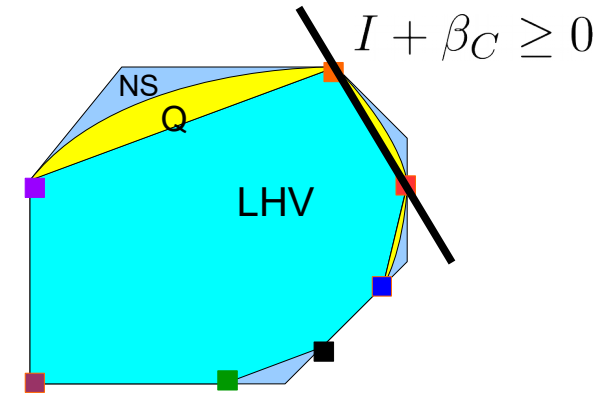
# Finding the classical bound

- Optimization over all LHV models
  - Linear programming (general case)



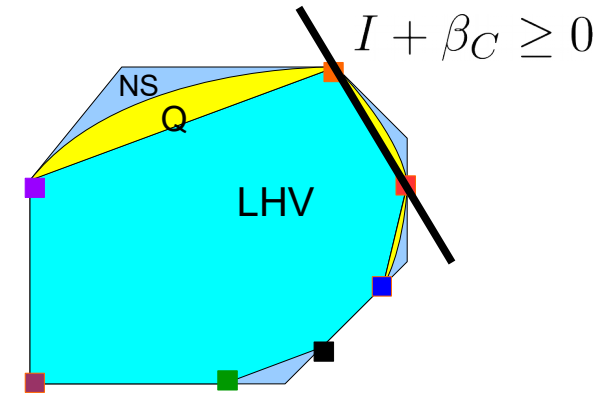
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# Finding the classical bound

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- Dynamic programming is extremely efficient for 1D-like BI



[N. Schuch, J. I. Cirac, Phys. Rev. A. **82**, 012314 (2010)]

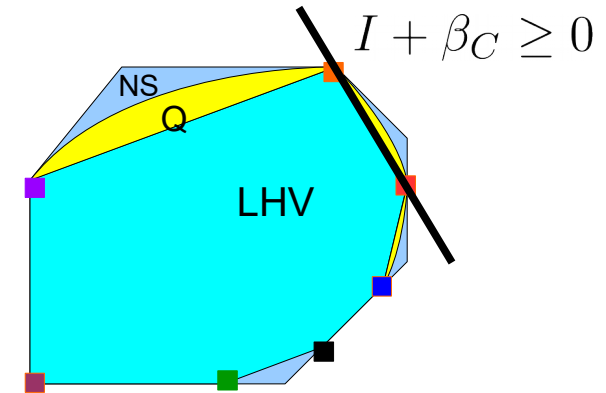


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[N. Schuch, J. I. Cirac, Phys. Rev. A. **82**, 012314 (2010)]

*Ingredients*



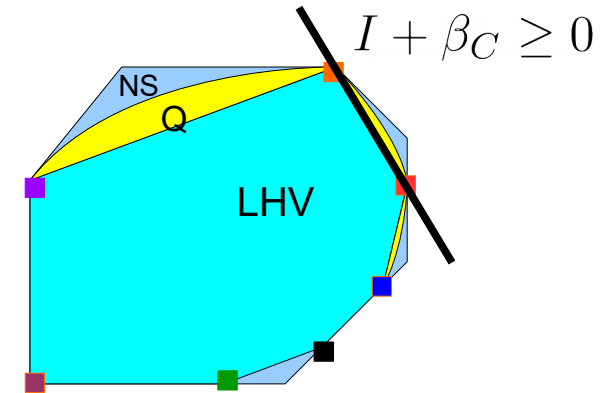
# Finding the classical bound

- Optimization over all LHV models
  - Linear programming (general case)
  - Impossible for many-body BI
- Dynamic programming is extremely efficient for 1D-like BI

[N. Schuch, J. I. Cirac, Phys. Rev. A. **82**, 012314 (2010)]

## *Ingredients*

- Recurrence relation





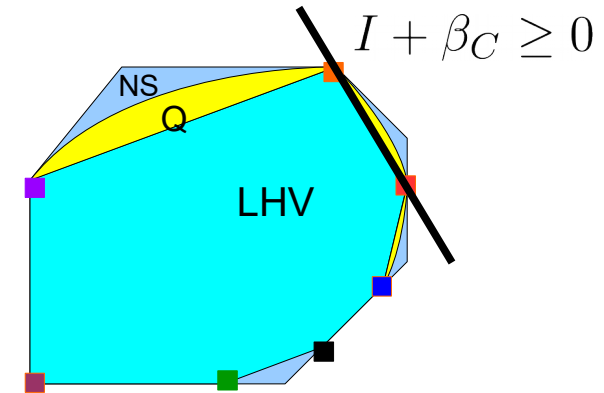
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- Compute & store intermediate sub-solutions



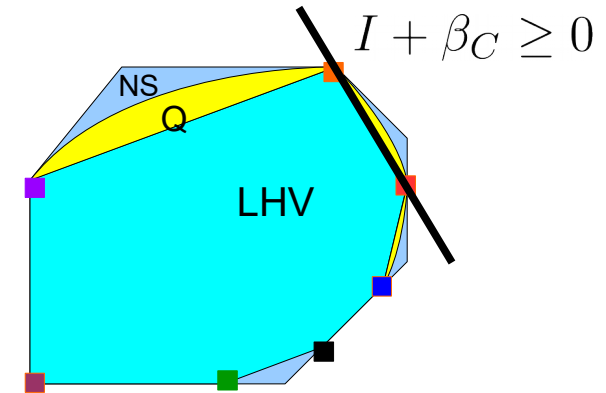
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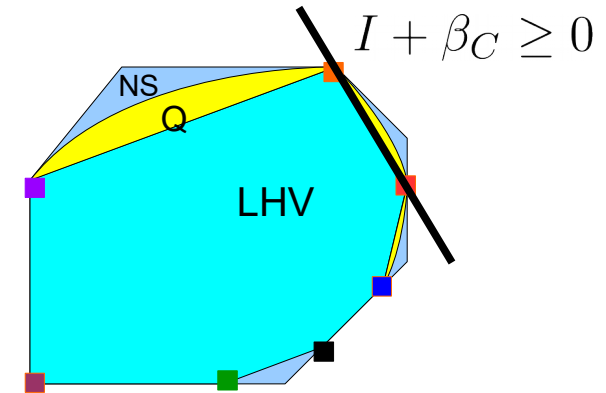
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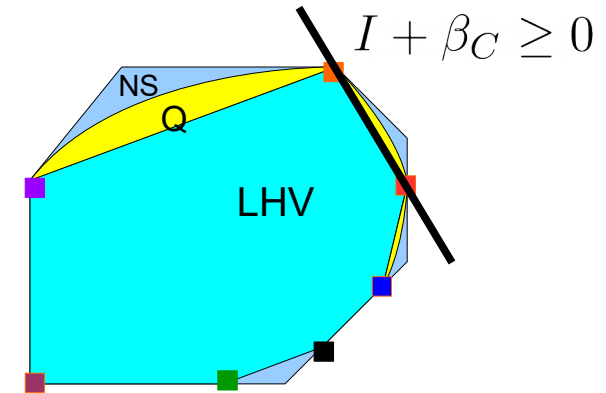
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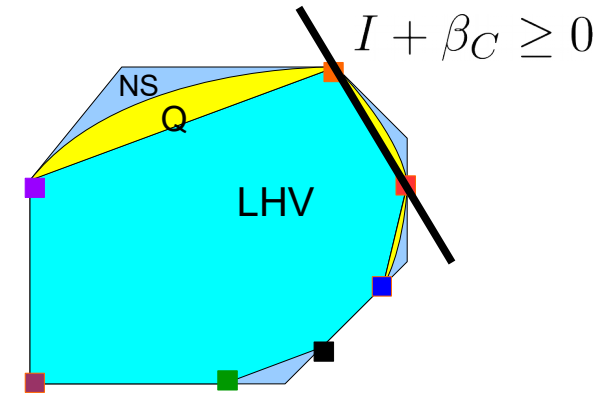
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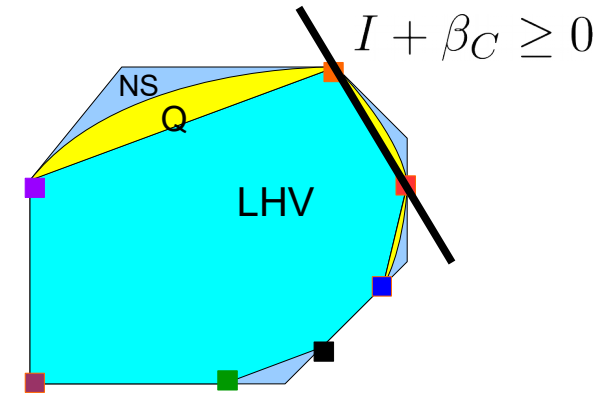
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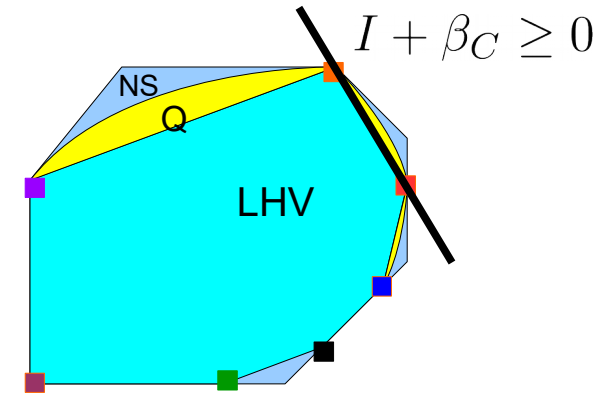
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- Polynomial scaling
- Constructive method of 1 optimal solution



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[N. Schuch, J. I. Cirac, Phys. Rev. A. **82**, 012314 (2010)]

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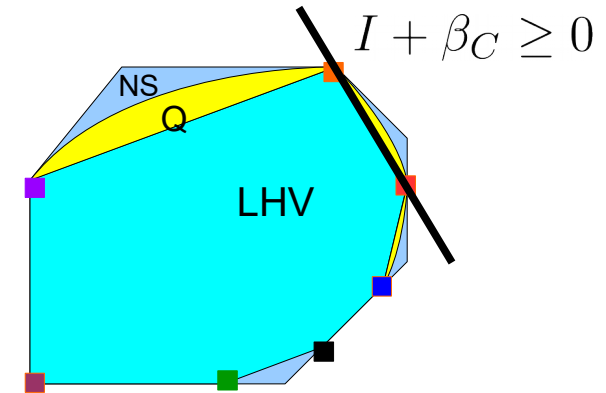
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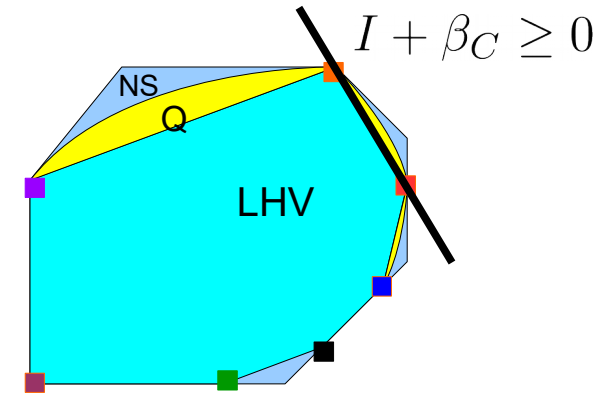
# of problems by category at [[TOPCODER](#)]<sup>®</sup>





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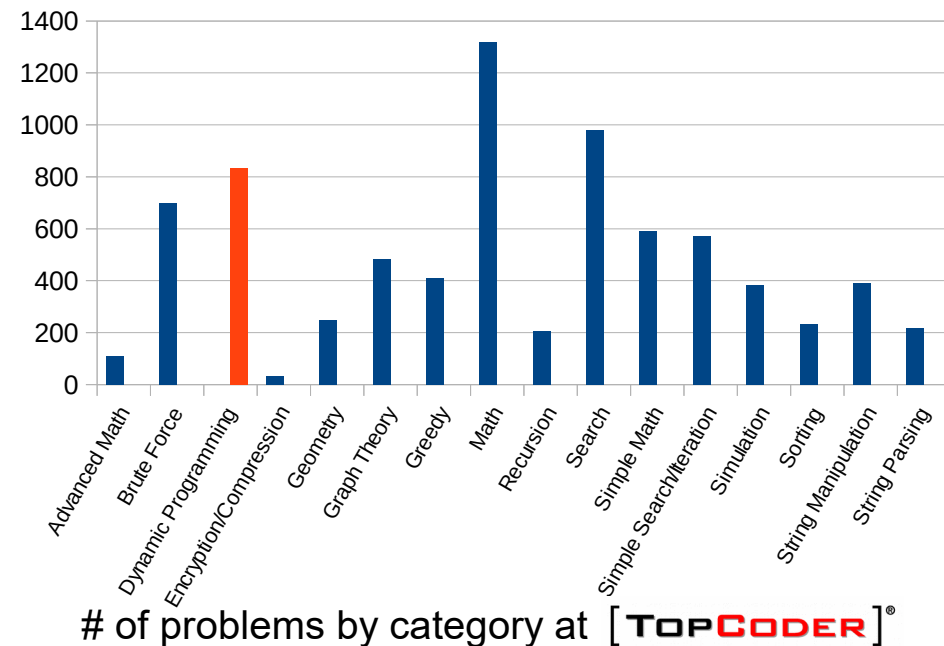
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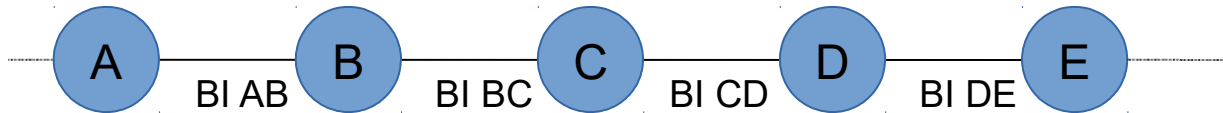


# Dynamic programming



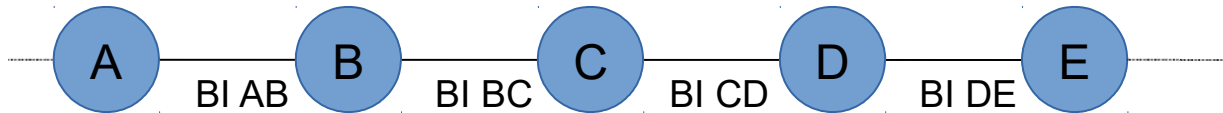
# Dynamic programming

- The Bell Inequality as a sum of smaller BI



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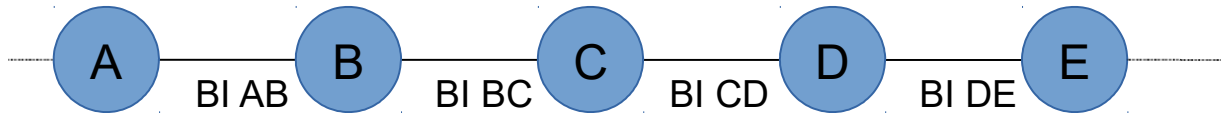


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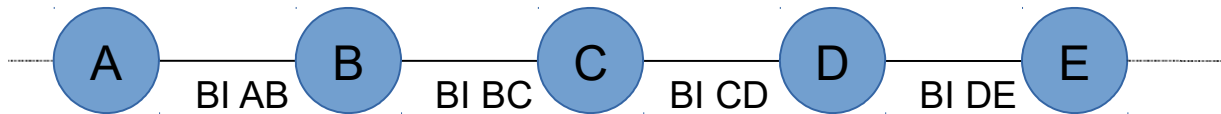


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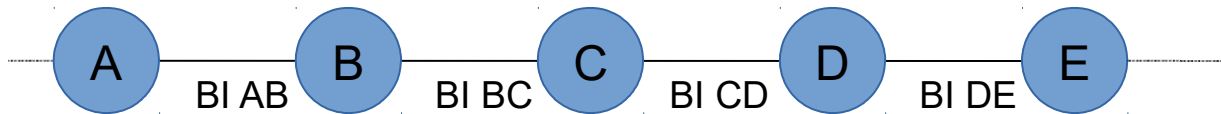


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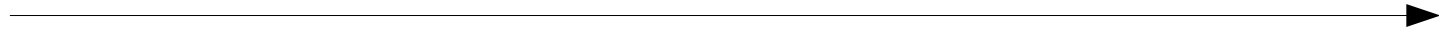
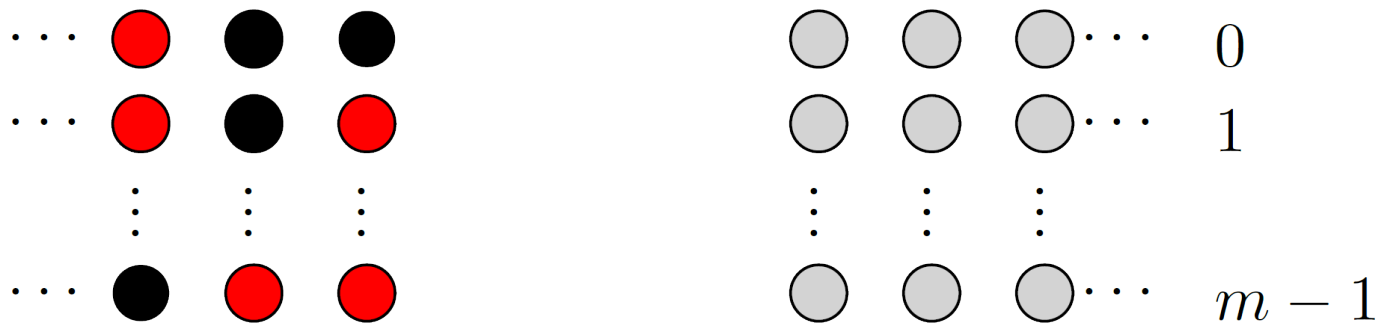


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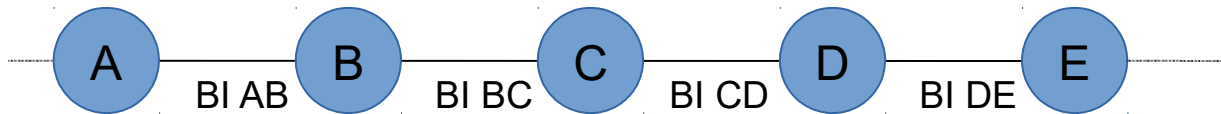


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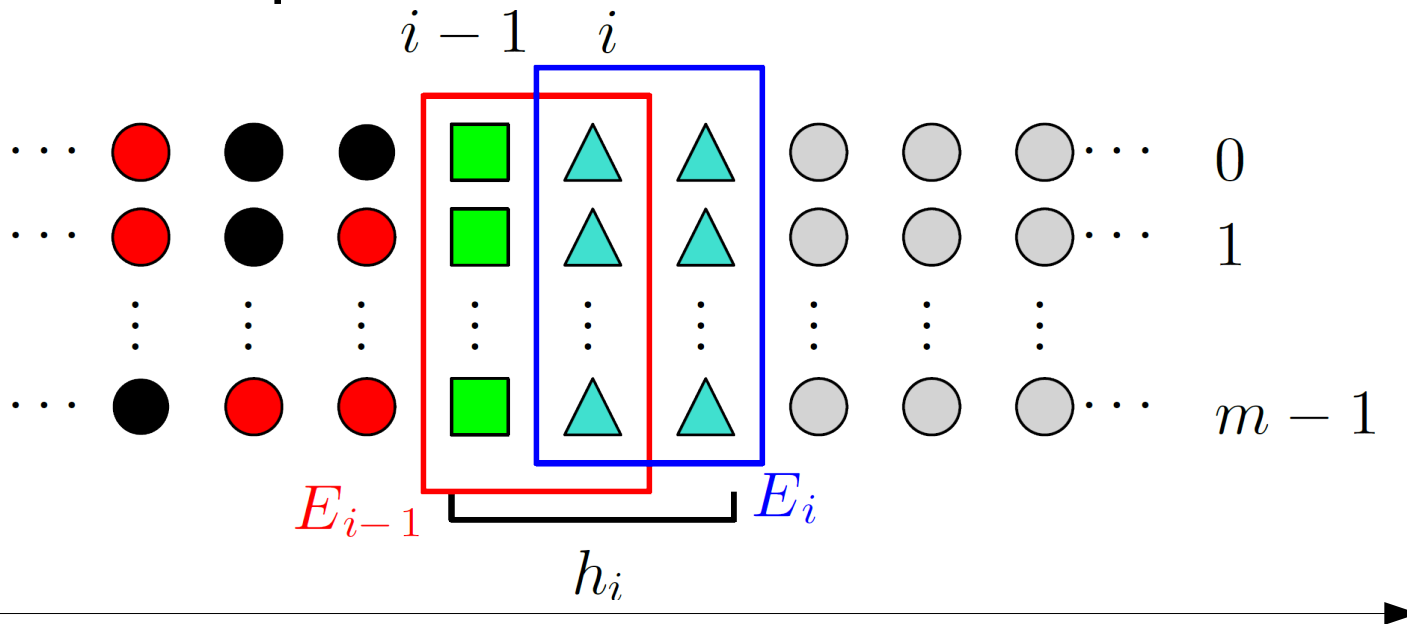


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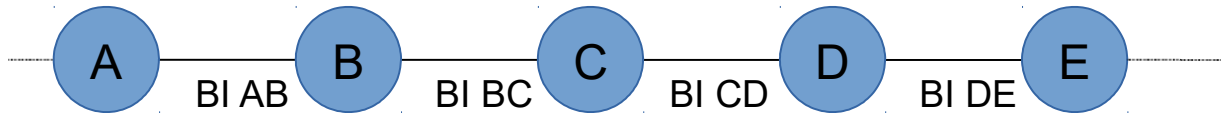
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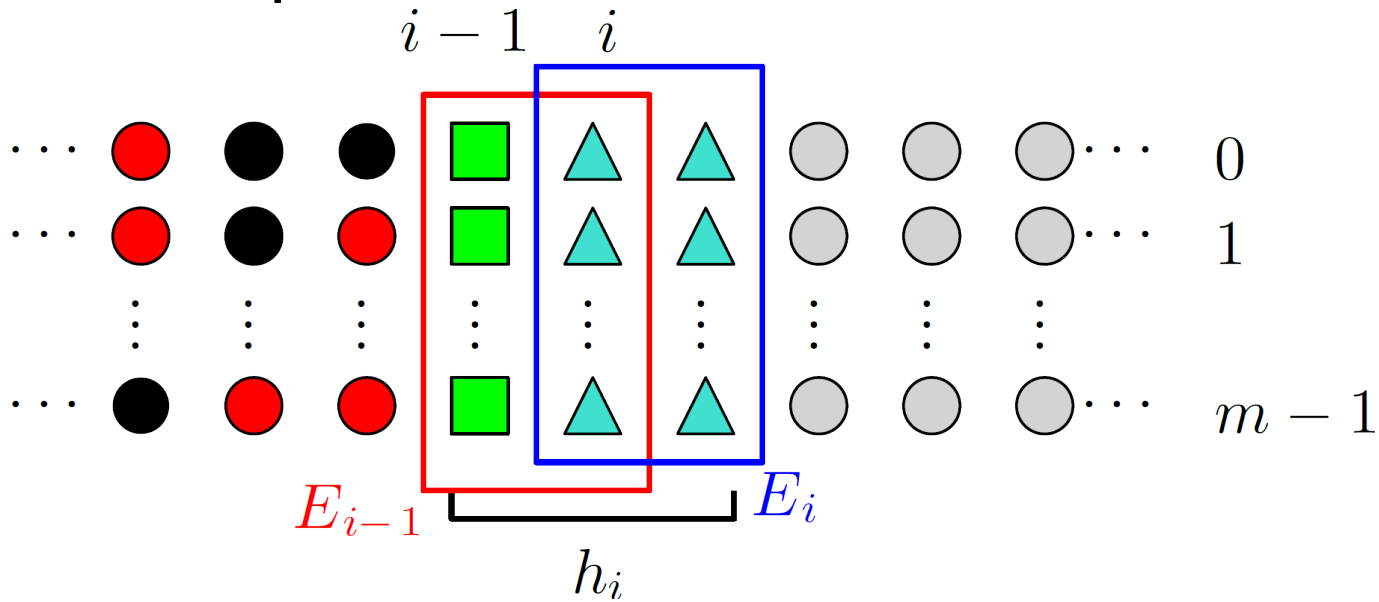


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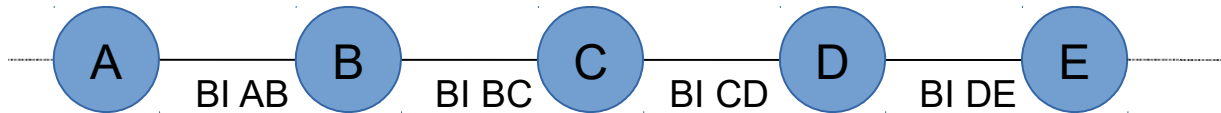
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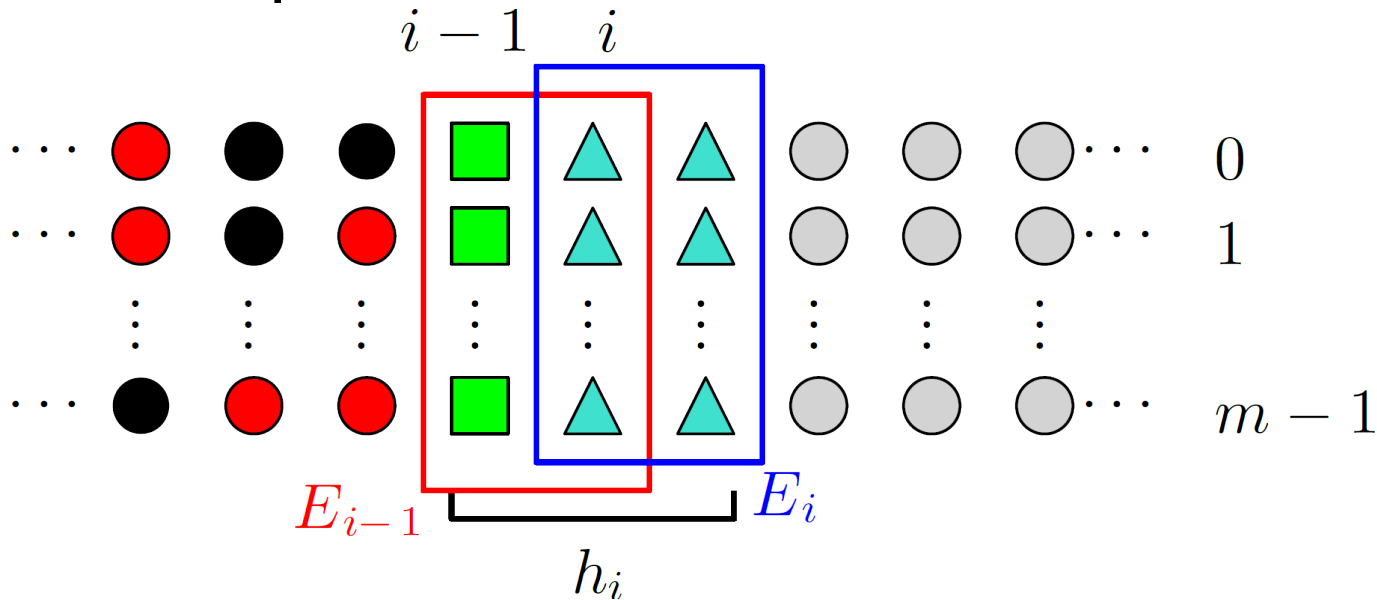
$$E_i(\triangle, \triangle) = \min_{\square} E_{i-1}(\square, \triangle) + h_i(\square, \triangle, \triangle)$$

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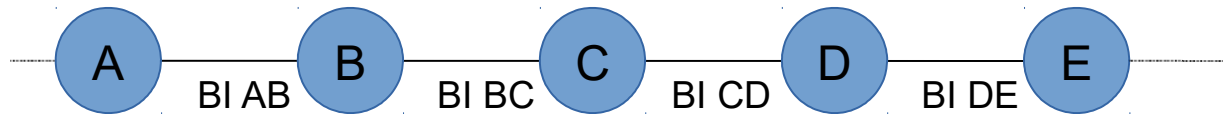
Classical bound at

$$\beta_C := E_n$$

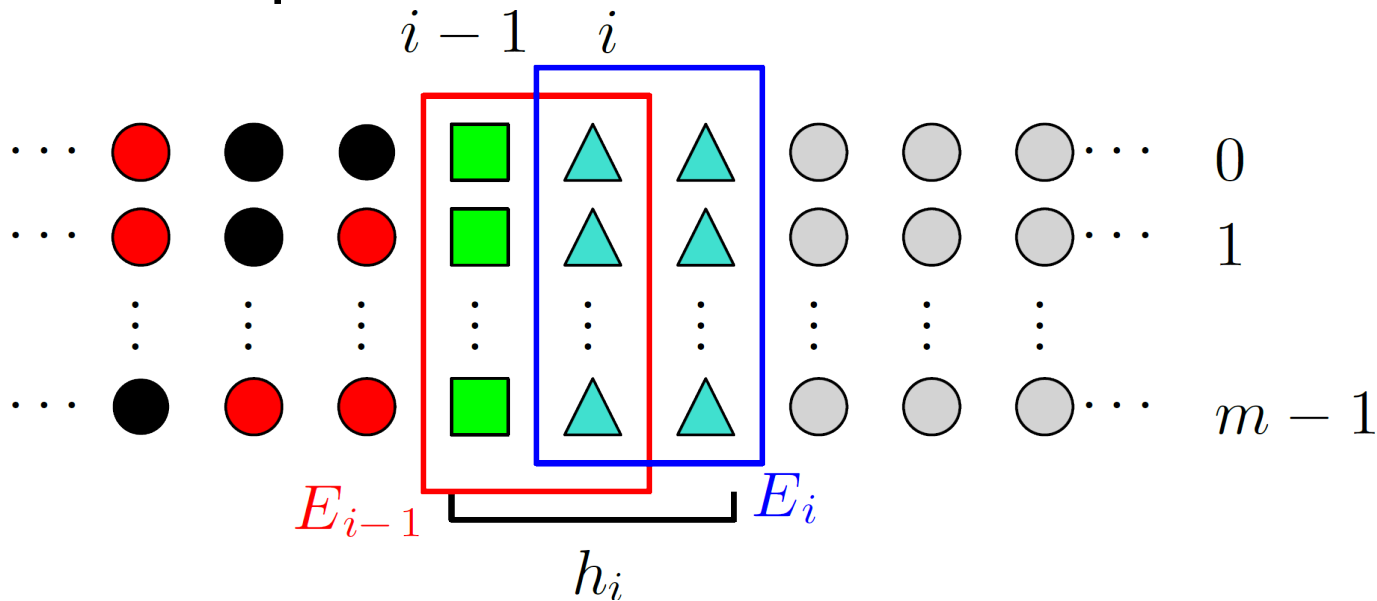
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Classical bound at

$$\beta_C := E_n$$

Overall complexity  $O(n)$

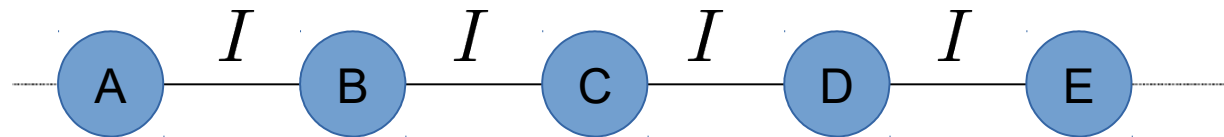
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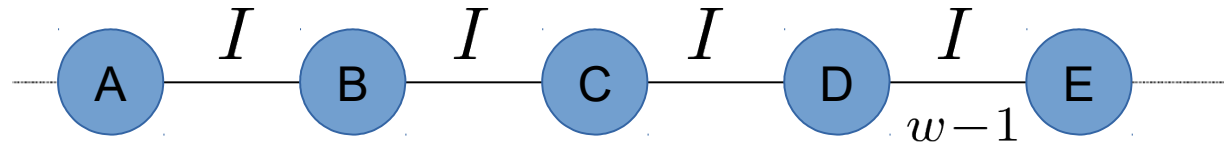
# Translationally invariant BI



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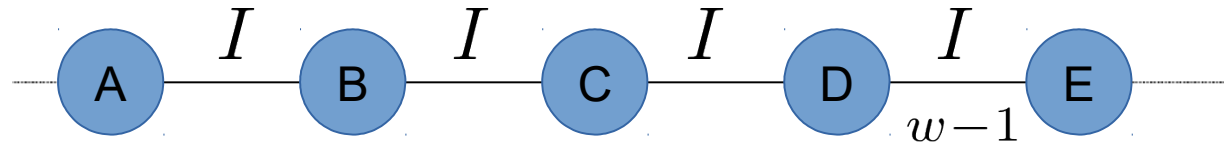
# Translationally invariant BI



- **Idea:** Minimize a function  $F = \min_{x_0, \dots, x_w} \sum_{j=0}^{w-1} f^{(0)}(x_j, x_{j+1})$



# Translationally invariant BI



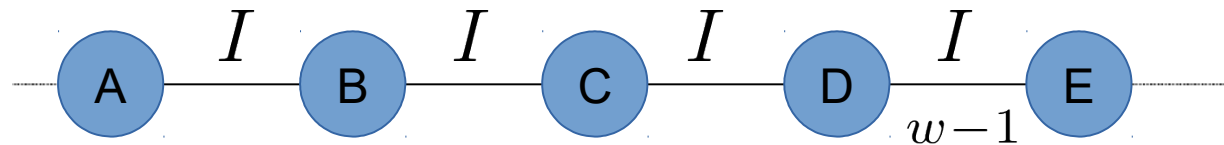
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by eliminating half of the variables at each step

$$f^{(t+1)}(x, y) = \min_z (f^{(t)}(x, z) + f^{(t)}(z, y))$$



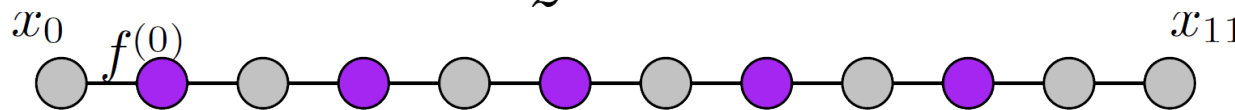
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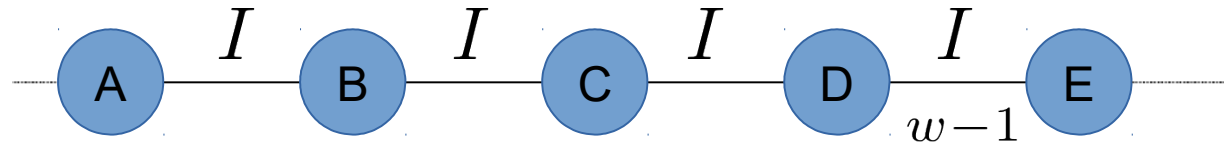
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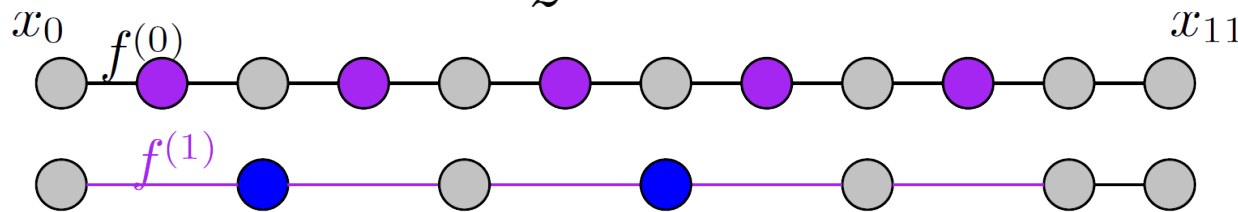
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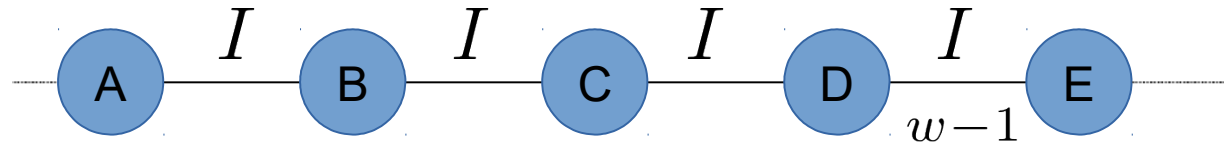
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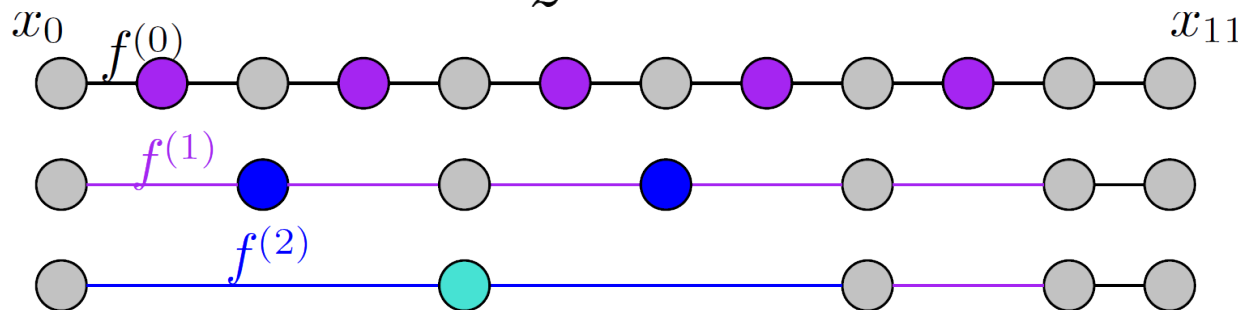
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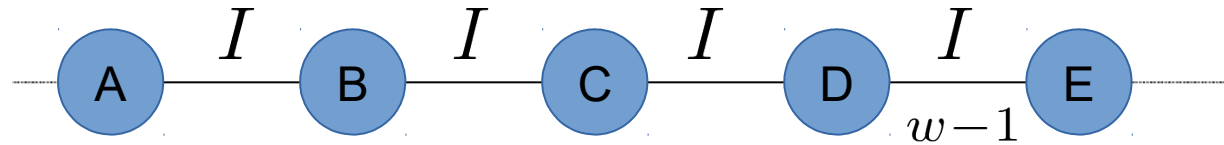
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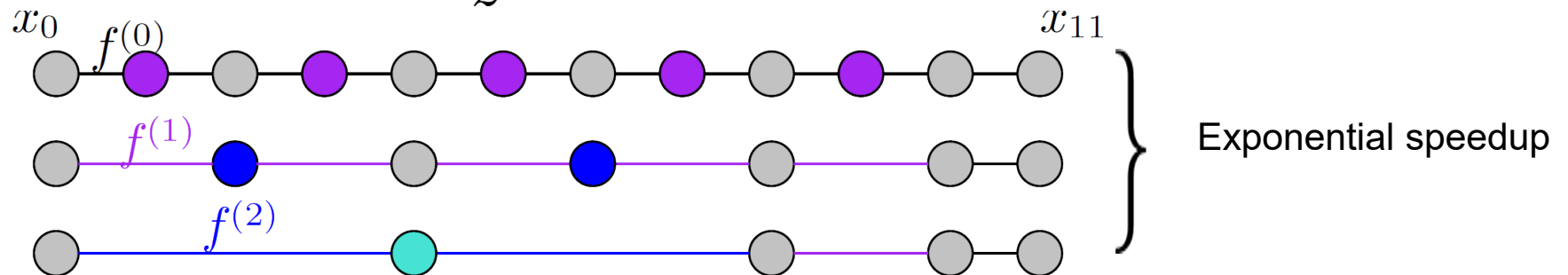
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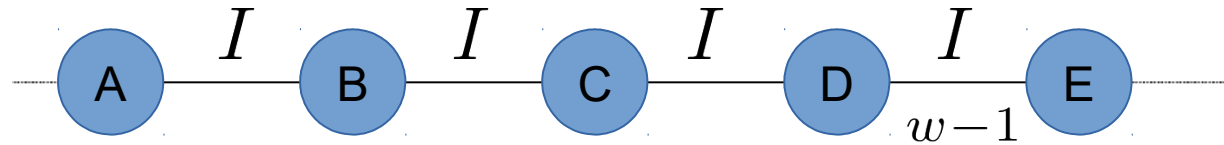
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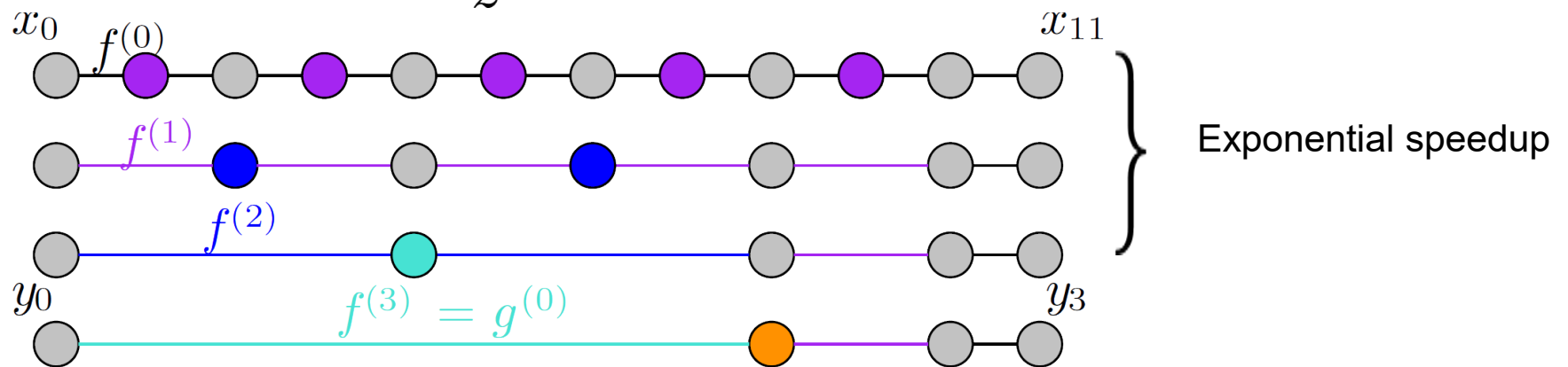
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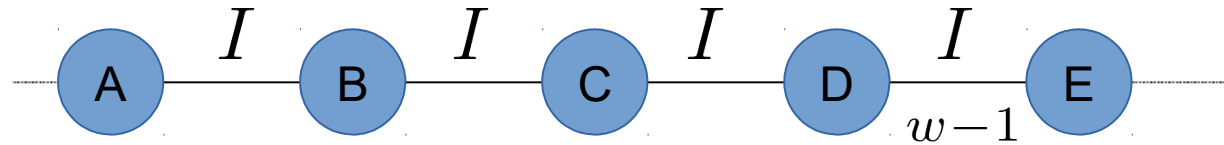
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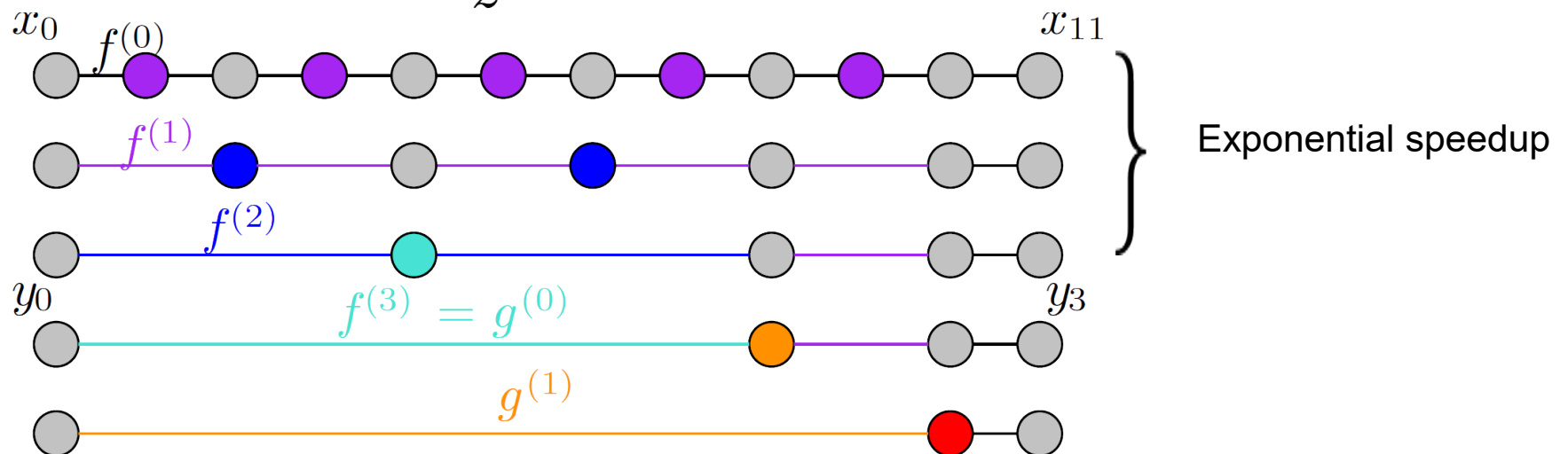
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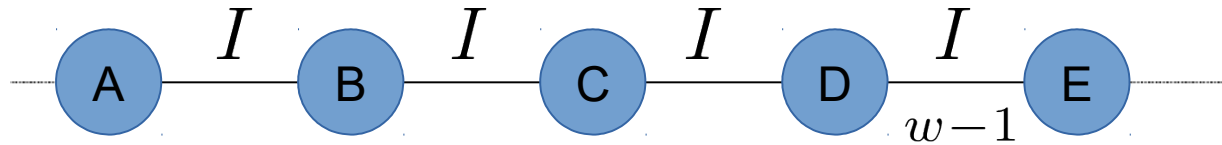
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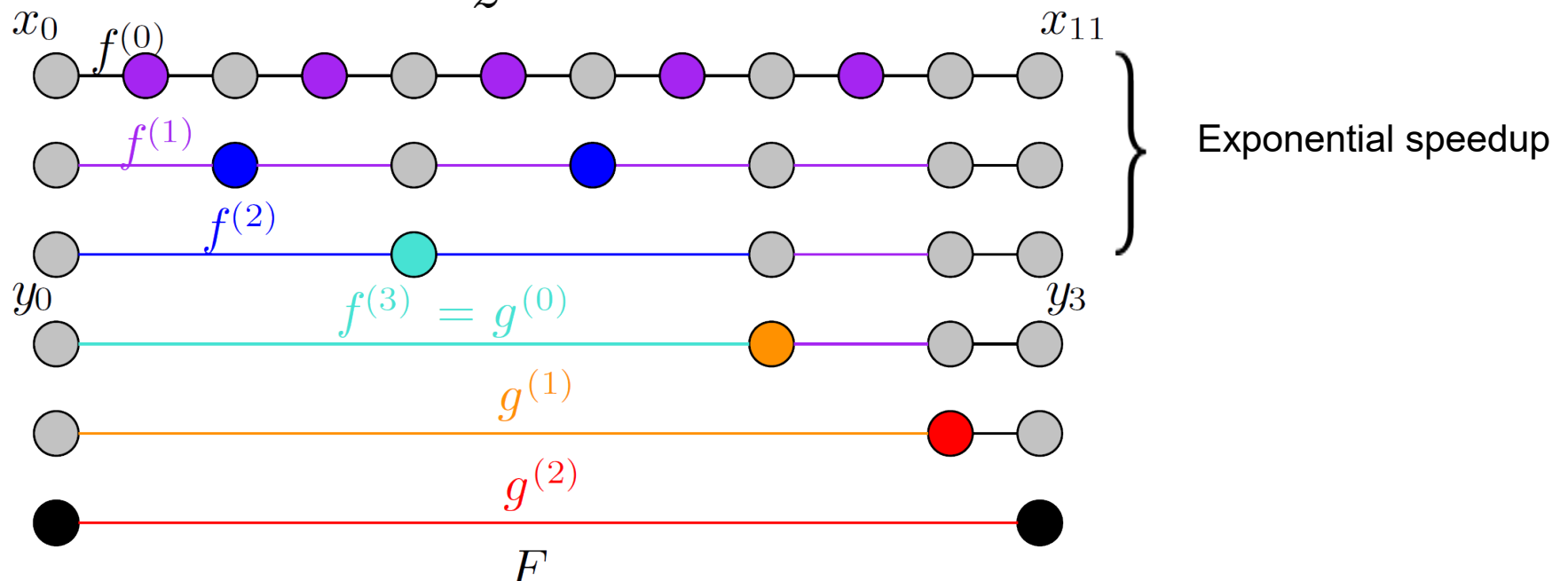
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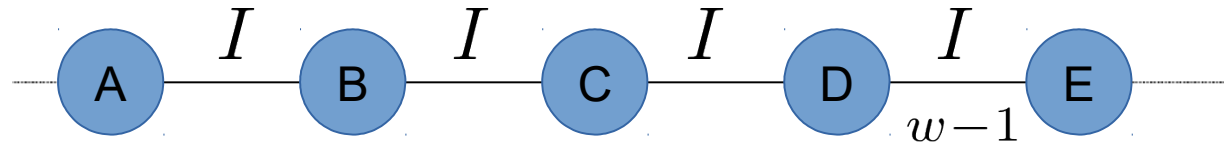
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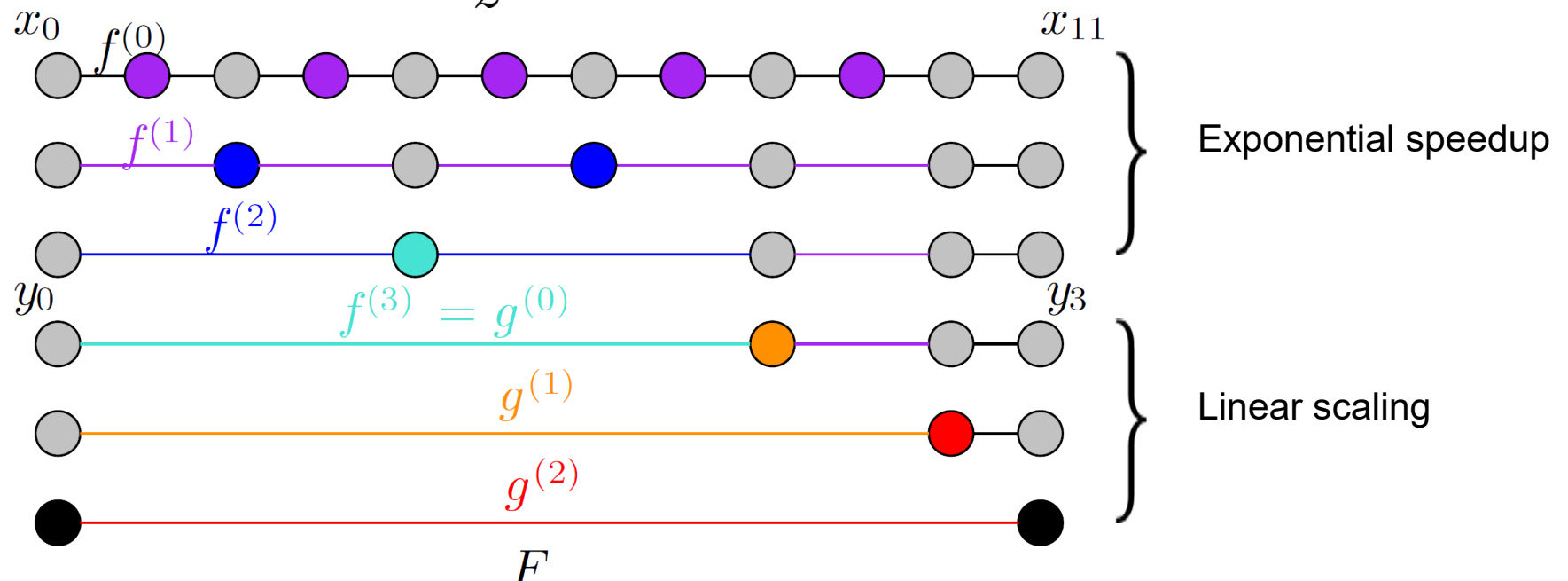
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# Application to an inequality with $R > 1$





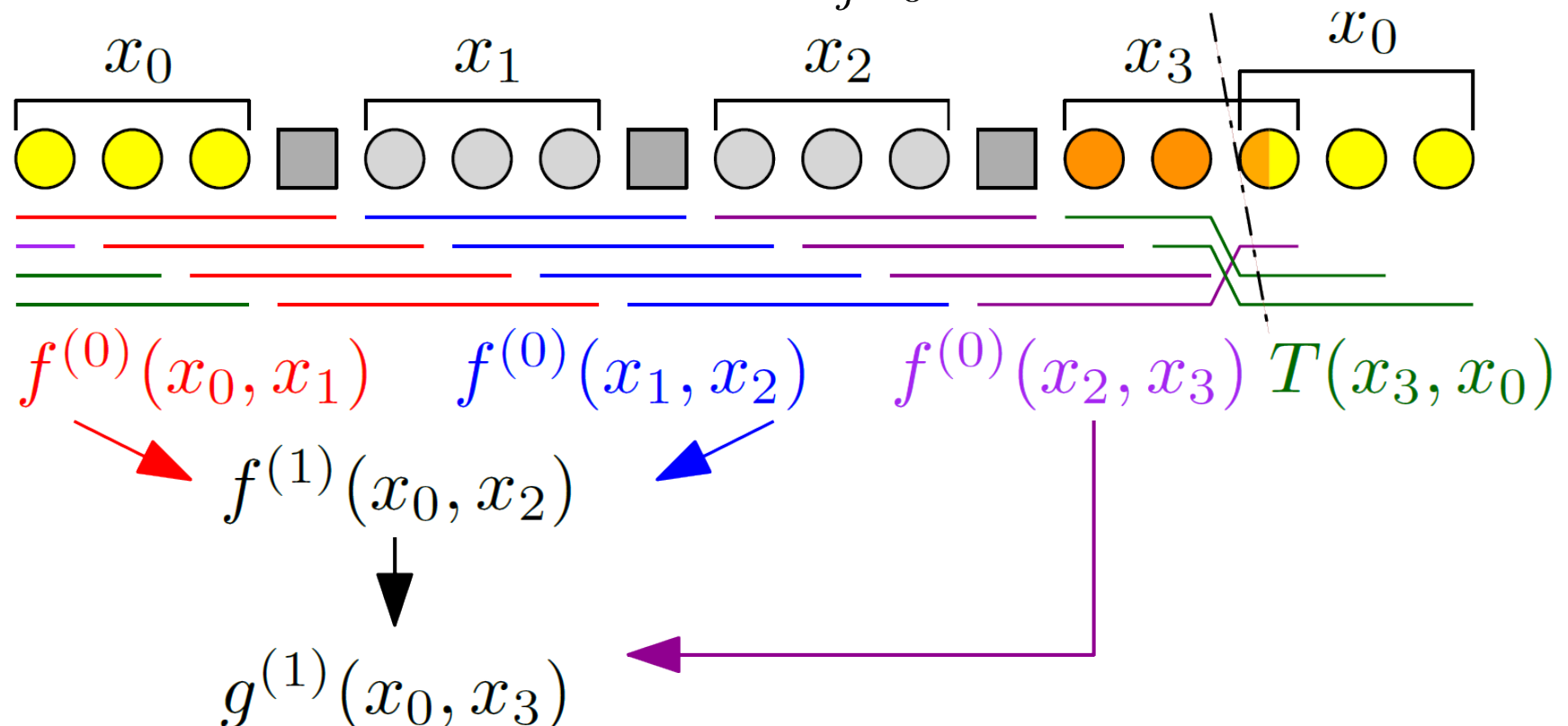
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- To reach the form  $F = \min_{x_0, \dots, x_w} \sum_{j=0}^{w-1} f^{(0)}(x_j, x_{j+1})$



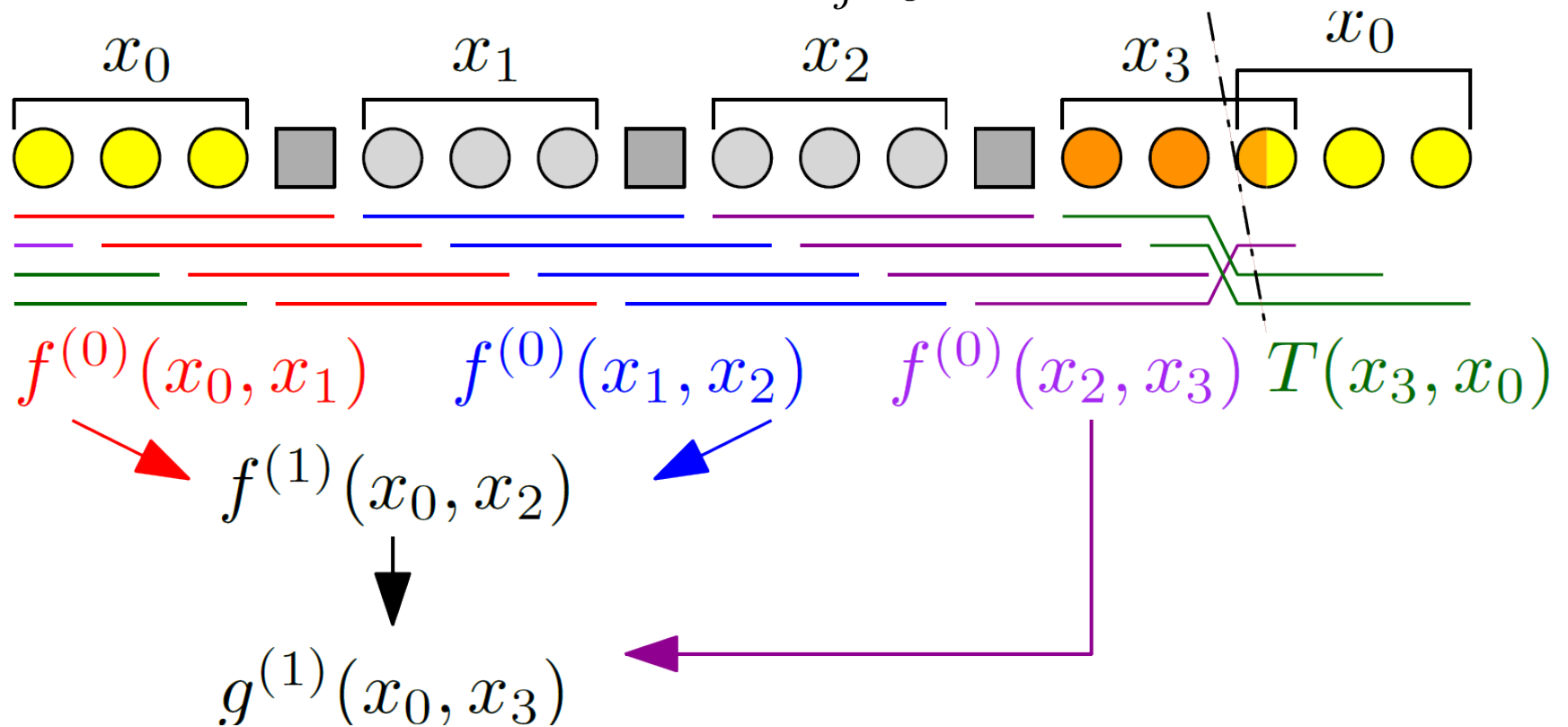
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$O(\log n)$  overall complexity

# Translationally invariant Hamiltonian (I)



# Translationally invariant Hamiltonian

$$\mathcal{H} = \sum_{i=0}^{n-1} \left( t^{(i)} \sigma_z^{(i)} + \sum_{r=1}^R \sum_{\alpha, \beta \in \{x, y\}} t_{\alpha, \beta}^{(i, r)} \text{Str}_{\alpha, \beta}^{(i, r)} \right) \quad (\text{I})$$



# Translationally invariant Hamiltonian

(I)

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- $t^{(i)}, t_{\alpha, \beta}^{(i, r)}$  are independent of  $i$



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# Translationally invariant Hamiltonian

(I)

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- $t^{(i)}$ ,  $t_{\alpha, \beta}^{(i, r)}$  are independent of  $i$

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- $H$  is real, anti-symmetric, **block-circulant**





# Translationally invariant Hamiltonian

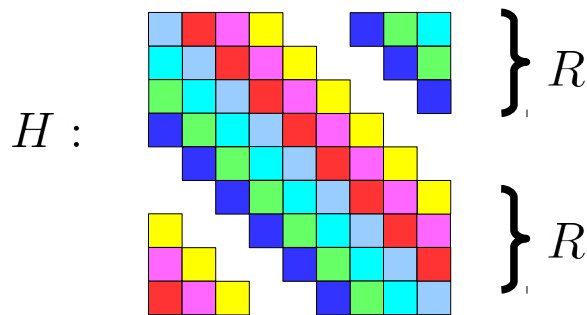
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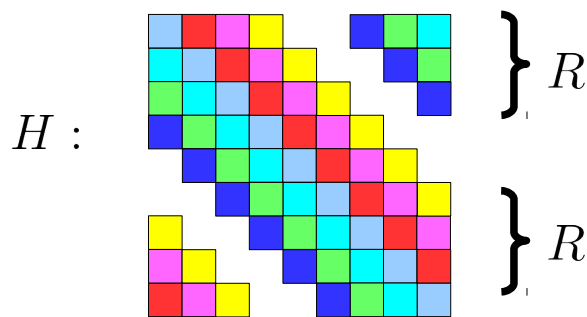
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- $H$  is real, anti-symmetric, **block-circulant**



If the fermion system has parity -1

# Translationally invariant Hamiltonian

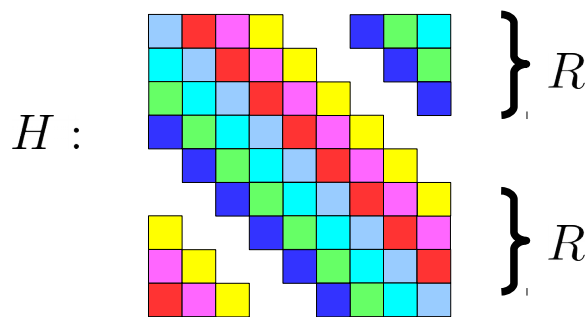
(I)

$$\mathcal{H} = \sum_{i=0}^{n-1} \left( t^{(i)} \sigma_z^{(i)} + \sum_{r=1}^R \sum_{\alpha, \beta \in \{x, y\}} t_{\alpha, \beta}^{(i, r)} \text{Str}_{\alpha, \beta}^{(i, r)} \right)$$

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# Translationally invariant Hamiltonian

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$$H : \left. \begin{array}{c} \begin{array}{cccccccc} \color{red}{\square} & \color{blue}{\square} & \color{green}{\square} & \color{magenta}{\square} & \color{yellow}{\square} & & & \color{cyan}{\square} \\ \color{cyan}{\square} & \color{red}{\square} & \color{blue}{\square} & \color{green}{\square} & \color{magenta}{\square} & \color{yellow}{\square} & & \color{cyan}{\square} \\ \color{green}{\square} & \color{cyan}{\square} & \color{red}{\square} & \color{blue}{\square} & \color{magenta}{\square} & \color{yellow}{\square} & & \color{cyan}{\square} \\ \color{blue}{\square} & \color{green}{\square} & \color{cyan}{\square} & \color{red}{\square} & \color{blue}{\square} & \color{magenta}{\square} & \color{yellow}{\square} & \color{cyan}{\square} \\ & \color{blue}{\square} & \color{green}{\square} & \color{cyan}{\square} & \color{red}{\square} & \color{blue}{\square} & \color{magenta}{\square} & \color{yellow}{\square} \\ & & \color{blue}{\square} & \color{green}{\square} & \color{cyan}{\square} & \color{red}{\square} & \color{blue}{\square} & \color{magenta}{\square} \\ & & & \color{blue}{\square} & \color{green}{\square} & \color{cyan}{\square} & \color{red}{\square} & \color{blue}{\square} \\ & & & & \color{blue}{\square} & \color{green}{\square} & \color{cyan}{\square} & \color{red}{\square} \\ & & & & & \color{blue}{\square} & \color{green}{\square} & \color{cyan}{\square} \\ & & & & & & \color{blue}{\square} & \color{green}{\square} \\ & & & & & & & \color{blue}{\square} \end{array} \\ \color{red}{\square} & \color{blue}{\square} & \color{green}{\square} & \color{magenta}{\square} & \color{yellow}{\square} & & & \color{cyan}{\square} \\ \color{cyan}{\square} & \color{red}{\square} & \color{blue}{\square} & \color{green}{\square} & \color{magenta}{\square} & \color{yellow}{\square} & & \color{cyan}{\square} \\ \color{green}{\square} & \color{cyan}{\square} & \color{red}{\square} & \color{blue}{\square} & \color{magenta}{\square} & \color{yellow}{\square} & & \color{cyan}{\square} \\ \color{blue}{\square} & \color{green}{\square} & \color{cyan}{\square} & \color{red}{\square} & \color{blue}{\square} & \color{magenta}{\square} & \color{yellow}{\square} & \color{cyan}{\square} \\ & \color{blue}{\square} & \color{green}{\square} & \color{cyan}{\square} & \color{red}{\square} & \color{blue}{\square} & \color{magenta}{\square} & \color{yellow}{\square} \\ & & \color{blue}{\square} & \color{green}{\square} & \color{cyan}{\square} & \color{red}{\square} & \color{blue}{\square} & \color{magenta}{\square} \\ & & & \color{blue}{\square} & \color{green}{\square} & \color{cyan}{\square} & \color{red}{\square} & \color{blue}{\square} \\ & & & & \color{blue}{\square} & \color{green}{\square} & \color{cyan}{\square} & \color{red}{\square} \\ & & & & & \color{blue}{\square} & \color{green}{\square} & \color{cyan}{\square} \\ & & & & & & \color{blue}{\square} & \color{green}{\square} \\ & & & & & & & \color{blue}{\square} \end{array} \right\} R$$

If the fermion system has parity -1  
Discrete Fourier Transform will diagonalize it

$$(\mathcal{F}_n)_{kl} := \frac{1}{\sqrt{n}} \omega^{k \cdot l}, \quad \omega^n = 1$$



# Translationally invariant Hamiltonian (II)

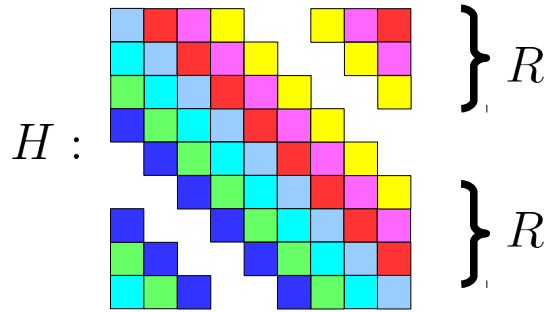


# Translationally invariant Hamiltonian (II)

If the fermion system has parity 1



# Translationally invariant Hamiltonian (II)

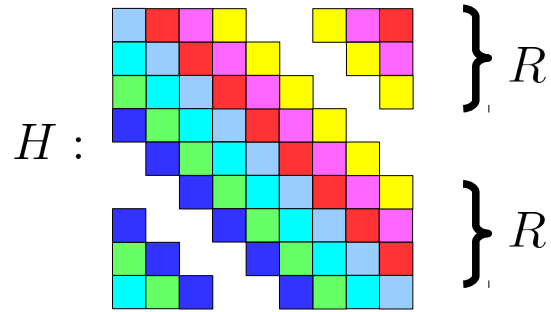


If the fermion system has parity 1  
it is no longer circulant, but



# Translationally invariant Hamiltonian

(II)



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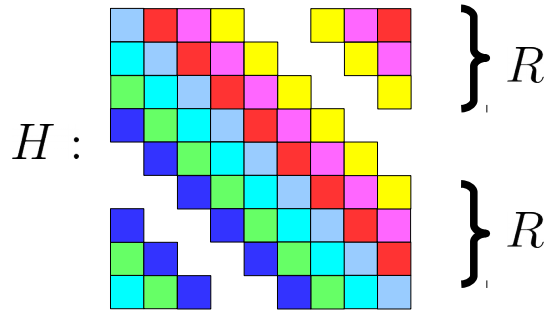
$$H \longrightarrow \begin{pmatrix} H & -H \\ -H & H \end{pmatrix} \text{ is.}$$





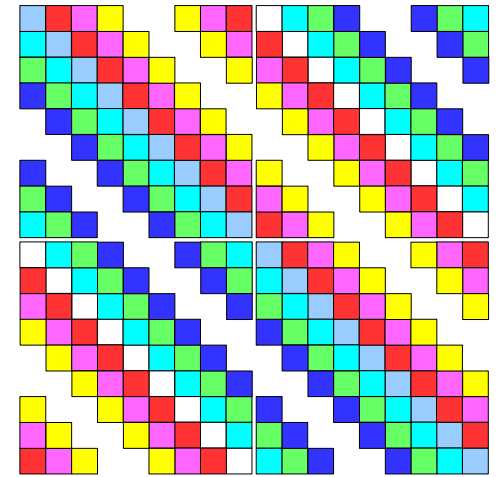
# Translationally invariant Hamiltonian

(II)



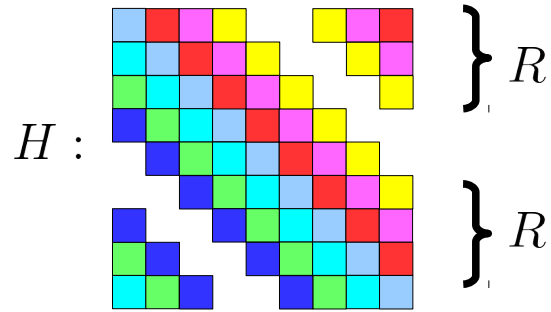
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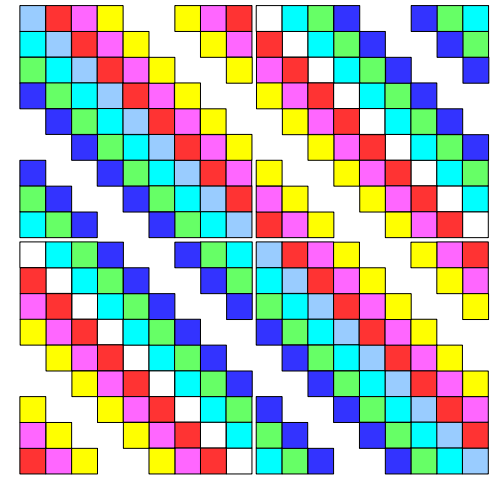
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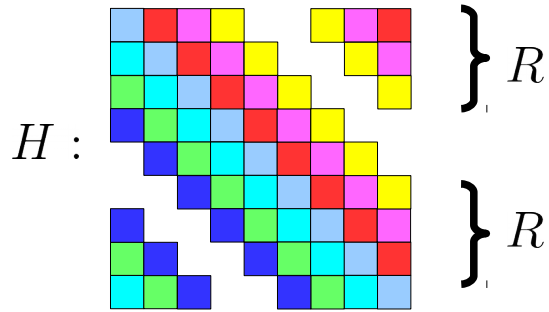
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Diagonalizable using  $\mathcal{F}_{2n}$



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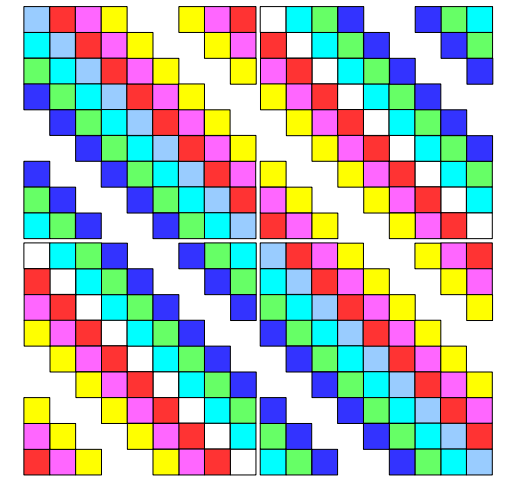
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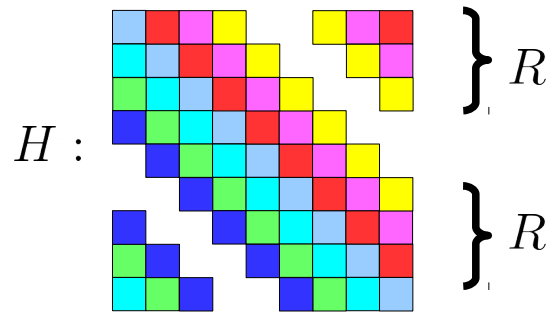
$$\text{diag}(\zeta^0, \dots, \zeta^{n-1}) \cdot \mathcal{F}_n$$

$\zeta^{2n} = 1$  Block-diagonalizes  $H$



# Translationally invariant Hamiltonian

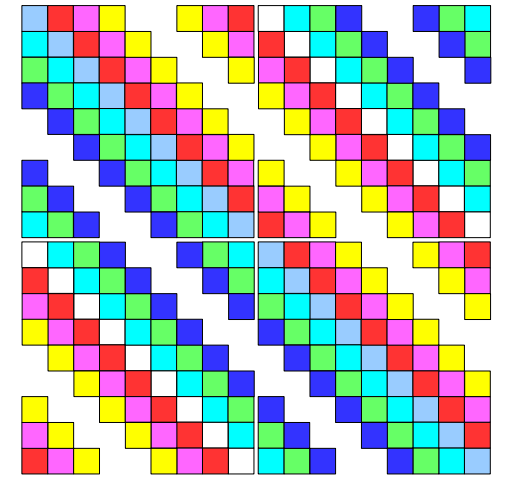
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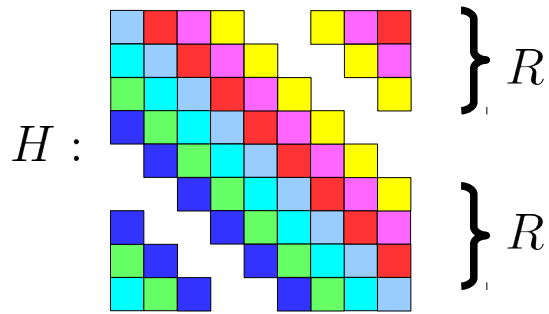
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- Simple super-selection rule

$$p = (-1)^{\lfloor \frac{n+(p-1)/2}{2} \rfloor} \prod_{k=0}^{n-1} s_k$$

# Translationally invariant Hamiltonian

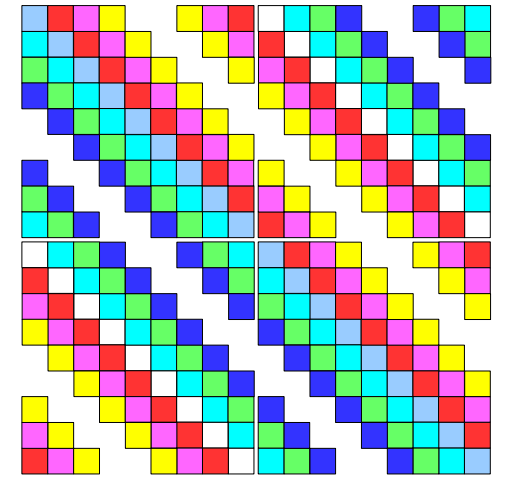
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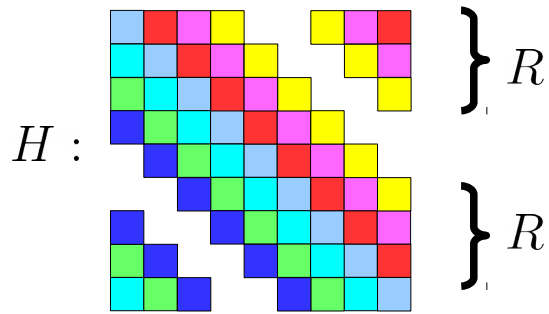
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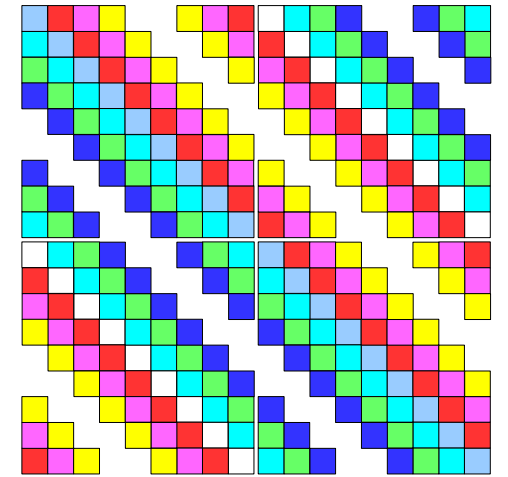
# Translationally invariant Hamiltonian

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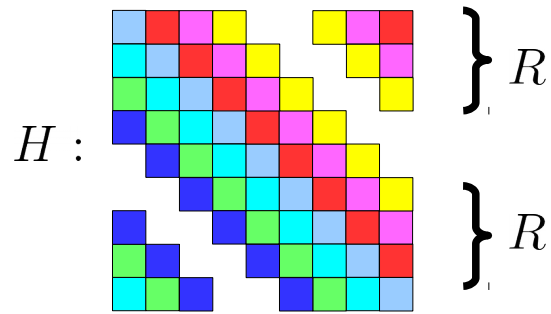
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- Analytical solution

$$\varepsilon_{k,\pm} = a_k + c_k \pm \sqrt{(a_k - c_k)^2 + 4(b_k^2 + x_k^2)}$$

# Translationally invariant Hamiltonian

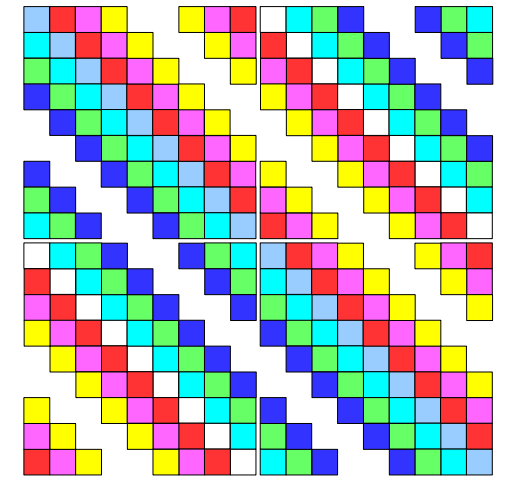
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$$\begin{cases} x_k &= H_{00;01} + \sum_{r=1}^R \cos(\nu_{k,r})(H_{00;r1} - H_{01;r0}) \\ a_k &= -2 \sum_{r=1}^R \sin(\nu_{k,r}) H_{00;r0} \\ b_k &= -\sum_{r=1}^R \sin(\nu_{k,r})(H_{00;r1} + H_{01;r0}) \\ c_k &= -2 \sum_{r=1}^R \sin(\nu_{k,r}) H_{11;r0} \end{cases}$$

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# Examples (Ia)





# Examples (Ia)

- The projected polytope approach



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- The projected polytope approach

Finding all Bell inequalities  $\longleftrightarrow$  Convex Hull problem



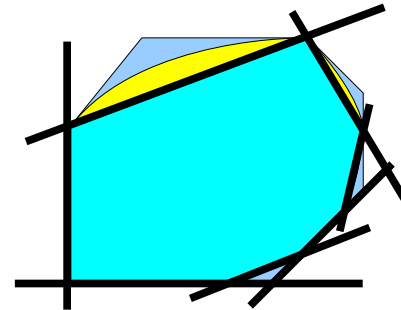
# Examples (1a)

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Finding all Bell inequalities



Convex Hull problem

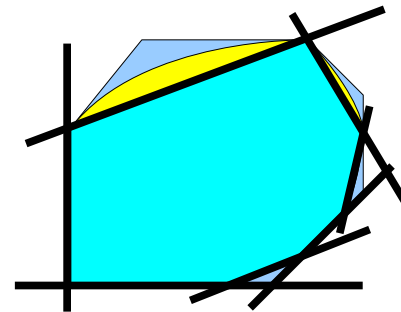


# Examples (Ia)

- The projected polytope approach

Finding all Bell inequalities  $\longleftrightarrow$  Convex Hull problem

$(n, m, d)$  scenario



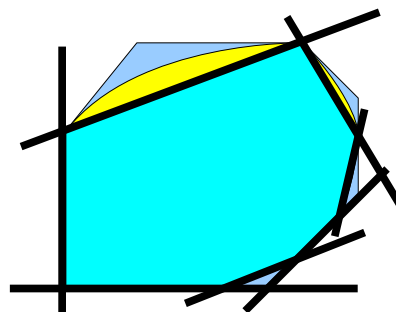
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Dimension of the Local Polytope  $D \approx (md)^n$



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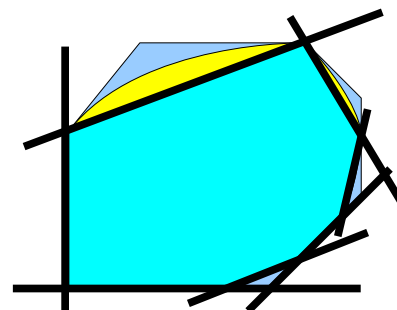
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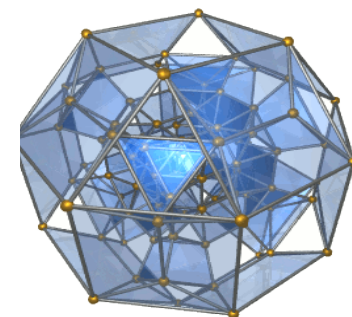
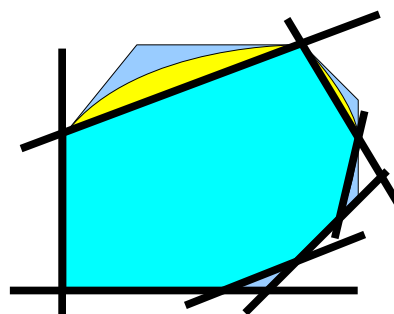
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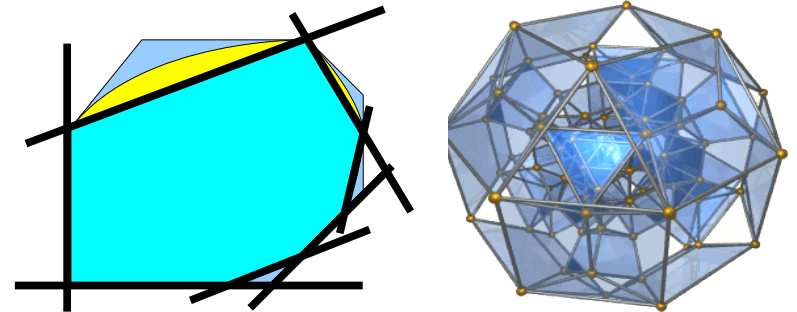
$(n, m, d)$  scenario

Dimension of the Local Polytope  $D \approx (md)^n$

Number of vertices  $v = d^{mn}$

Complexity of dual description:  $O(v^{\lfloor D/2 \rfloor} + v \log v)$

[B. Chazelle, *An optimal convex hull algorithm in any fixed dimension*, *Discrete Comput. Geom.* **10** 377409 (1993)]





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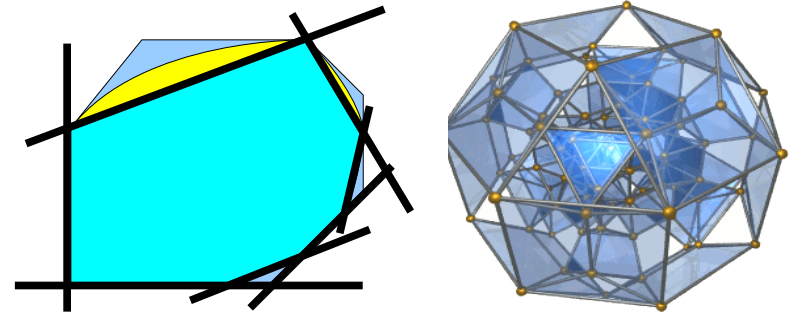
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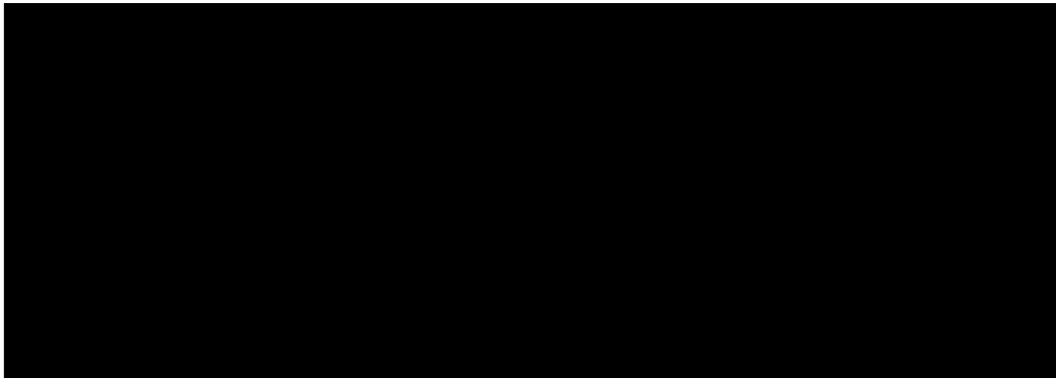
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## Examples



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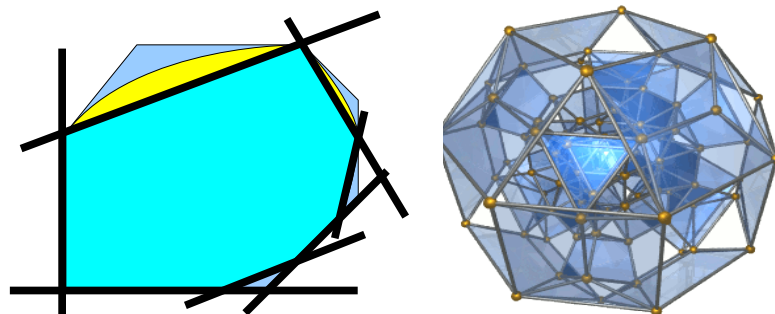
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## Examples

$(2, 2, 2) \rightarrow O(\text{ms})$



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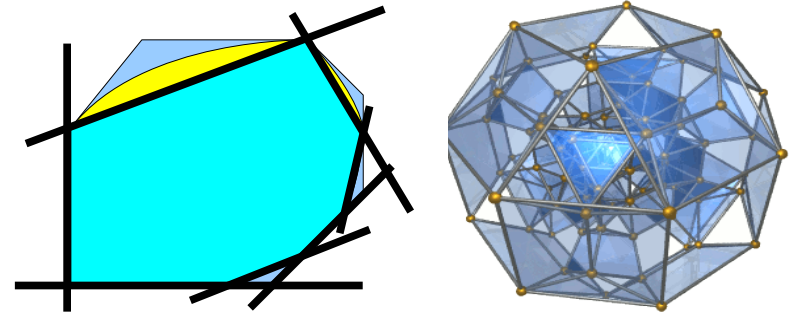
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## Examples

$(2, 2, 2) \longrightarrow O(\text{ms})$   
 $(3, 2, 2) \longrightarrow 5'$

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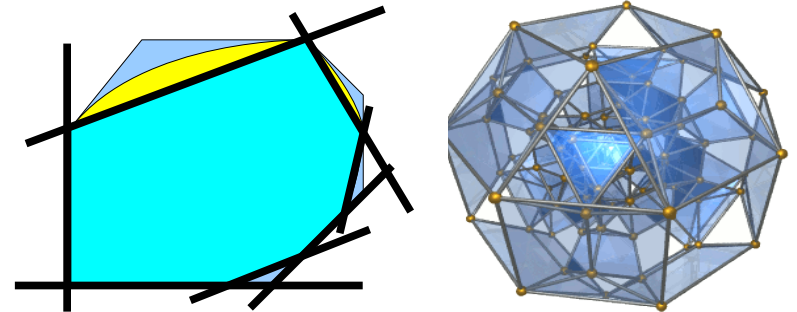
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## Examples

$(2, 2, 2) \longrightarrow O(\text{ms})$

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$(4, 2, 2) \longrightarrow 10^{67} \text{ years}$



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Finding all Bell inequalities  $\longleftrightarrow$  Convex Hull problem

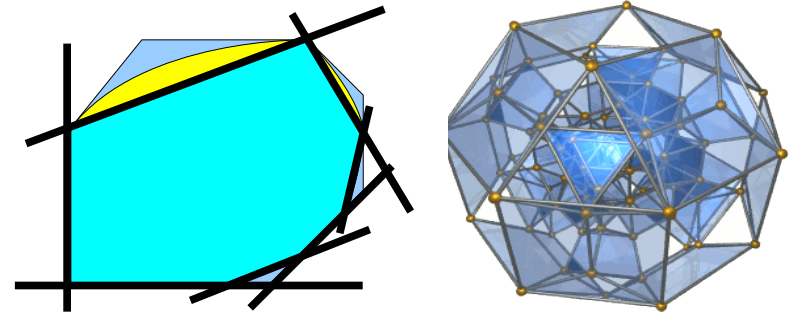
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## Examples

$(2, 2, 2) \longrightarrow O(\text{ms})$   
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 $\vdots$



# Examples (Ia)

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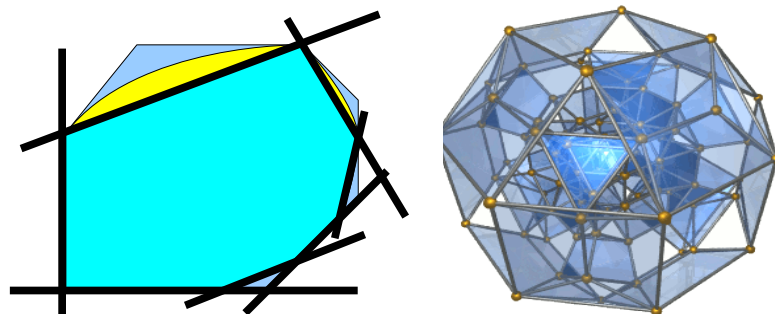
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Number of vertices  $v = d^{mn}$

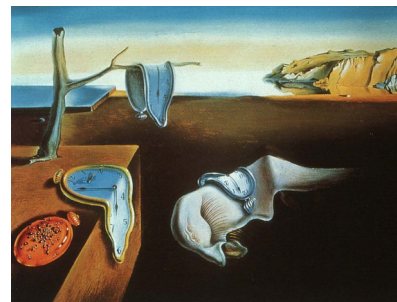
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[B. Chazelle, *An optimal convex hull algorithm in any fixed dimension*, *Discrete Comput. Geom.* **10** 377409 (1993)]



## Examples

$(2, 2, 2) \rightarrow O(\text{ms})$   
 $(3, 2, 2) \rightarrow 5'$   
 $(4, 2, 2) \rightarrow 10^{67}$  years  
 $\vdots$   
 $(10^4, 2, 2) \rightarrow 10^{10^{10^{4.67867\dots}}}$  basically any timescale you want



[S. Dalí *The persistence of memory* (1931)]



# Examples (Ib)





# Examples (Ib)

- Projecting  $\mathbb{P}_L$  to the space of few-body, TI BI





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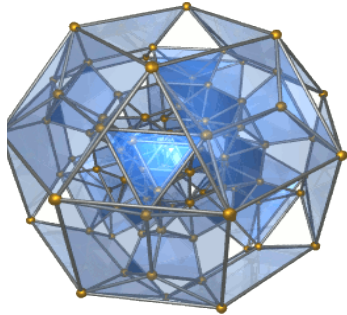
$$I = \gamma \mathcal{T}_2 + \sum_{k,l \in \{0,1\}} (\gamma_{k,l} \mathcal{T}_{k,l} + \gamma_{k,2,l} \mathcal{T}_{k,2,l}) \quad T_{k_1, \dots, k_r} = \sum_{i=0}^{n-1} M_{(k_1, \dots, k_r)}^{(i,r)}$$



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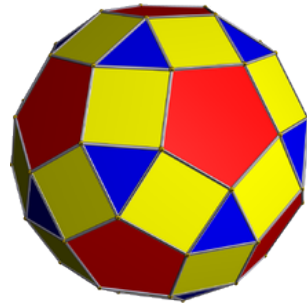
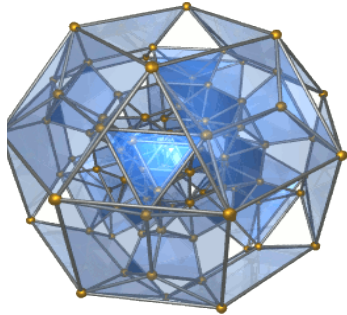
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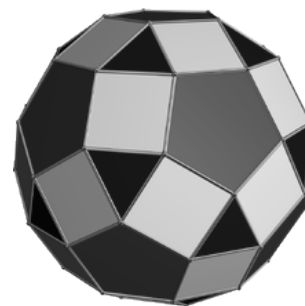
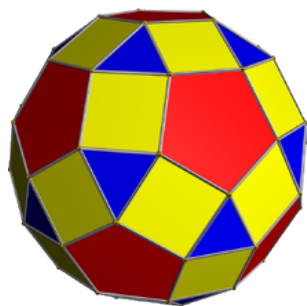
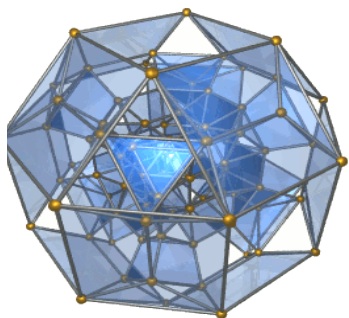
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# Examples (Ib)

- Projecting  $\mathbb{P}_L$  to the space of few-body, **TI/BI**

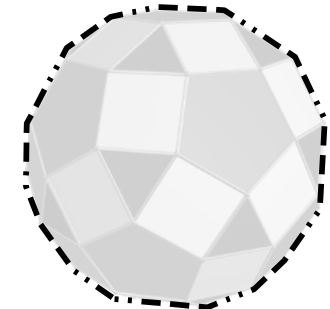
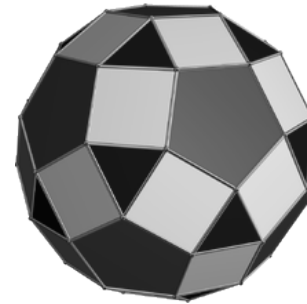
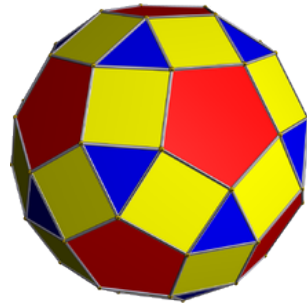
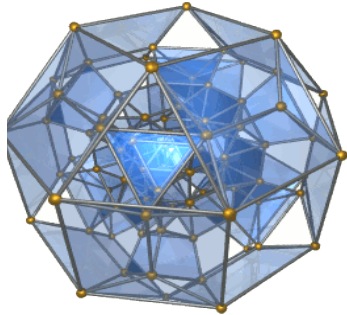
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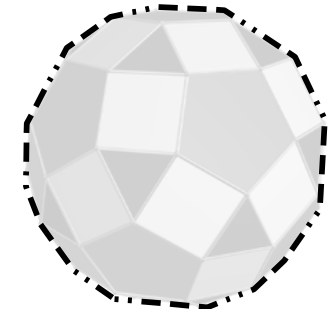
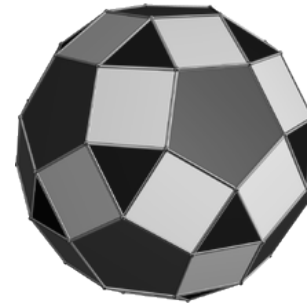
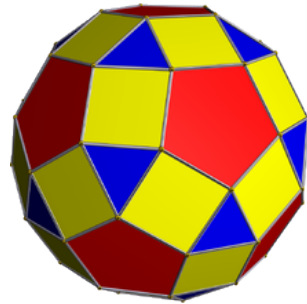
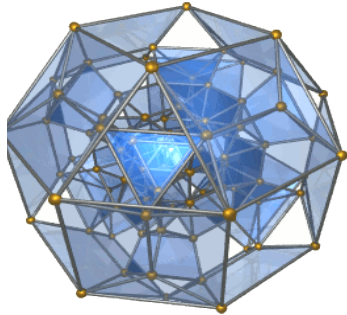
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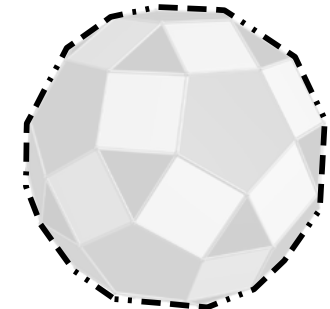
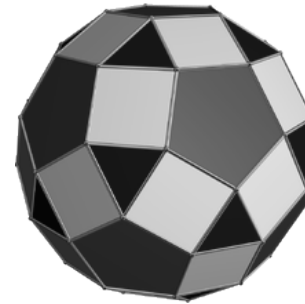
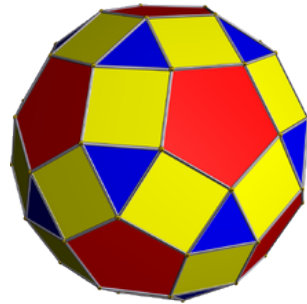
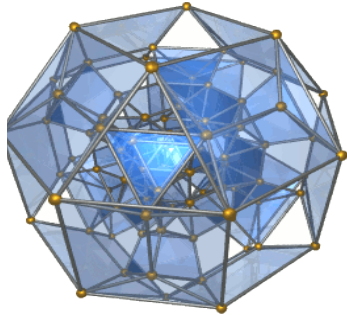
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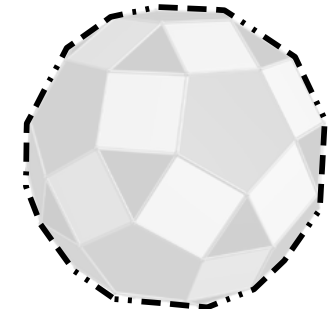
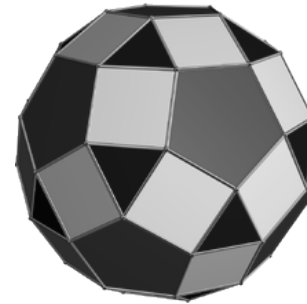
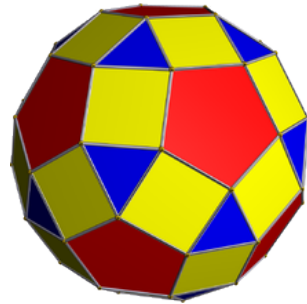
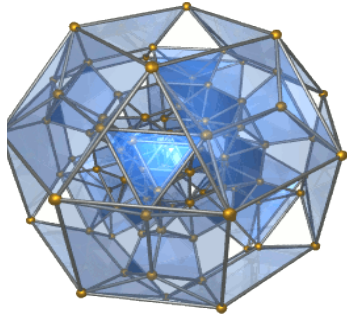
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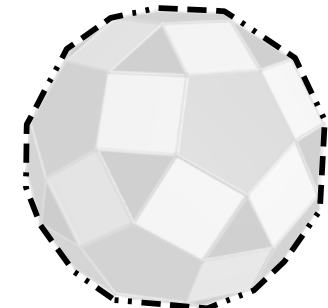
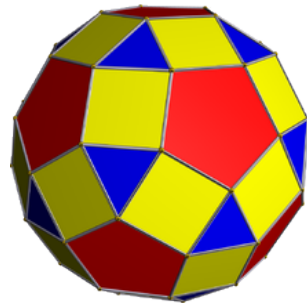
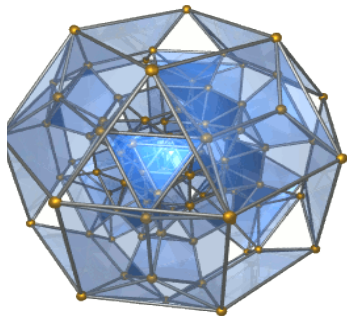




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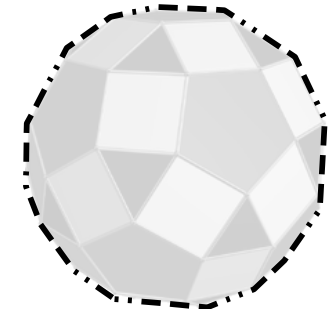
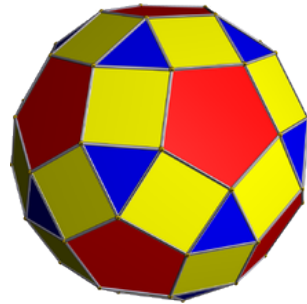
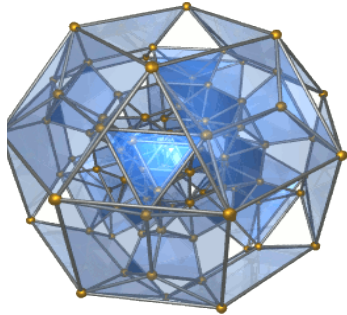
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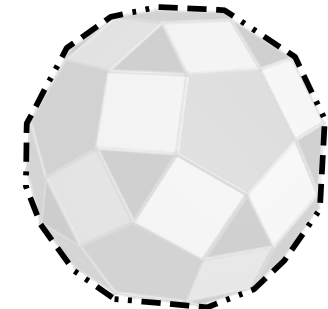
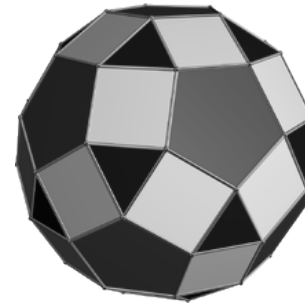
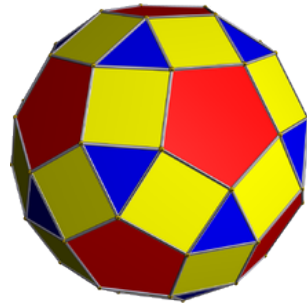
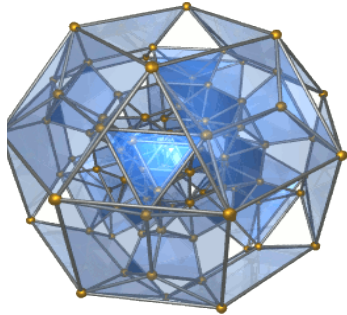
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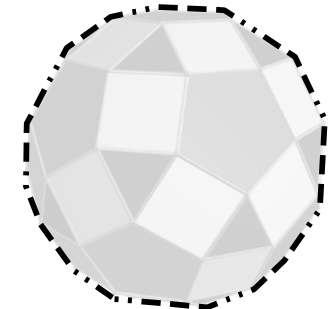
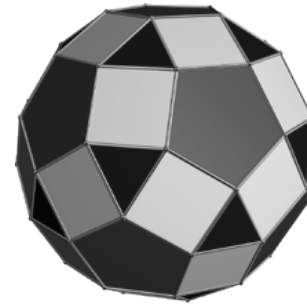
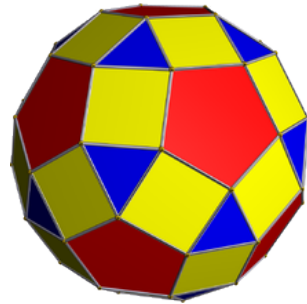
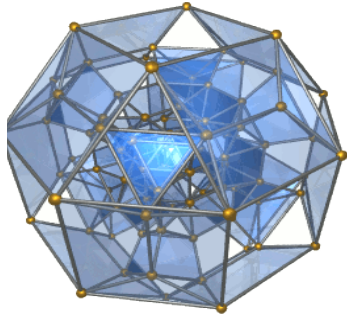
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# Examples (IIa)



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- Building a quasi-TI class



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  - Uniparametric  $\varepsilon$



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  - Uniparametric  $\varepsilon$
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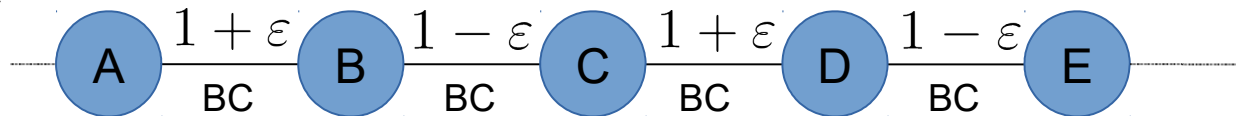




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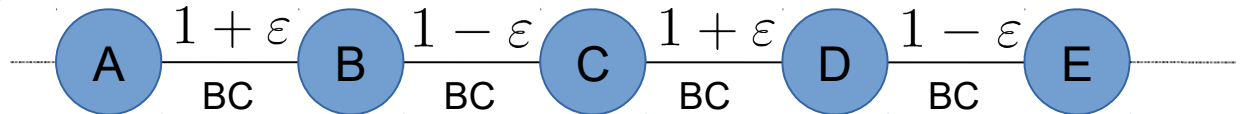


# Examples (IIa)

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Take the Braunstein-Caves (BC) chained inequality for  $m$  measurement settings

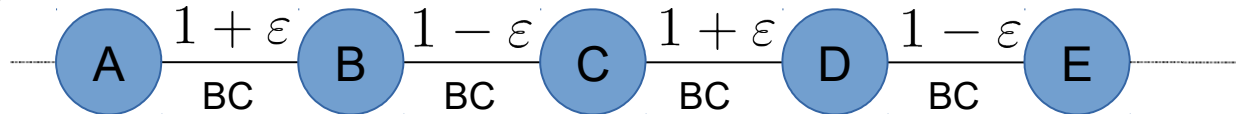
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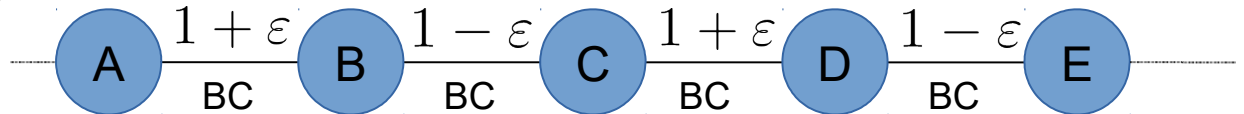
$$I_{BC} = \sum_{k=0}^{m-1} (A_{m-k-2} B_k + A_{m-k-1} B_1) \quad [Braunstein and Caves, Ann. Phys. \mathbf{202}, 22 (1990)] \quad A_{-1} = -A_{m-1}$$



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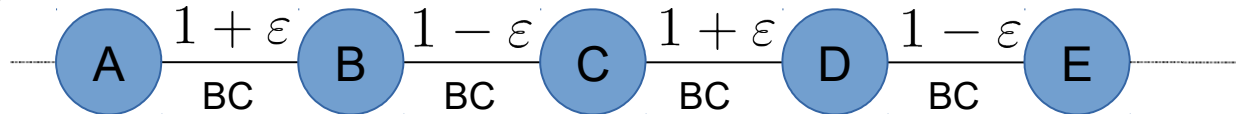
For  $m = 2$ , it is simply the CHSH inequality  $A_0 B_0 + A_0 B_1 + A_1 B_0 - A_1 B_1$   
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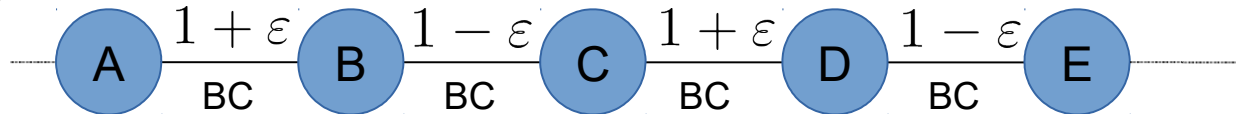
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- Always nonlocal when  $\varepsilon = \pm 1$
- Monogamy of correlations dominates when  $\varepsilon = 0$

[Wang et al., arXiv:1608.03485v3 (2016)]

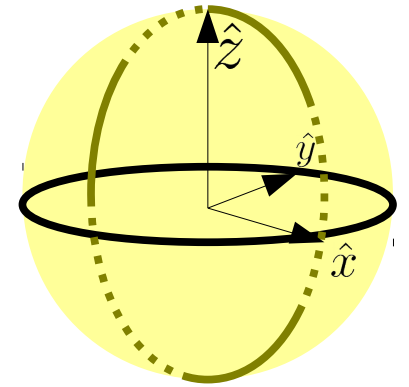


# Examples (IIb)



# Examples (IIb)

- Bell operator is an XY-like Hamiltonian

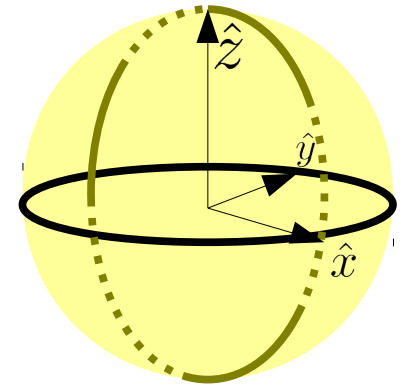




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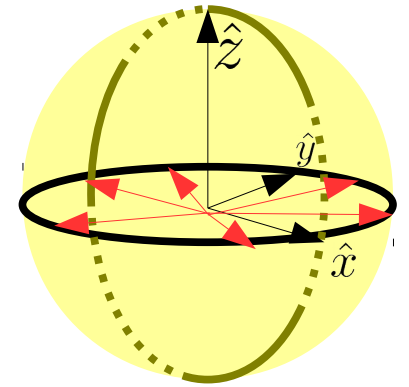
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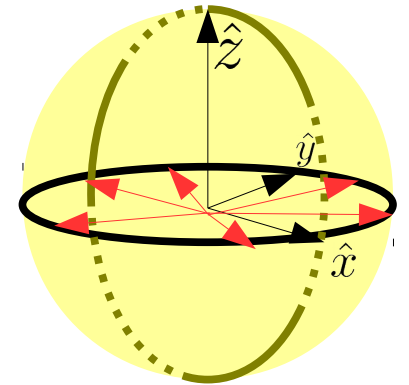


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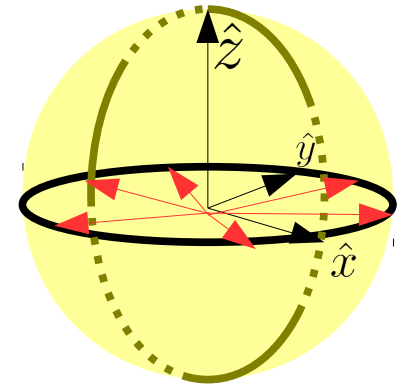
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Asymptotic contributions per particle  
to quantum value and classical bound



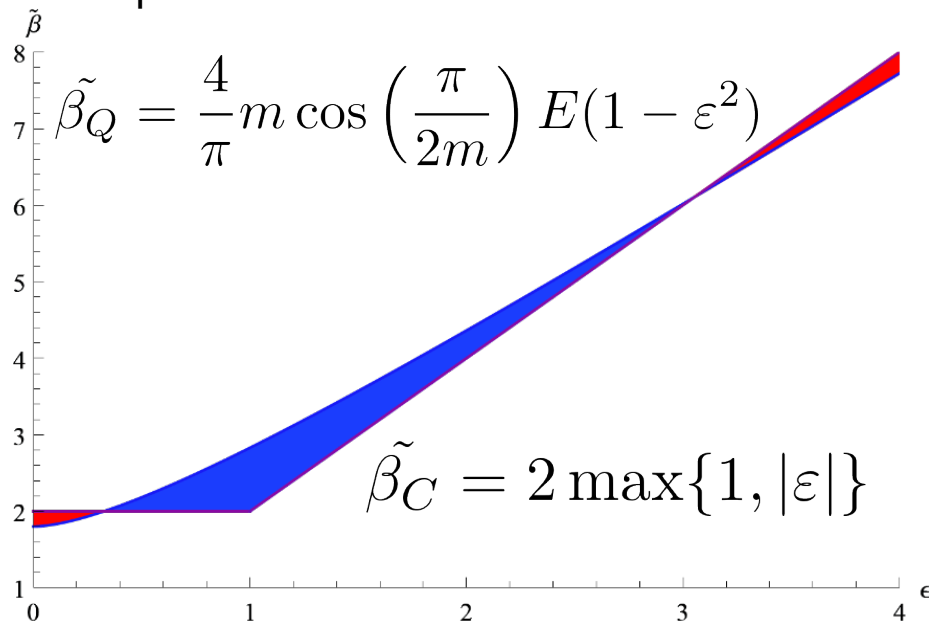
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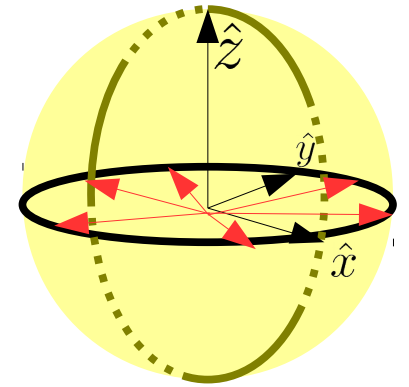
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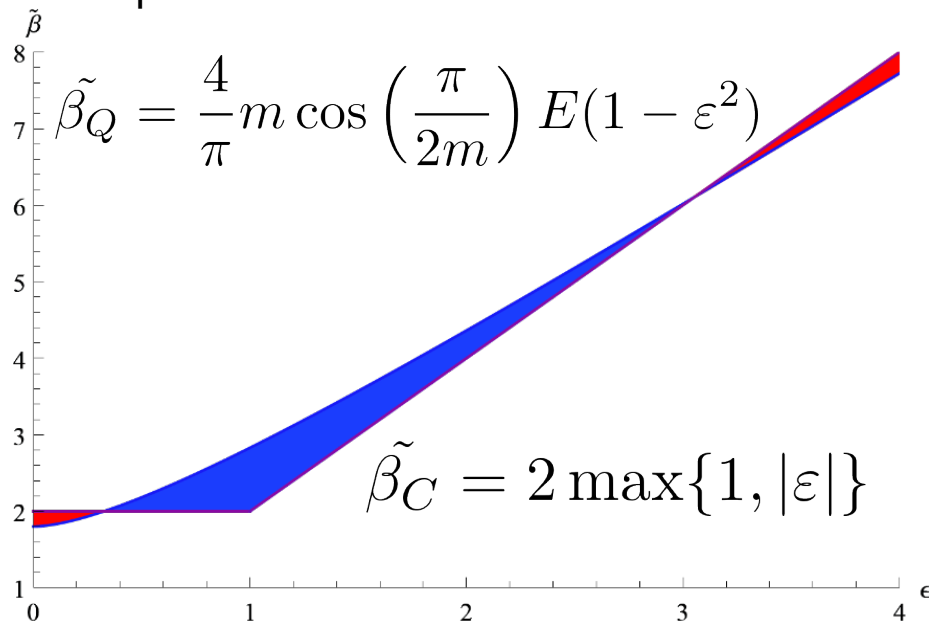
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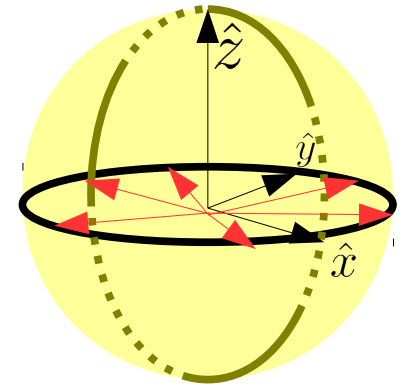
$E(t) \rightarrow$  Elliptic integral of  
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Ground state is nonlocal  
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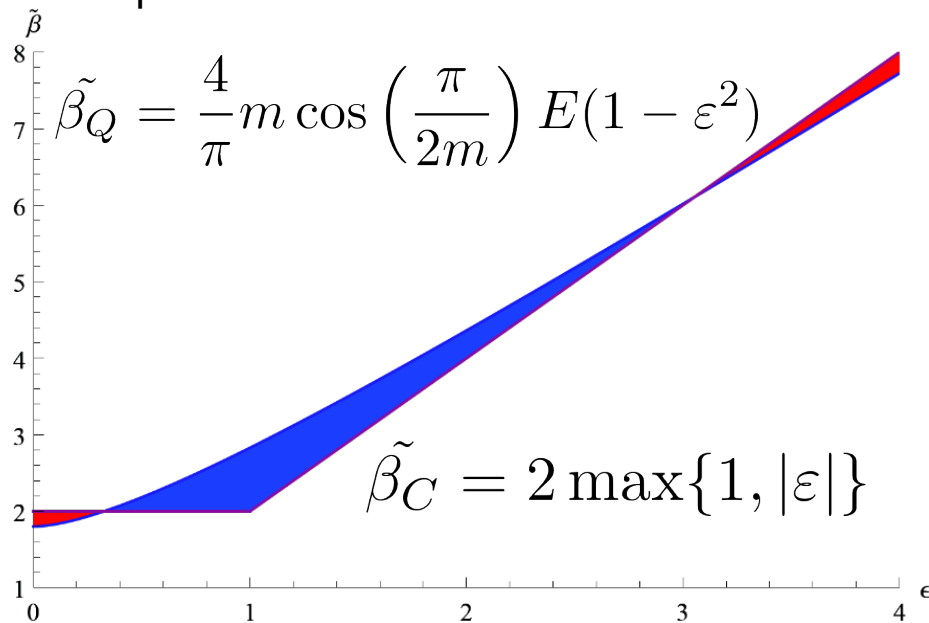
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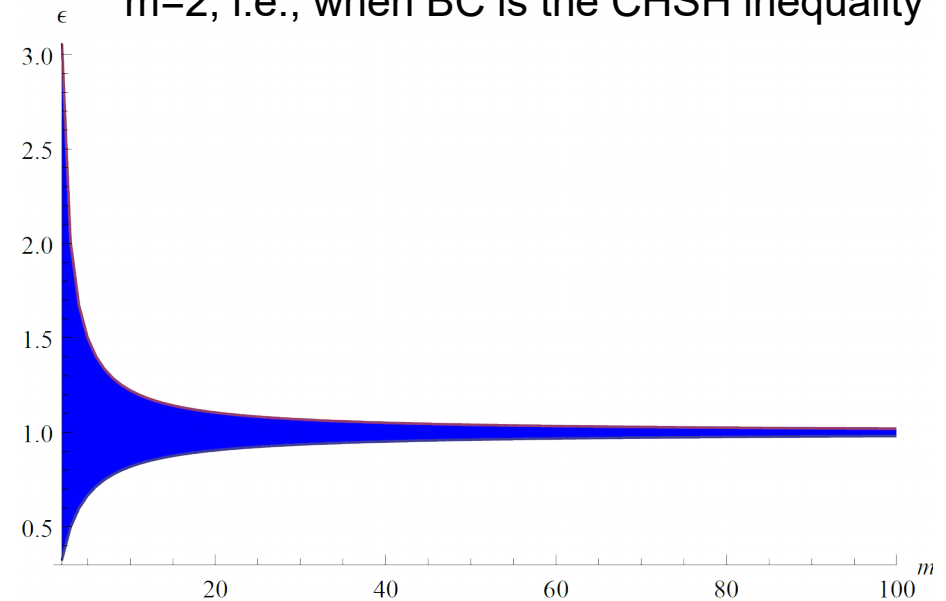


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Asymptotic contributions per particle to quantum value and classical bound



The optimal number of measurements is  $m=2$ , i.e., when BC is the CHSH inequality



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# Examples (III)



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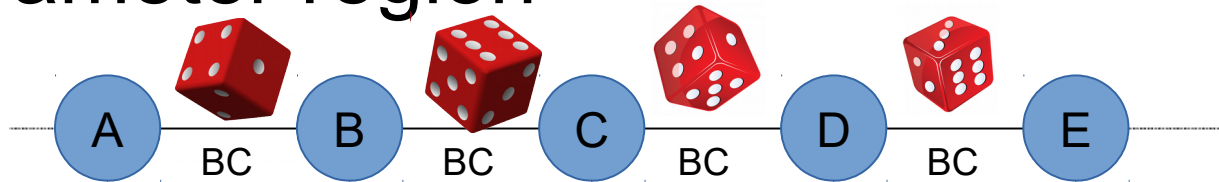
- Spin glass displays Bell correlations in some parameter region





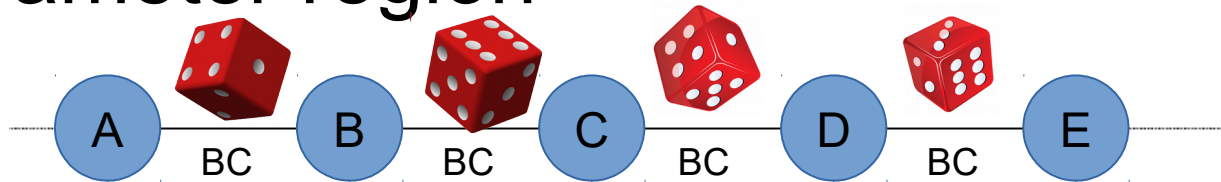
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# Examples (III)

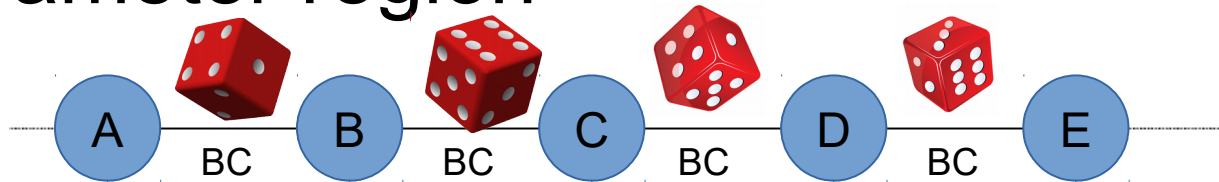
- Spin glass displays Bell correlations in some parameter region



$$\mathcal{H} = \sum_{i=0}^{n-1} J_{\mu,\sigma}^{(i)} \left( \sigma_{\pi/4}^{(i)} \sigma_{\pi/4}^{(i+1)} - \sigma_y^{(i)} \sigma_y^{(i+1)} \right)$$

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100 spins

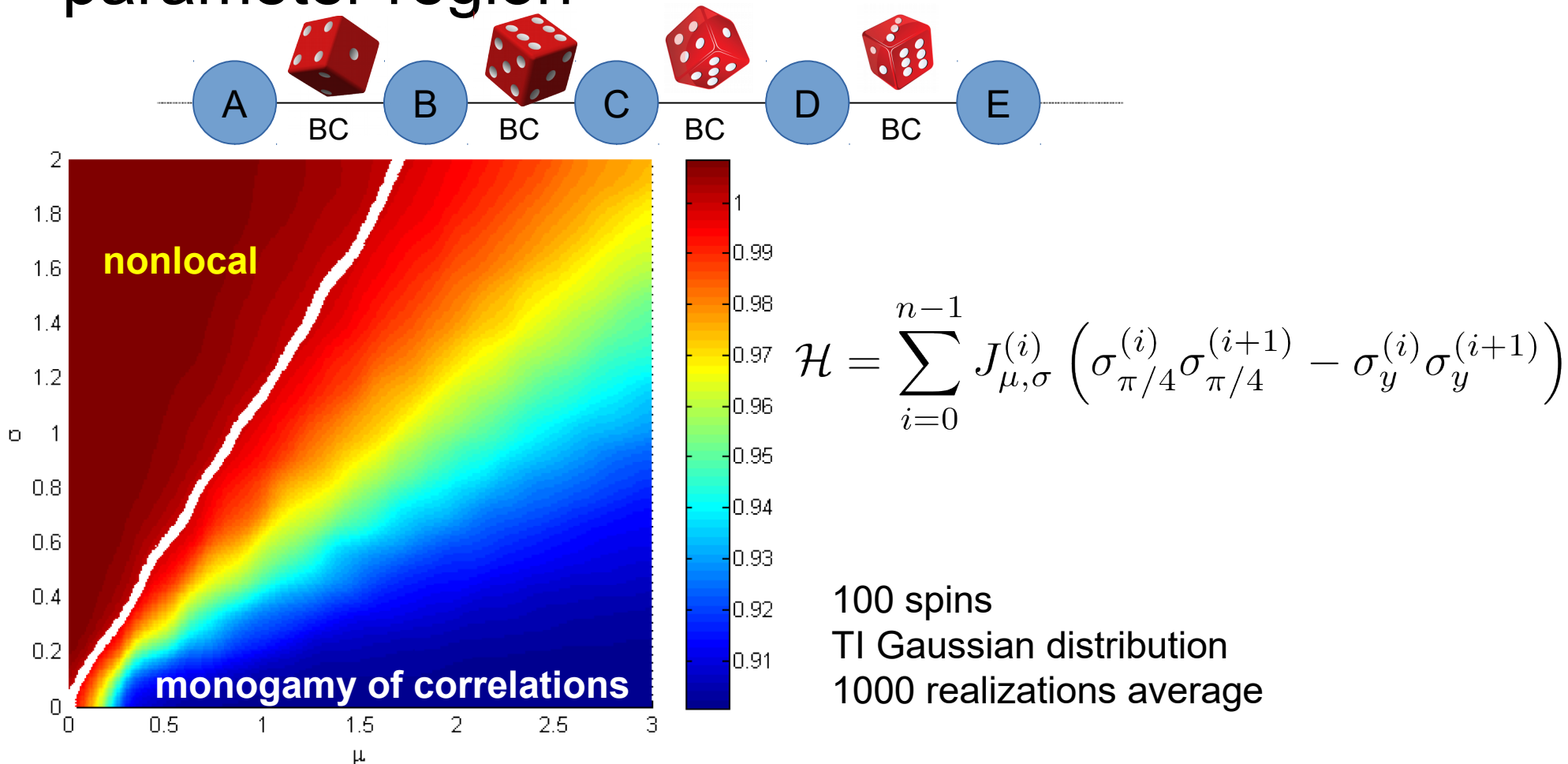
TI Gaussian distribution

1000 realizations average



# Examples (III)

- Spin glass displays Bell correlations in some parameter region



# Let's generalize



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- Up to one's imagination!



# Examples (IVa)



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- The XXZ-model and Gisin's *elegant* inequality



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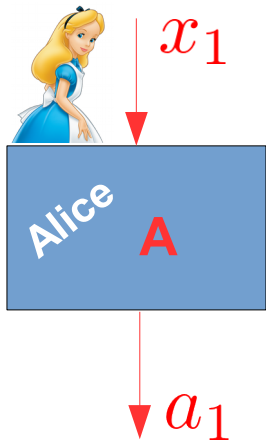
$$I = ( A_0 \quad A_1 \quad A_2 \quad A_3 ) \begin{pmatrix} 1 & 1 & \Delta \\ 1 & -1 & -\Delta \\ -1 & 1 & -\Delta \\ -1 & -1 & \Delta \end{pmatrix} \begin{pmatrix} B_0 \\ B_1 \\ B_2 \end{pmatrix}$$



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$$I = \left( \begin{array}{cccc} A_0 & A_1 & A_2 & A_3 \end{array} \right) \left( \begin{array}{ccc} 1 & 1 & \Delta \\ 1 & -1 & -\Delta \\ -1 & 1 & -\Delta \\ -1 & -1 & \Delta \end{array} \right) \left( \begin{array}{c} B_0 \\ B_1 \\ B_2 \end{array} \right)$$

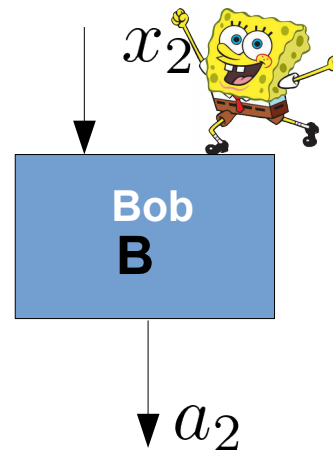
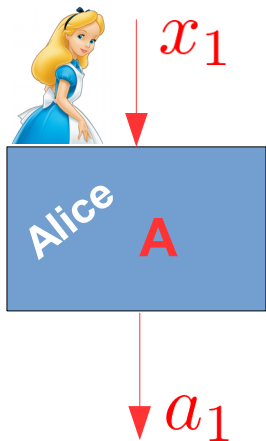




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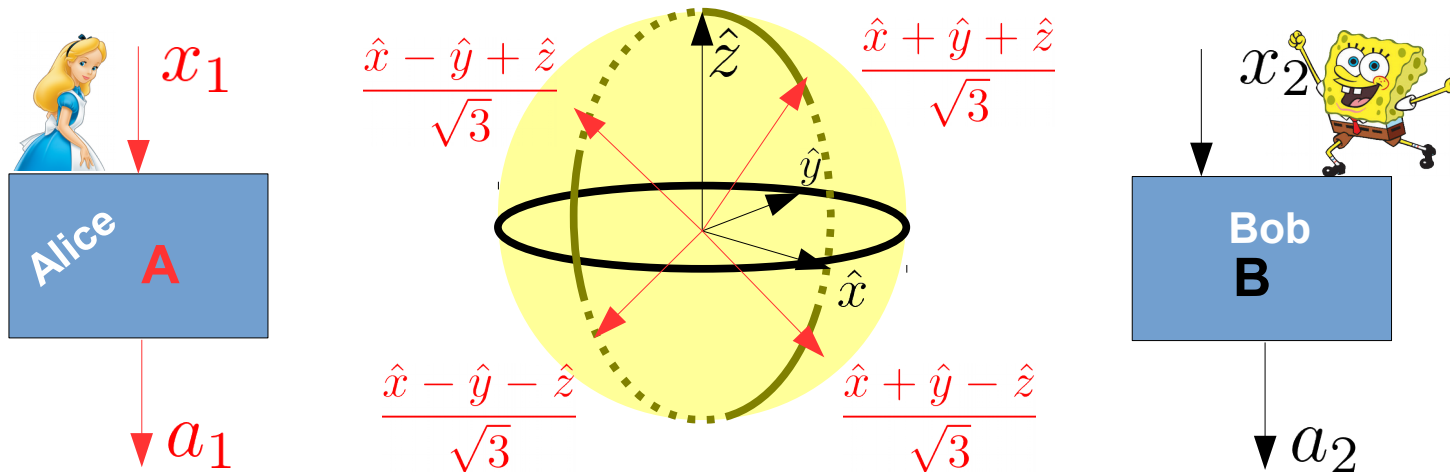
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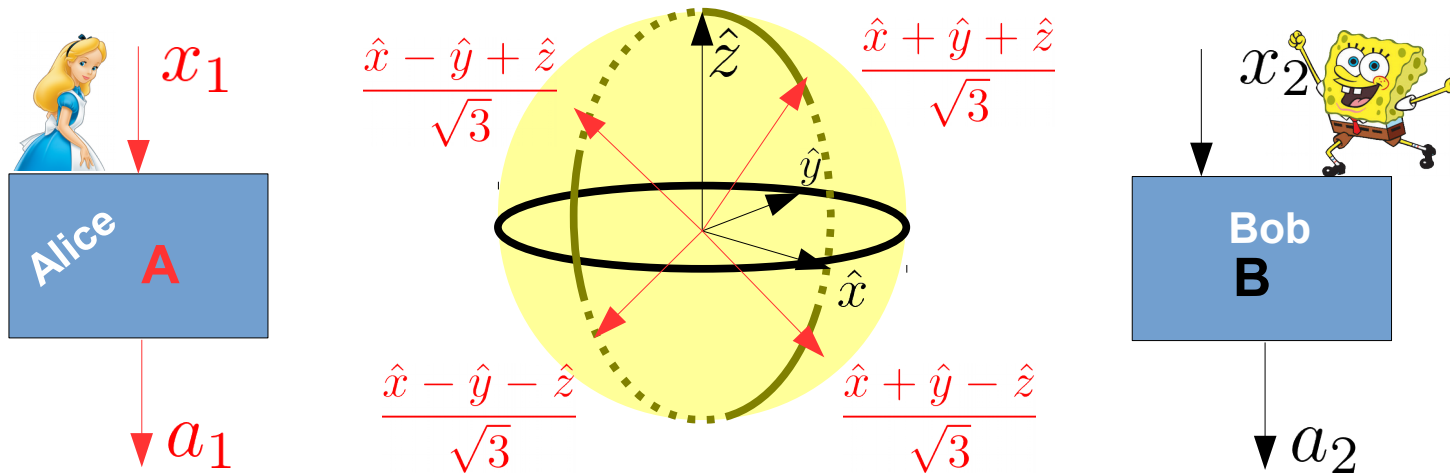
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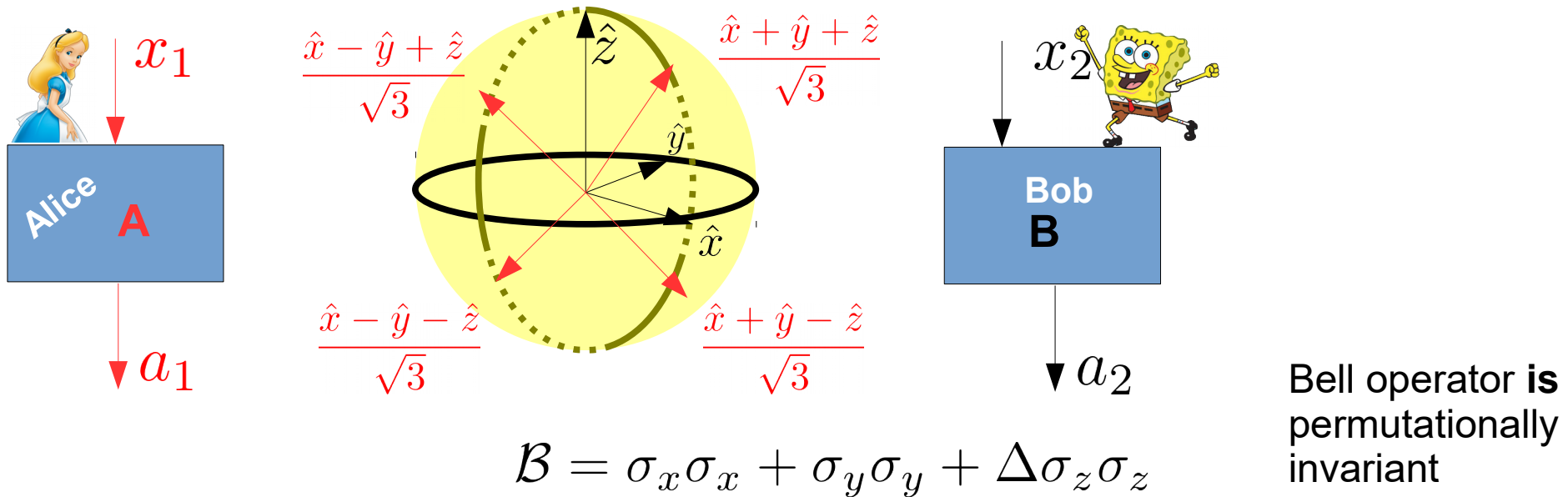


$$\mathcal{B} = \sigma_x \sigma_x + \sigma_y \sigma_y + \Delta \sigma_z \sigma_z$$

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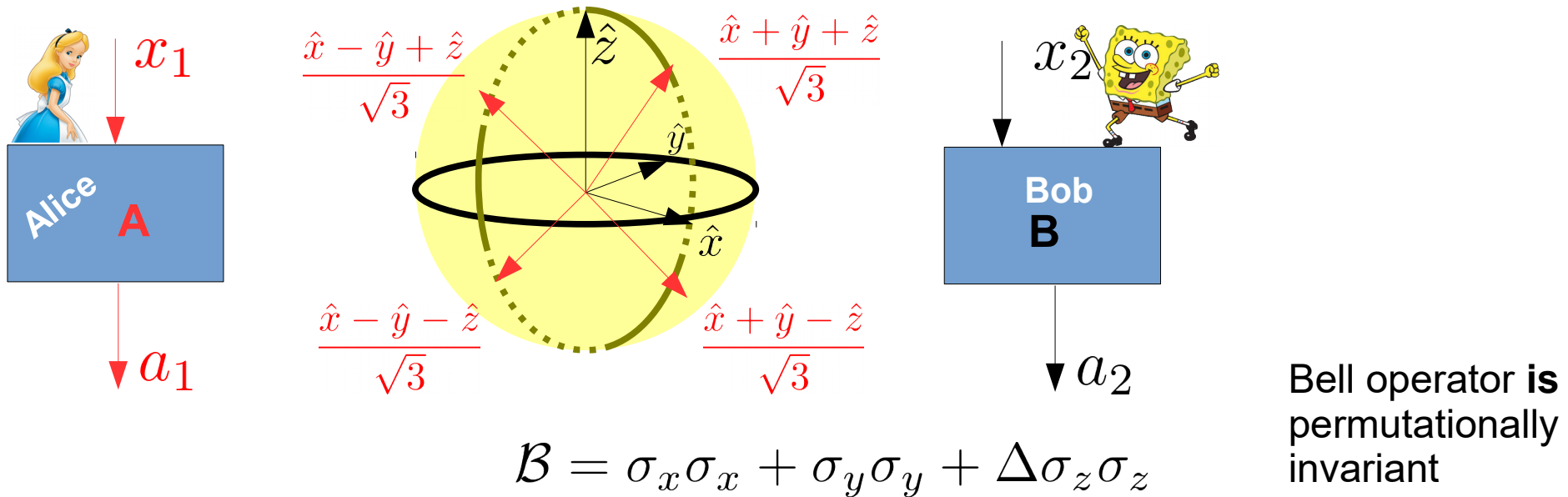
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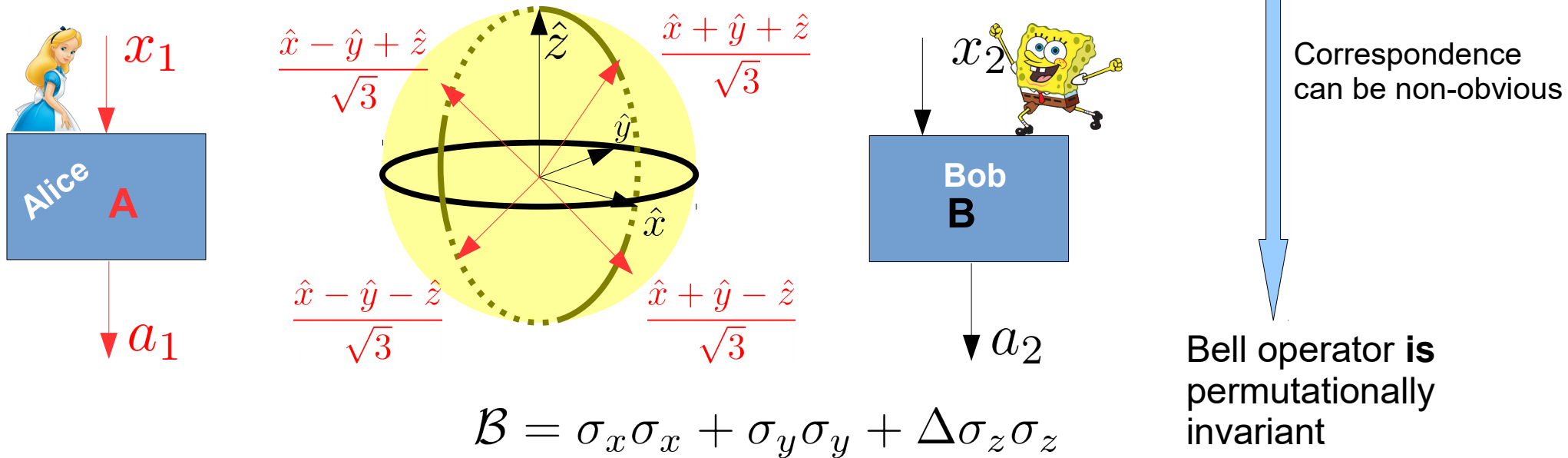
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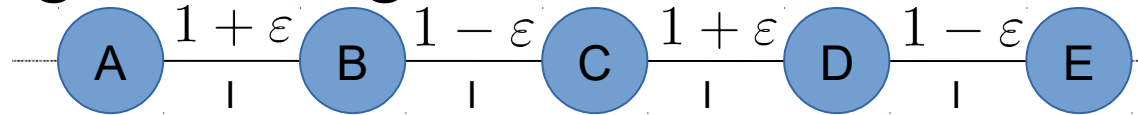


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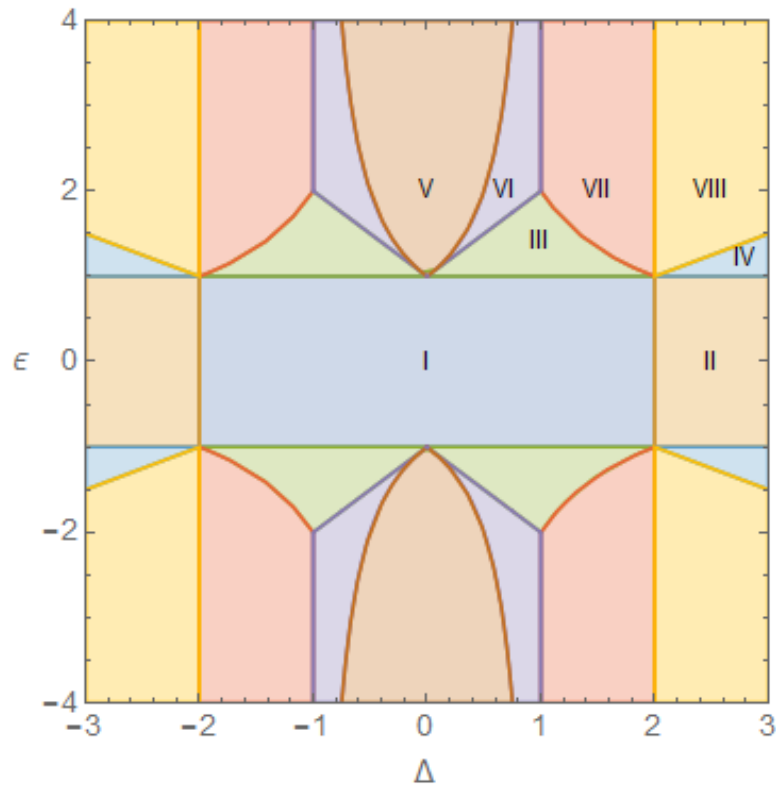
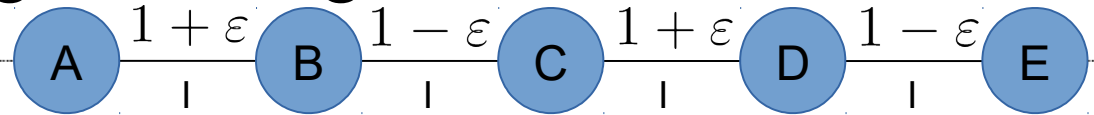
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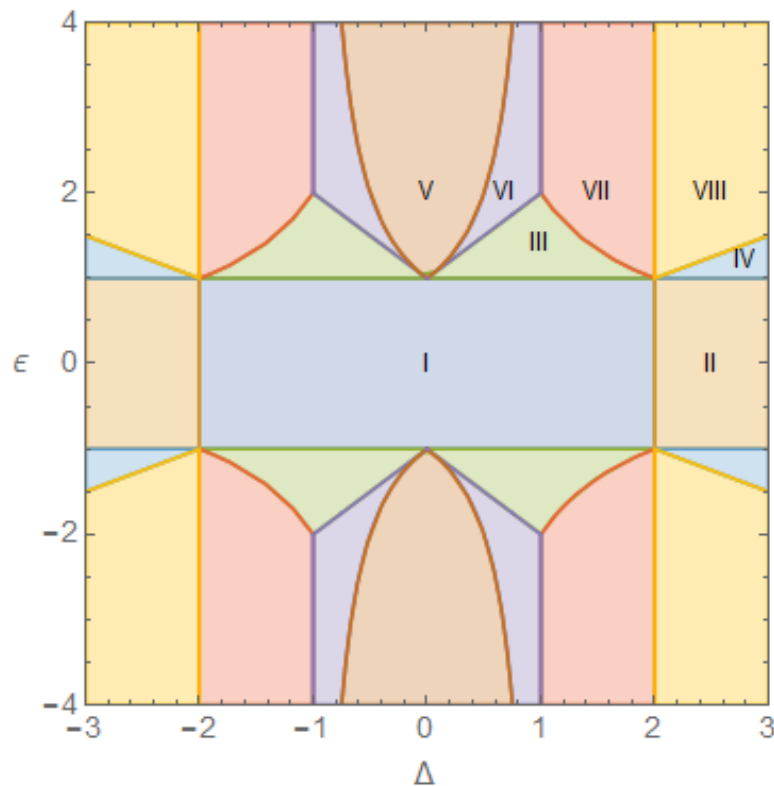
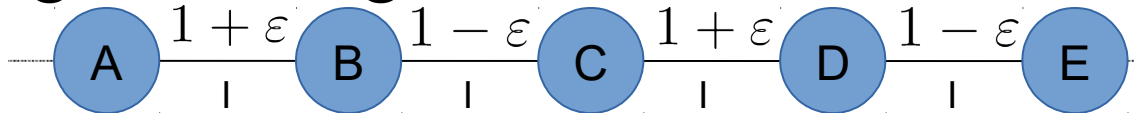
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$$\begin{aligned}
 \beta_{C,I} &= -n(4 + 2|\Delta|) \\
 \beta_{C,II} &= -4n|\Delta| \\
 \beta_{C,III} &= -8 - 4|\Delta| - (4n - 8)|\epsilon| - (2n - 4)|\Delta||\epsilon| \\
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 \beta_{C,V} &= -4n|\epsilon| - (2n - 8)|\epsilon||\Delta| \\
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 \beta_{C,VII} &= -4|\Delta| - (4n - 8)|\epsilon| - 2n|\epsilon||\Delta| \\
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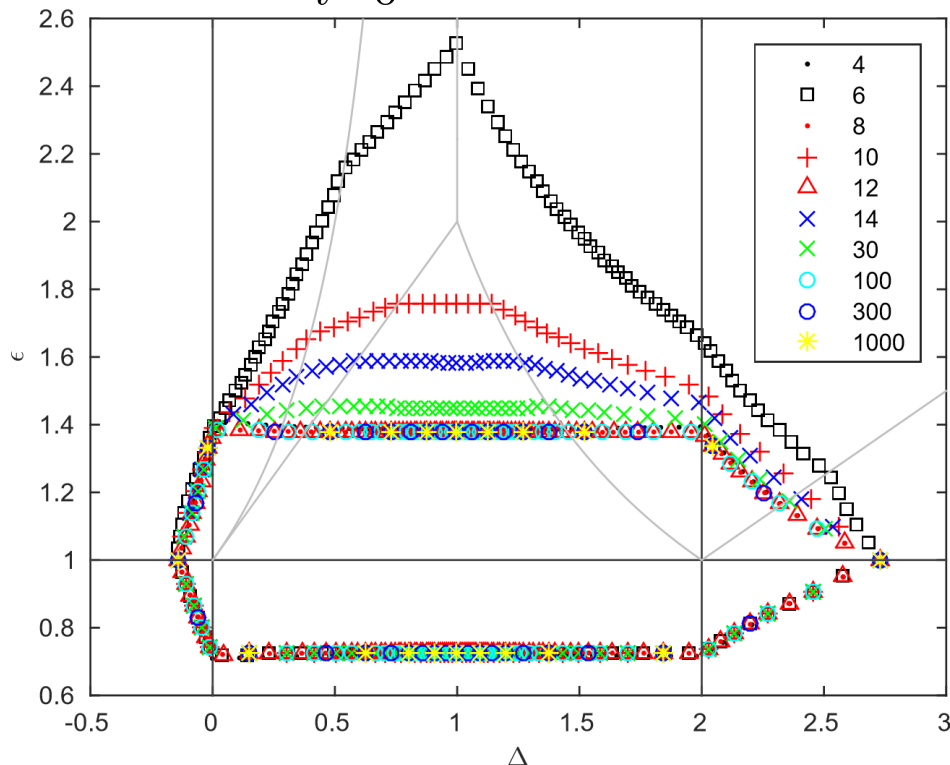
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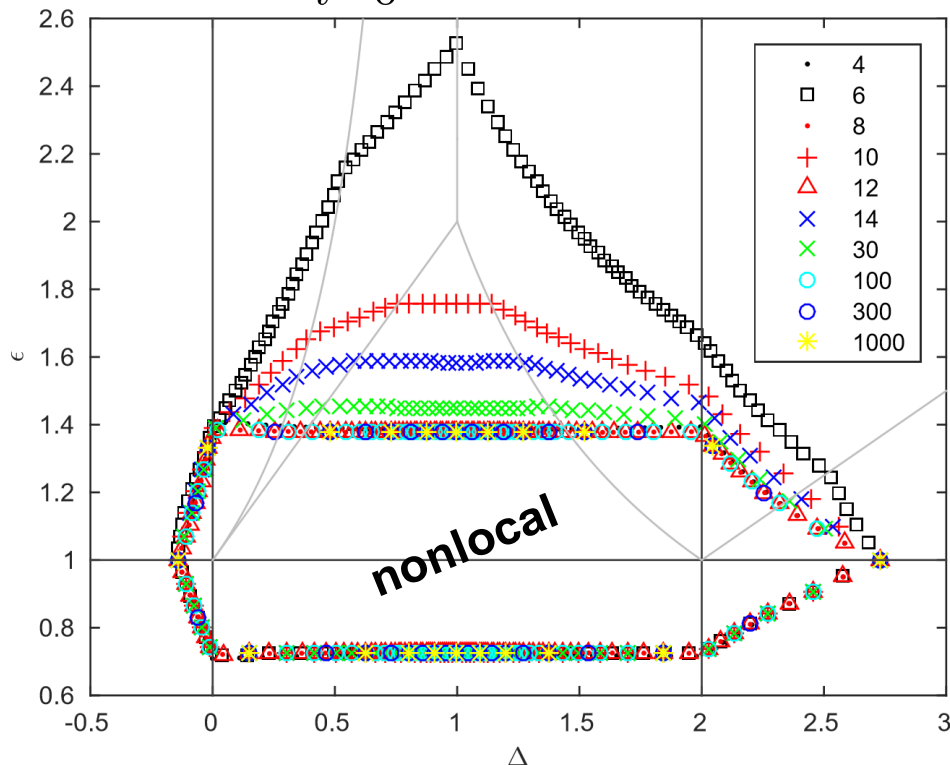
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- Study **persistence of nonlocality**



# Thanks for your attention!



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